The exciting physics with exotics

Ulrich Wiedner (Ruhr-University Bochum)

SFAIR-Meeting Uppsala, 10.11.2014

### Standard Model of FUNDAMENTAL PARTICLES AND INTERACTIONS

The Standard Model summarizes the current knowledge in Particle Physics. It is the quantum theory that includes the theory of strong interactions (quantum chromodynamics or QCD) and the unified theory of weak and electromagnetic interactions (electroweak). Gravity is included on this chart because it is one of the fundamental interactions even though not part of the "Standard Model."

#### FERMIONS

#### matter constituents spin = 1/2, 3/2, 5/2, ...

Leptons spin = 1/2			Quarks spin = 1/2		
Flavor	Mass GeV/c <sup>2</sup>	Electric charge	Flavor	Approx. Mass GeV/c <sup>2</sup>	Electric charge
$\nu_e$ electron neutrino	<1×10 <sup>-8</sup>	0	U up	0.003	2/3
e electron	0.000511	-1	<b>d</b> down	0.006	-1/3
$\nu_{\mu}$ muon neutrino	<0.0002	0	C charm	1.3	2/3
$\mu$ muon	0.106	-1	S strange	0.1	-1/3
$ u_{\tau}^{tau}_{neutrino}$	<0.02	0	t top	175	2/3
au tau	1.7771	-1	<b>b</b> bottom	4.3	-1/3

Spin is the intrinsic angular momentum of particles. Spin is given in units of h, which is the quantum unit of angular momentum, where  $\hbar = h/2\pi = 6.58 \times 10^{-25}$  GeV s =  $1.05 \times 10^{-34}$  J s.

Electric charges are given in units of the proton's charge. In SI units the electric charge of the proton is 1.60×10<sup>-19</sup> coulombs.

The energy unit of particle physics is the electronvolt (eV), the energy gained by one electron in crossing a potential difference of one volt. **Masses** are given in GeV/c<sup>2</sup> (remember  $E = mc^2$ ), where 1 GeV =  $10^9 \text{ eV} = 1.60 \times 10^{-10}$  joule. The mass of the proton is 0.938 GeV/c<sup>2</sup> = 1.67×10<sup>-27</sup> kg.

Baryons qqq and Antibaryons qqq Baryons are fermionic hadrons. There are about 120 types of baryons.					
Symbol Name Quark Content Electric Mass GeV/c <sup>2</sup> Spin					
р	proton	uud	1	0.938	1/2
p	anti- proton	ūūd	-1	0.938	1/2
n	neutron	udd	0	0.940	1/2
Λ	lambda	uds	0	1.116	1/2
Ω-	omega	SSS	-1	1.672	3/2

#### Matter and Antimatter

For every particle type there is a corresponding antiparticle type, denoted by a bar over the particle symbol (unless + or - charge is shown). Particle and antiparticle have identical mass and spin but opposite charges. Some electrically neutral bosons (e.g.,  $Z^0$ ,  $\gamma$ , and  $\eta_r = c\bar{c}$ , but not  $K^0 = ds$ ) are their own antiparticles.

#### Figures

These diagrams are an artist's conception of physical processes. They are not exact and have no meaningful scale. Green shaded areas represent the cloud of gluons or the gluon field, and red lines the quark paths.



then the quarks and electrons would be less than 0.1 mm in size and the entire atom would be about 10 km across.

