

# SymTFTs with Interfaces in Geometric Engineering

Ling Lin

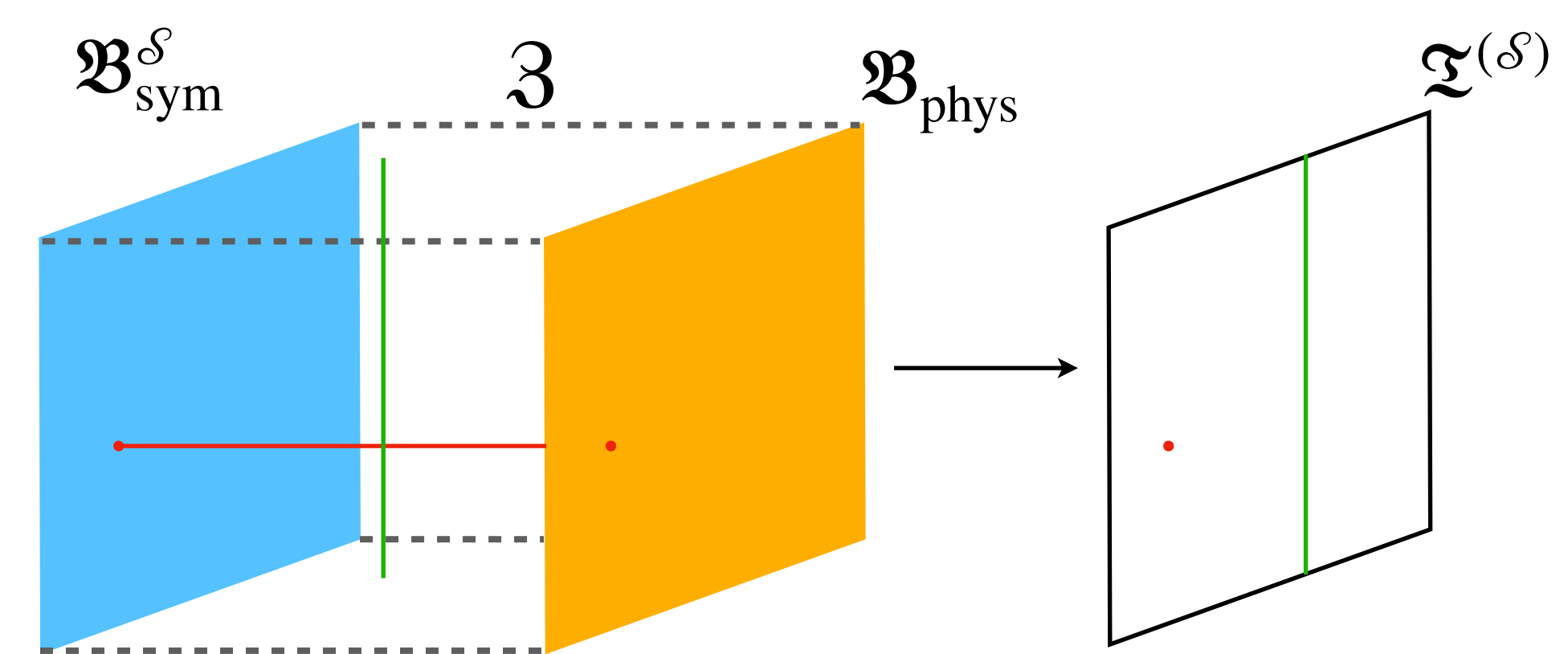
University and INFN Bologna



Work to appear with  
Daniel Robbins, Subham Roy

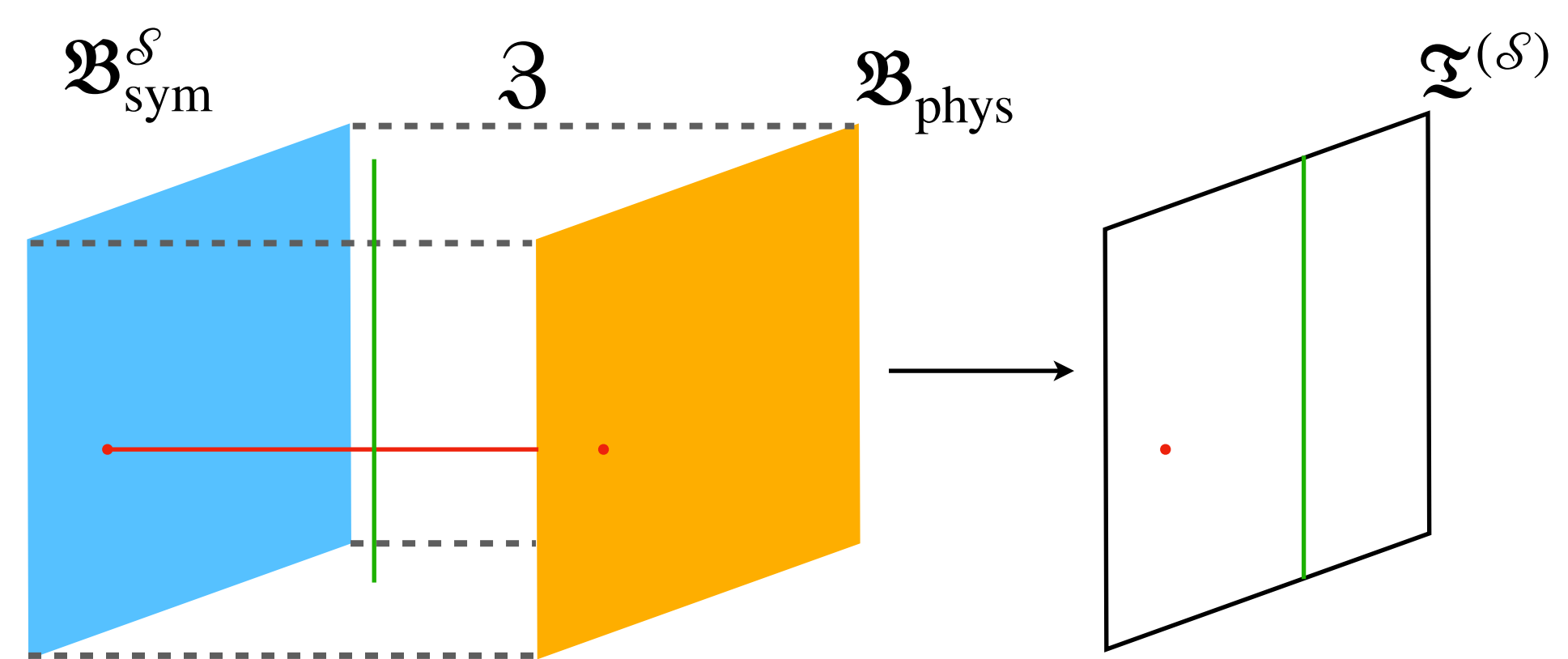
Strings & Geometry 2026  
Uppsala, May 19