#### 5

In	teraction	Gravit	Weak	Electromagnetic	Str	ong
operty	-	Gib	(Electr	-41	Fundamental	
Acts on:		Mass – Energy	Flavo	Electric Charge	Color Charge	See Residual Strong Interaction Note
Particles experience	ing:	All	Quarks Jeptons	Electrically charged	Quarks, Gluons	Hadrons
Particles mediation	ng:	(not yet observed)	₩ W- Z <sup>0</sup>	γ	Gluons	Mesons
ength relative to electromag	10 <sup>-18</sup> m	10 <sup>-41</sup>	0.8	1	25	Not applicable
two u quarks at:	3×10 <sup>−17</sup> m	10 <sup>-41</sup>	10 <sup>-4</sup>	1	60	to quarks
two protons in nucle	us	10 <sup>-36</sup>	10 <sup>-7</sup>	1	Not applicable to hadrons	20



#### The Particle Adventure

Visit the award-winning web feature The Particle Adventure at http://ParticleAdventure.org

This chart has been made possible by the generous support of:

U.S. Department of Energy U.S. National Science Foundation Lawrence Berkeley National Laboratory Stanford Linear Accelerator Center American Physical Society, Division of Particles and Fields BURLE INDUSTRIES, INC.

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#### http://CPEPweb.org

#### force carriers BOSONS

<b>Unified Electroweak</b> spin = 1			
Name	Mass Electri GeV/c <sup>2</sup> charge		
$\gamma$ photon	0	0	
W-	80.4	-1	
W+	80.4	+1	
Z <sup>0</sup>	91.187	0	

spin = 0, 1, 2, ...

troweak	spin = 1	Strong (cold	
Mass Electric GeV/c <sup>2</sup> charge		Name	N G
0	0	<b>g</b> gluon	
80.4	-1 Color Char		
80.4 +1		Each quark carries one "strong charge," also	

of three types of lled "color charge. These charges have nothing to do with the colors of visible light. There are eight possible types of color charge for gluons. Just as electri-

spin = 1

lass V/c<sup>2</sup>

0

Electric

charge

0

cally-charged particles interact by exchanging photons, in strong interactions color-charged particles interact by exchanging gluons. Leptons, photons, and W and Z bosons have no strong interactions and hence no color charge.

#### Quarks Confined in Mesons and Baryons

One cannot isolate quarks and gluons; they are confined in color-neutral particles called **hadrons**. This confinement (binding) results from multiple exchanges of gluons among the color-charged constituents. As color-charged particles (quarks and gluons) move apart, the energy in the color-force field between them increases. This energy eventually is converted into additional guark-antiguark pairs (see figure below). The guarks and antiguarks then combine into hadrons; these are the particles seen to emerge. Two types of hadrons have been observed in nature: mesons gg and baryons ggg.

#### **Residual Strong Interaction**

The strong binding of color-neutral protons and neutrons to form nuclei is due to residual strong interactions between their color-charged constituents. It is similar to the residual electrical interaction that binds electrically neutral atoms to form molecules. It can also be viewed as the exchange of mesons between the hadrons.

PROPERTI	ES OF TH	2	INTERACTION	
ion	Weak		Electromagnetic	

pe⁻⊽ <sub>e</sub>	e <sup>+</sup> e <sup>-</sup> → B <sup>0</sup> B <sup>0</sup>
d w- e-	e <sup>+</sup> v e <sup>-</sup> z

n ->

W boson. This is neutron B decay.

An electron and positron A neutron decays to a proton, an electron, (antielectron) colliding at high energy can and an antineutrino via a virtual (mediating) annihilate to produce B<sup>0</sup> and B<sup>0</sup> mesons via a virtual Z boson or a virtual photon



Two protons colliding at high energy can produce various hadrons plus very high mass particles such as Z bosons. Events such as this one are rare but can vield vital clues to the structure of matter.

hadrons

Z<sup>0</sup>

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Flavor

v muon

v tau



### BOSONS



force carriers spin = 0, 1, 2, ...



Each quark carries one of three types of "strong charge," also called "color charge." These charges have nothing to do with the colors of visible light. There are eight possible types of color charge for gluons. Just as electri-

cally-charged particles interact by exchanging photons, in strong interactions color-charged particles interact by exchanging gluons. Leptons, photons, and W and Z bosons have no strong interactions and hence no color charge.

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One cannot isolate guarks and gluons; they are confined in color-neutral particles called hadrons. This confinement (binding) results from multiple exchanges of gluons among the color-charged constituents. As color-charged particles (quarks and gluons) move apart, the energy in the color-force field between them increases. This energy eventually is converted into additional quark-antiquark pairs (see figure below). The quarks and antiquarks then combine into hadrons; these are the particles seen to emerge. Two types of hadrons have been observed in nature: **mesons**  $q\bar{q}$  and **baryons** qqq.

and



#### Positronium





## X and Y mesons



GeV/c<sup>2</sup>

M(D\*D\*)

 $M(\Lambda_{c}\Lambda_{c})$ 

M(DD\*) Ulrich Wiedner

GeV/c<sup>2</sup>





	x(3872) molecule		Sign in
 Search	About 298,000 results (0.22 seconds) $\leftarrow About \ 29800$	0 results (0.2	22 seconds)
Everything	The X(3872) particle - The DZero Experiment - Fermilab www-d0.fnal.gov/Run2Physics/WWW/results/final/B//B04A.htm		
Images Maps	Apr 15, 2004 – The <b>X</b> ( <b>3872</b> ) particle What is it? April 15 Some theories have predicted that the <b>X</b> ( <b>3872</b> ) is a new type of particle called a meson- <b>molecule</b> .		
Videos	[PDF] D (2700), D (2860) and the open-charm system X(3872): molecu.		
News	web.na.infn.it/fileadmin/b-physics-workshop-2//colangelo.pdf File Format: PDF/Adobe Acrobat - Quick View		
Shopping	sJ. (2860) and the open-charm system. <b>X</b> ( <b>3872</b> ): <b>molecule</b> vs charmonium with Fulvia De Fazio, Rossella Ferrandes, Floriana Giannuzzi and Stefano Nicotri	I	
More	X (3872) as a DD* molecule bound by quark exchange forces		
Philadelphia, PA       arxiv.org > hep-ph         Dec 31, 2011 - Abstract: The Bethe-Salpeter equation for the T-Matrix of D-D* scattering is solved with a meson-meson potential that results from 2nd order		ing	
Show search tools	The X (3872) boson: Molecule or charmonium arxiv.org > hep-ph by M Suzuki - 2005 - Cited by 103 - Related articles Aug 24, 2005 – Abstract: It has been argued that the mystery boson X(3872) is a molecule state consisting of primarily D0-D0*bar + D0bar-D*0. In contrast		
	Spin-parity analysis of the X(3872) « A Quantum Diaries Survivor dorigo.wordpress.com/2006/06//spin-parity-analysis-of-the-x3872/ Jun 9, 2006 – Two possible spin-parity assignments of the X(3872) are equally probabl in particular, the X may be indeed a <b>molecular</b> bound state of two	le:	
	Phys. Rev. D 72, 114013 (2005): X(3872) boson: Molecule or link.aps.org > Journals > Phys. Rev. D > Volume 72 > Issue 11 by M Suzuki - 2005 - Cited by 103 - Related articles Dec 19, 2005 – It has been argued that the mystery boson X(3872) is a molecule state consisting of primarily D0D *0+D 0D*0. In contrast, apparent puzzles	e	
	<u>Charm meson molecules and the X(3872)</u> drc.ohiolink.edu//7166?X(3872)1 Title: Charm meson molecules and the X(3872). Author: Kusunoki, Masaoki. Description: The recently discovered resonance X(3872) is interpreted as a		
	PROPERTIES OF X(3872) AS A HADRONIC MOLECULE WITH A www.worldscinet.com/ijmpcs/02/0201//S2010194511000857.pdf by M HARADA - Related articles		

We discuss the possible interpretation of V(2972) as a DD bedrapic melocyle with ID