# Symmetry TFT with interfaces



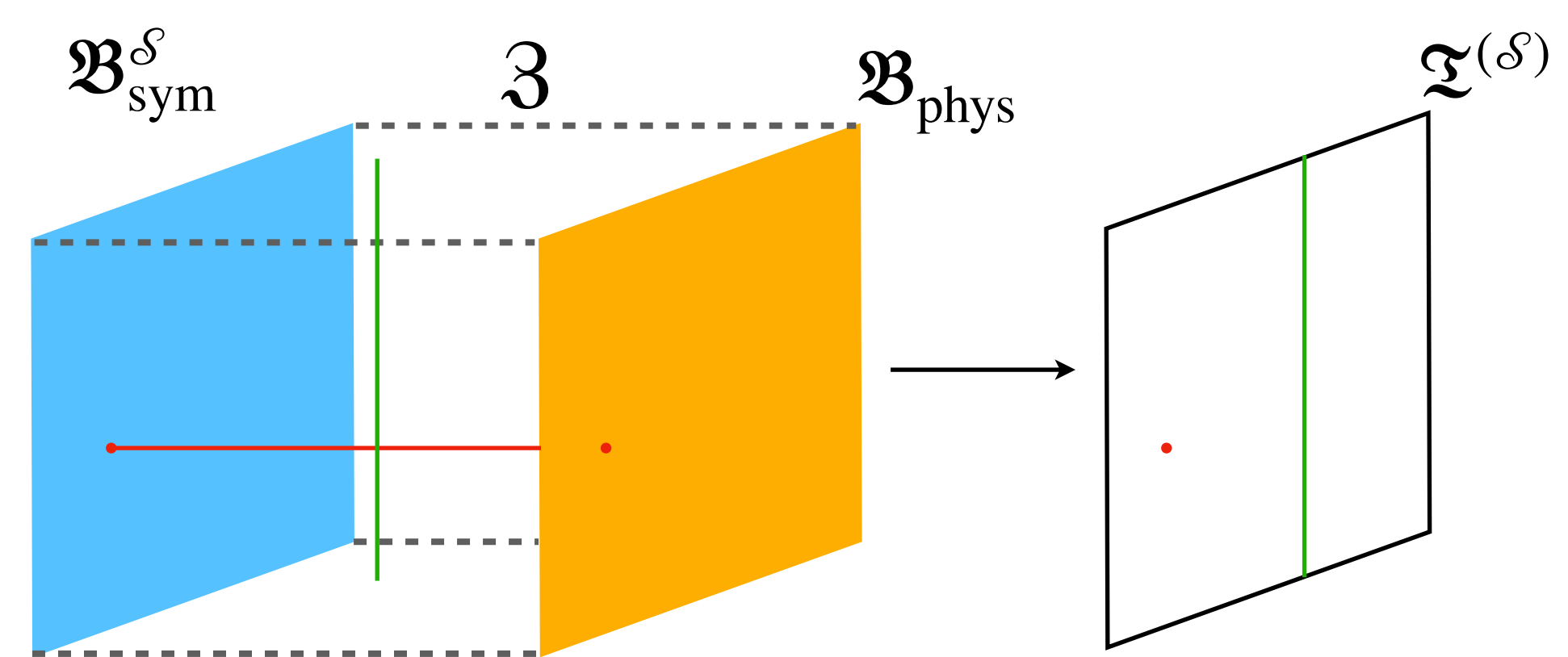
- Symmetry operators and charged defects in  $d$  dimensions can be encoded in a  $(d+1)$ -dim *Symmetry Topological Field Theory (SymTFT)* [Ji/Wen '19, Gaiotto/Kulp '20, Apruzzi/Bonetti/García Etxebarria/Hosseini/Schafer-Nameki '21, Freed/Moore/Teleman '22]

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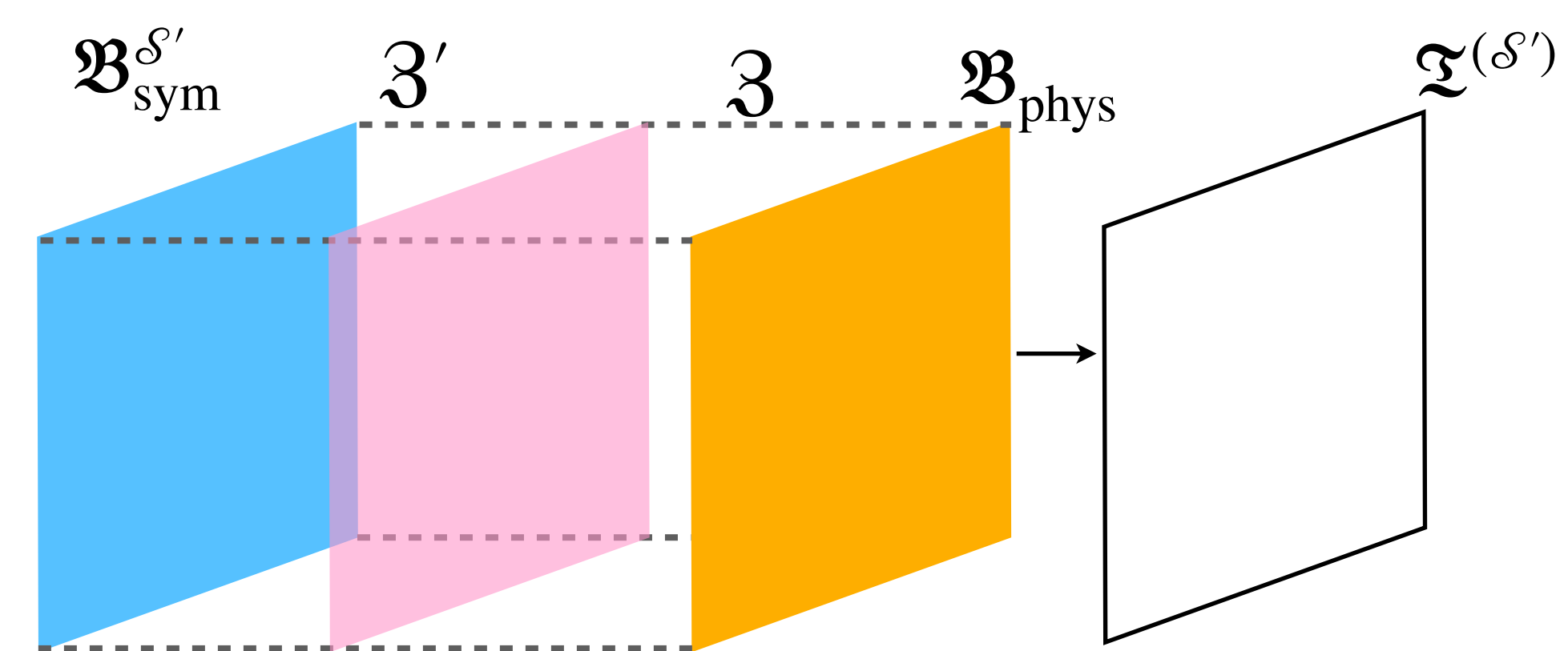
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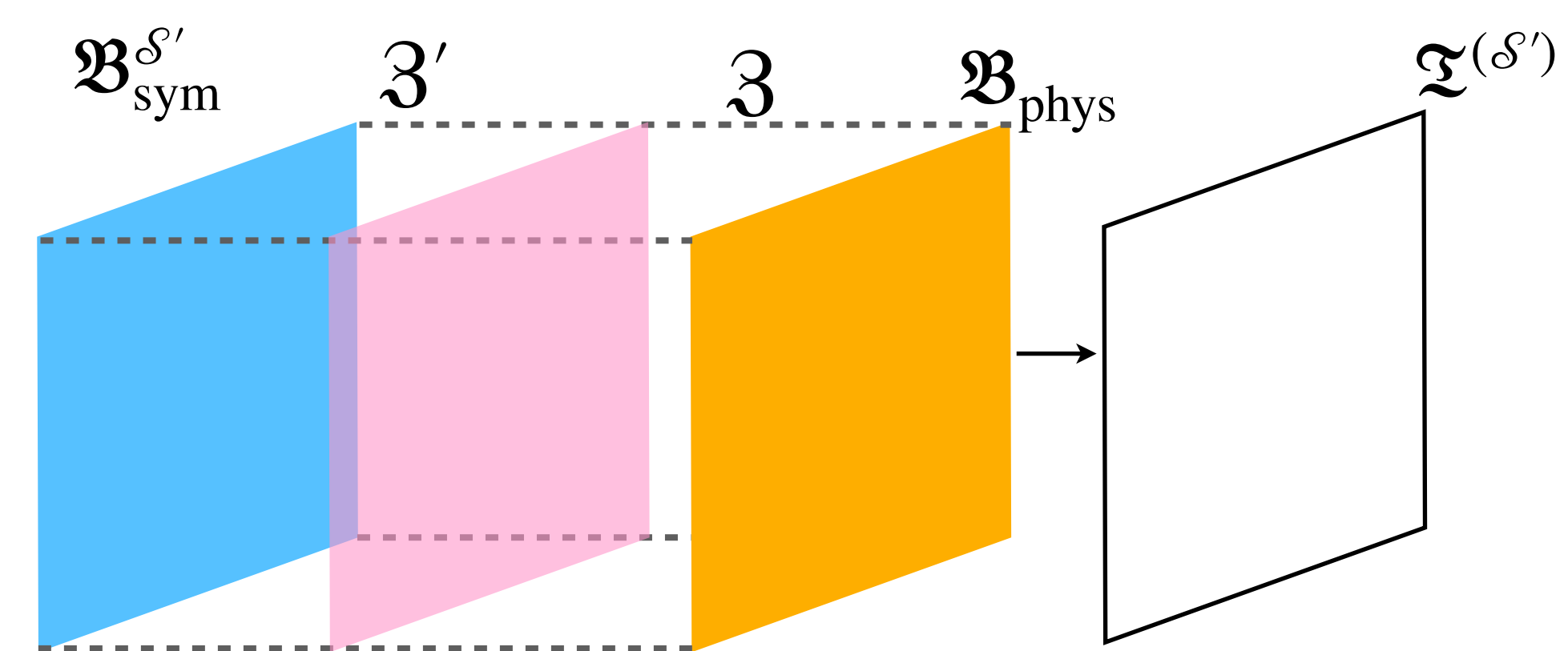