 $\psi' {\rightarrow} \gamma X$ 



## The X(3872)



$$B \to KX; \ p\bar{p}$$
  

$$X \to \pi^{+}\pi^{-}J/\psi$$
  

$$X \to \pi^{+}\pi^{-}\pi^{0}J/\psi$$
  

$$X \to \gamma J/\psi; \ X \to \gamma \psi(2S)$$
  

$$X(3875) \to D^{0}\bar{D}^{0}\pi^{0}$$

$$J^{PC} = ? (1^{++})$$
  
M = 3871.68 ± 0.17 MeV  
 $\Gamma < 1.2$  MeV  
> 10  $\sigma$ 





Observed decay  $X(3872) \rightarrow \gamma J/\psi : \Rightarrow C=+$ 



## What is the nature of the X(3872)?



This agrees more with models favoring a charmonium state or a mixture of a charmonium/molecular state solutions than with a pure molecular interpretation.

... but still likely to be exotic:

Di-pion mass is dominated by the  $\rho(770) \Rightarrow I=1$  BELLE and BaBar see decay  $X(3872 \rightarrow J/\psi\omega)$  Ratio ~ 1  $\blacktriangleright$  huge isospin violation



How to progress further in the understanding of the new states? It is important to determine the resonance curve precisely ...



The line shapes for virtual state and bound state are the same *above* threshold but differ dramatically *below* threshold.

Analysis of  $J/\psi \, \pi^+\pi^-$  and  $D^0 \bar{D}^0 \pi^0$  Decays of the X(3872)

Eric Braaten and James Stapleton

Physics Department, Ohio State University, Columbus, Ohio 43210, USA (Dated: July 17, 2009)

Phys.Rev. D81 (2010) 014019

## The PANDA Detector





# Resonance scan with varying p momentum at PANDA (possible for states with all quantum numbers)



Measure rate of final state under study:

$$\mathbf{R}_{i} = \mathbf{L}_{0} \bullet \boldsymbol{\sigma}(\mathbf{p}_{i}) \bullet \mathbf{K} \ (\Delta \mathbf{p}/\mathbf{p}, |\mathbf{p}_{i} - \mathbf{p}_{R}|)$$

(K takes overlap between beam and resonance into account)



## PANDA reconstruction of X(3872) mass and width



### **BELLE II:**

An advanced width determination using measured masses, the beam energy and momentum of the B allowed BELLE to reduce the width determination from originally <2.4 GeV to <1.2 GeV in a 3-dimensional fit (Phys. Rev. D84(2011)052004).

With the higher statistics (350 events expected for the X(3872) in 2020) in BELLE II  $\Gamma$ <110 keV might be achievable (S. Lange - FAIR conference Worms 2014).

A precise determination of the resonance curve form is not possible in BELLE.

## The Y story

Using ISR (Initial State Radiation) to find states:



#### The respective cross sections:

BELLE: PRL110, 252002 (2013). BaBar: PRD86, 051102 (2013). 120 **(b)** 80 (b) 100 70  $\sigma(e^+e^- \rightarrow J/\psi \pi^+\pi)(pb)$ 60  $\sigma(\pi^{+}\pi^{-}J/\psi)$  (pb) 80 50 60 40 30 40 20 10 0 0 4.2 4.8 5.2 5.4 3.8 4.4 4.6 5 3.8 5.2 5.4 5 8 Ecm(GeV) E<sub>cm</sub> (GeV)

Y(4260):  $M \approx 4260$  MeV,  $\Gamma \approx 100$  MeV

Conventional wisdom and potential models: charmonia above threshold decay to open charm  $\psi(4040), \psi(4160), \psi(4415)$ 

The Y(4260) has a large decay width to  $\pi^+\pi^- J/\psi$ 



**BES**II

$$e^+e^- \rightarrow \gamma \pi^+ \pi^- J/\psi$$



#### Observation of the X(3872) in the radiative decay of the Y(4260)



BESIII: PRL112, 092001 (2014).

What does nature want to tell us?