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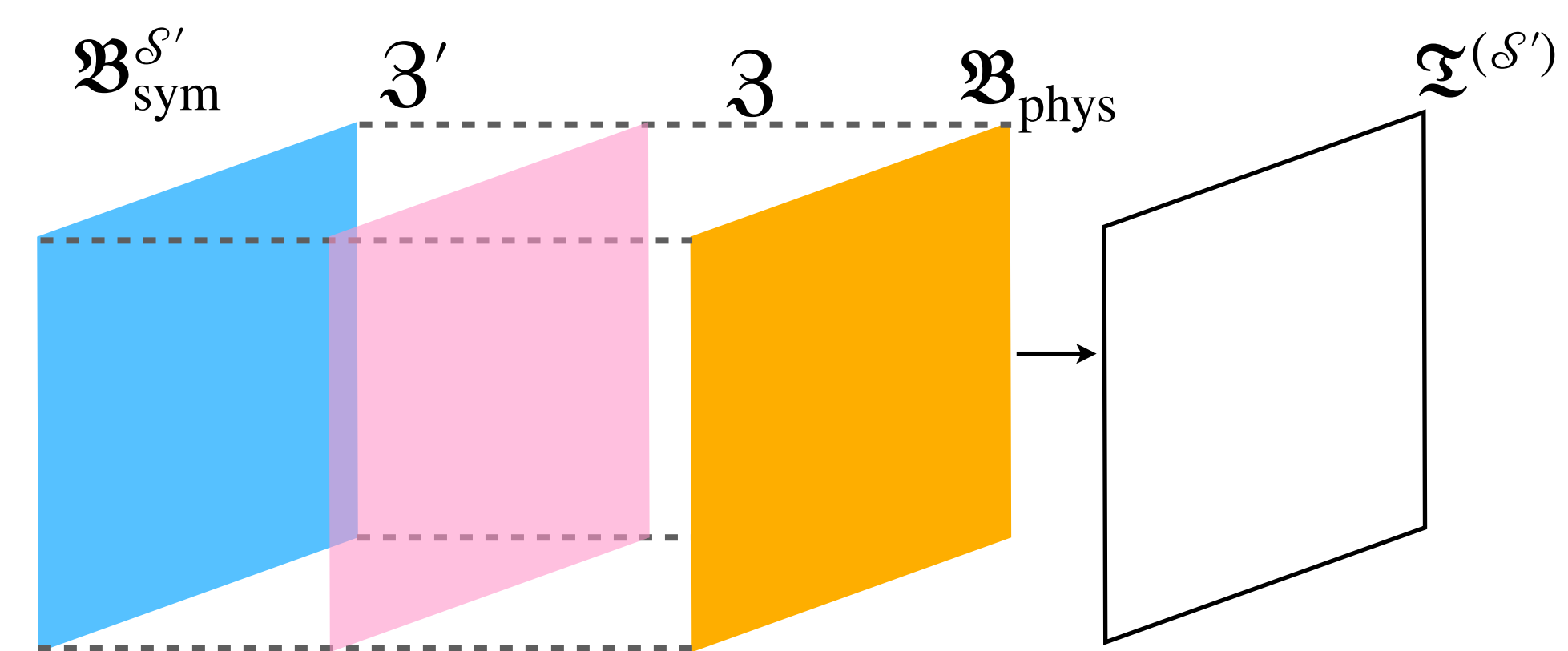
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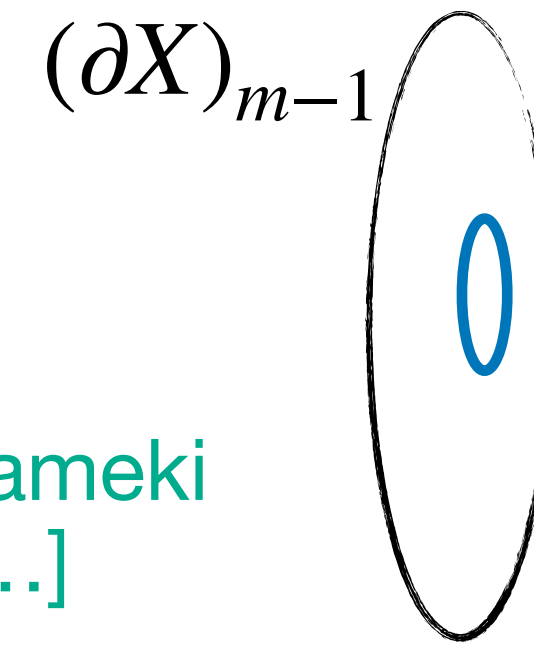
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- How are these features geometrically realized in string theory?

# Plan of the talk

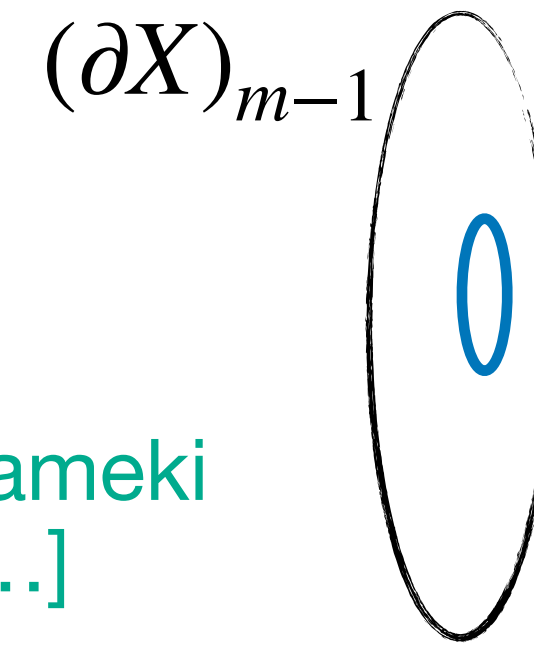
1. Topological SymTFT-interface in M-theory geometric engineering
2. Interfaces in SymTFT of “frozen”  $d \geq 6$   $\mathfrak{sp}(N)$  SYM
3. Summary & Outlook