Additional members of the ISR Y family seen by BELLE and BaBar:



## The first obvious exotic: Z(4430)

Z<sup>+</sup> (4430) - a new state of matter (tetraquark) decaying into  $\pi^+\psi'$ 



 $M = (4.433 \pm 0.004 \text{ (stat)} \pm 0.001 \text{ (syst)}) \text{ GeV}$   $\Gamma = (0.044^{+0.017}_{-0.01} \text{ (stat)}^{+0.030}_{-0.01} \text{ (syst)}) \text{ GeV}$   $\mathcal{E}(B \to KZ(4430) \times \mathcal{E}(Z \to \pi^{+}\psi') = (4.1 \pm 1.0 \text{ (stat)} \pm 1.3 \text{ (syst)}) \times 10^{-5}$ PRL 10

PRL 100, 142001 (2008) arXiv:0708.1790 [hep-ex]

### Confirmation by LHCb

![](_page_28_Figure_1.jpeg)

M =  $4475 \pm 7_{-25}^{+15}$  MeV  $\Gamma = 172 \pm 13_{-34}^{+37}$  MeV Significance: >13.9 $\sigma$ 

PANDA:  $\overline{p}p \rightarrow Z^+(4430) + \pi^-$ 

![](_page_29_Picture_1.jpeg)

![](_page_29_Figure_2.jpeg)

![](_page_30_Figure_0.jpeg)

$\sqrt{s}$ (GeV)	$\mathcal{L}$ (pb <sup>-1</sup> )
3.900	52.8
4.009	482.0
4.090	51.0
4.190	43.0
4.210	54.7
4.220	54.6
4.230	1090.0
4.245	56.0
4.260	826.8
4.310	44.9
4.360	544.5
4.390	55.1
4.420	44.7

### $Y(4260) \rightarrow \pi^+\pi^- J/\psi \rightarrow \pi^+\pi^- e^+ e^- (\pi^+\pi^-\mu^+\mu^-)$ : Straightforward analysis with 4 tracks

![](_page_31_Figure_1.jpeg)

![](_page_32_Picture_0.jpeg)

Observation of the Z<sub>c</sub>(3900) in BESIII

![](_page_32_Figure_2.jpeg)

![](_page_33_Picture_0.jpeg)

Observation of Z<sub>c</sub>(3900) at BESIII

![](_page_33_Figure_2.jpeg)

![](_page_34_Picture_0.jpeg)

Observation of  $Z_c(3885)$  in  $e^+e^- \rightarrow \pi^- (D^*D)^+$ 

![](_page_34_Figure_2.jpeg)

Phys. Rev. Lett 112, 022001 (2014) / 1310.1163

#### M = $3883.9 \pm 1.5 \pm 4.2$ MeV; $\Gamma = 24.8 \pm 3.3 \pm 11.0$ MeV

 $Z_{c}(3885) = Z_{c}(3900) \text{ but large yield of } \sim 6 \text{ for } \frac{\Gamma(DD^{*})}{\Gamma(\pi^{\pm}J/\psi)}$ 

![](_page_35_Picture_0.jpeg)

 $e^+e^- \rightarrow \pi Z_c(4020) \rightarrow \pi^+\pi^-h_c$ 

BESIII: 1309.1896

![](_page_35_Figure_3.jpeg)

Simultaneous fit to 4.23/4.26/4.36 GeV data,  $16 \eta_c$  decay modes:

M =  $4022.9 \pm 0.8 \pm 2.7 \text{ MeV/c}^2$   $\Gamma = 7.9 \pm 2.7 \pm 2.6 \text{ MeV}$ 

![](_page_36_Picture_0.jpeg)

## $e^+e^- \rightarrow \pi Z_c(4025) \rightarrow \pi^- (D^*\overline{D}^*)^+$

![](_page_36_Figure_2.jpeg)

BESIII: 1308.2760

Fit to  $\pi^{\pm}$  recoil mass yields  $401\pm47 \ Zc(4025) \ events \implies >10\sigma$   $M(Z_c(4025)) = 4026.3\pm2.6\pm3.7 \ MeV$  $\Gamma(Z_c(4025)) = 24.8\pm5.6\pm7.7 \ MeV$ 

![](_page_37_Picture_0.jpeg)

$$e^+e^- \rightarrow \pi Z_c(4025) \rightarrow \pi^- (D^*\overline{D}^*)^+$$

BESIII: 1308.2760

![](_page_37_Figure_3.jpeg)

Fit to  $\pi^{\pm}$  recoil mass yields 401±47 Zc(4025) events  $\Rightarrow >10\sigma$ M(Z<sub>c</sub>(4025)) = 4026.3±2.6±3.7 MeV;  $\Gamma$ (Z<sub>c</sub>(4025)) = 24.8±5.6±7.7 MeV

$$R = \frac{\sigma(e^+e^- \to \pi^{\pm}Z_c^{\mp}(4025) \to \pi^{\pm}(D^*\overline{D}^*)^{\mp})}{e^+e^- \to \pi^{\pm}(D^*\overline{D}^*)^{\mp}} = (65 \pm 9 \pm 6)\%$$

![](_page_38_Picture_0.jpeg)

### Notes from the Editors: Highlights of the Year

Published December 30, 2013 | Physics 6, 139 (2013) | DOI: 10.1103/Physics.6.139

#### Physics looks back at the standout stories of 2013.

As 2013 draws to a close, we look back on the research covered in *Physics* that really made waves in and beyond the physics community. In thinking about which stories to highlight, we considered a combination of factors: popularity on the website, a clear element of surprise or discovery, or signs that the work could lead to better technology. On behalf of the *Physics* staff, we wish everyone an excellent New Year.

![](_page_38_Picture_5.jpeg)

Images from popular Physics stories in 2013.

- Matteo Rini and Jessica Thomas

#### Four-Quark Matter

Quarks come in twos and threes—or so nearly every experiment has told us. This summer, the BESIII Collaboration in China and the Belle Collaboration in Japan reported they had sorted through the debris of high-energy electron-positron collisions and seen a mysterious particle that appeared to contain four quarks. Though other explanations for the nature of the particle, dubbed  $Z_c(3900)$ , are possible, the "tetraquark" interpretation may be gaining traction: BESIII has since seen a series of other particles that appear to contain four quarks.

#### Article Options

![](_page_38_Picture_11.jpeg)

#### **New in Physics**

Crisis Averted for the Bose Glass Synopsis | Jun 3, 2014

Unexpected Impact from Medium-Sized Solar Flare Synopsis | Jun 2, 2014

Scalable Imaging of Superresolution Viewpoint | Jun 2, 2014

Electrons Not the Cause of Charged Grains Focus | May 30, 2014

Seeing Just One Photon Synopsis | May 29, 2014

![](_page_39_Picture_0.jpeg)

Search for neutral partner of the  $Z_c(3900)$ :  $e^+e^- \rightarrow \pi^0\pi^0 J/\psi$ 

![](_page_39_Figure_2.jpeg)

![](_page_40_Picture_0.jpeg)

Search for neutral partner of the  $Z_c(3900)$ :  $e^+e^- \rightarrow \pi^0\pi^0 J/\psi$ 

![](_page_40_Figure_2.jpeg)

A new class of particles have been observed:

- At least 4-quarks
- Charged
- Close to DD thresholds
- They couple to DD final states larger than to charmonia

## 4-quark state

## D-D-"molecule"

![](_page_42_Figure_2.jpeg)

## Transition from color forces to colorless nuclear forces ?

## The future: PANDA

Proton-Antiproton contains already a 4-Quark-System

Idea: Dilepton-Tag from Drell-Yan-Production

### Advantages

- Trigger
- less J<sup>PC</sup>-Ambiguities
- 1200 E./day @ 12 GeV
- 300 E./day @ 5-8 GeV antiproton-Beam (for L=10<sup>32</sup>cm<sup>-2</sup>s<sup>-1</sup>)

![](_page_43_Picture_8.jpeg)

Bannikov, Gornuschkin, Kopeliovich, Krumshtein and Sapozhnikov, JINR E1-92-344 (1992)

## Other QCD states: Glueballs

A possible glueball spectrum

![](_page_44_Figure_2.jpeg)

## Glueballs $\rightarrow$ Creation of Mass

A few % of a hadron (proton) mass is generated due to the Higgs mechanism.

Most of the proton mass is created by the strong interaction.

**Glueballs gain their mass solely by the strong interaction** and are therefore an unique approach to the mass creation by the strong interaction.

The structure of glueballs

![](_page_46_Picture_1.jpeg)

Glueball (gg)

Are glueballs configurations of twisted or knotted colored flux?

![](_page_46_Picture_4.jpeg)

GLUEBALLS, FLUXTUBES AND η(1440). L. Fadeev, A. Niemi and U. Wiedner Phys.Rev.D70:114033, 2004

![](_page_47_Picture_0.jpeg)

## Hadron World

![](_page_47_Picture_2.jpeg)

meson

![](_page_47_Picture_4.jpeg)

glueball?

![](_page_48_Figure_0.jpeg)

Harvey B. Meyer, Michael J. Teper; Phys.Lett. B605 (2005) 344-354

Ulrich Wiedner

G. S. Bali et al.; arXiv:1302.1502

# Hadron physics is the place on earth to study non-Abelian massless gauge boson - gauge boson interaction in a controlled manner.

#### Feynman lectures on gravitation:

In fact, his work led to two sets of very useful results. The first, purely pedagogical, is embodied in the *Feynman Lectures on Gravitation* (publication [123]). In those lectures, Feynman develops the quantum field theory of a neutral massless spin 2 particle (the *graviton*), emphasizing the special features that arise, in comparison to theories of spin 0 and spin 1 particles, as well as the complications that result for a zero-mass particle in trying to create a self-consistent theory. As in the case of spin 1, masslessness results in redundant degrees of freedom, since Lorentz invariance requires that a *massless* particle can spin only along or opposite to its direction of momentum (positive or negative *chirality*), while a massive spin 2 particle may take up five different orientations relative to any arbitrary quantization direction. Eliminating the unwanted degrees of freedom is achieved by imposing certain "gauge conditions," which in the gravitational case brings about nonlinearity in the form of graviton-graviton interaction. Feynman shows that the classical limit of a properly gauged massless spin 2 theory is described by the Einstein gravitational field equations.<sup>3</sup>

### The difference between $e^+e^-$ colliders and PANDA

![](_page_50_Figure_1.jpeg)

## There exist also very narrow charged Z<sub>b</sub> states

$$\begin{split} M_1 &= 10607.2 \pm 2.0 \text{ MeV} \\ \Gamma_1 &= 18.4 \pm 2.4 \text{ MeV} \\ M_{Zb} &- (M_B + M_{B^*}) = +2.6 \pm 2.1 \text{ MeV} \end{split}$$

$$\begin{split} M_2 &= 10652.2 \pm 1.5 \text{ MeV} \\ \Gamma_2 &= 11.5 \pm 2.2 \text{ MeV} \\ M_{Zb} - (M_B + M_{B^*}) &= +1.8 \pm 1.7 \text{ MeV} \end{split}$$

![](_page_51_Figure_3.jpeg)

![](_page_52_Picture_0.jpeg)

![](_page_52_Picture_1.jpeg)

![](_page_52_Picture_2.jpeg)

![](_page_52_Picture_3.jpeg)

![](_page_53_Picture_0.jpeg)

![](_page_53_Picture_1.jpeg)

![](_page_53_Picture_2.jpeg)

![](_page_53_Picture_3.jpeg)

![](_page_54_Figure_0.jpeg)

Thank you very much for your attention.