# SymTFT in M-theory geometric engineering

[Apruzzi/Bonetti/García Etxebarria/Hosseini/Schafer-Nameki '21; Hübner/Morrison/Schafer-Nameki/Wang '22; Heckman/Hübner/Torres/Zhang '22; vanBeest/Gould/Schafer-Nameki/Wang '22, ...]

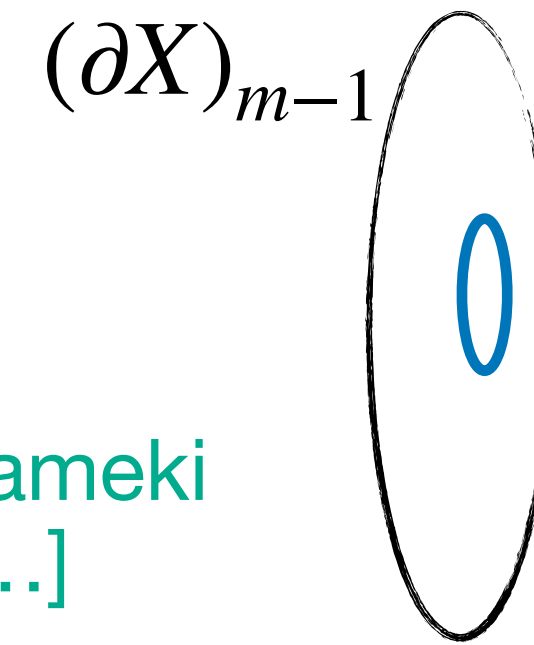
- Topological operators in  $W_{D+1=12-m}$  from “branes at infinity”:
  - ▶ M2s on  $H_1^{\text{tor}}(\partial X) \rightsquigarrow$  generators for magnetic (D-3)-form symmetry,
  - ▶ M5s on  $H_{m-3}^{\text{tor}}(\partial X) \rightsquigarrow$  generators for electric 1-form symmetry,
  - ▶ operator linking from geometric linking  $\ell : H_1^{\text{tor}}(\partial X) \times H_{m-3}^{\text{tor}}(\partial X) \rightarrow \mathbb{Q}/\mathbb{Z}$ .



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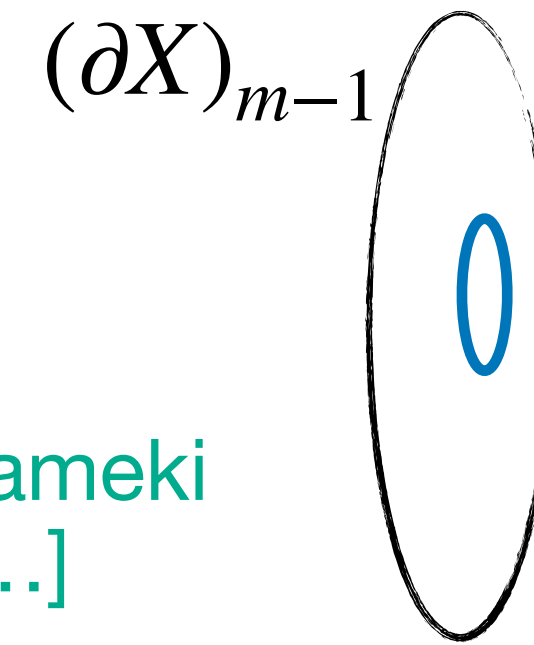
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- SymTFT action from dual cohomology:  $G_4 = B_2 e^{(2)}$ ,  $G_7 = B_{9-m} e^{(m-2)}$   
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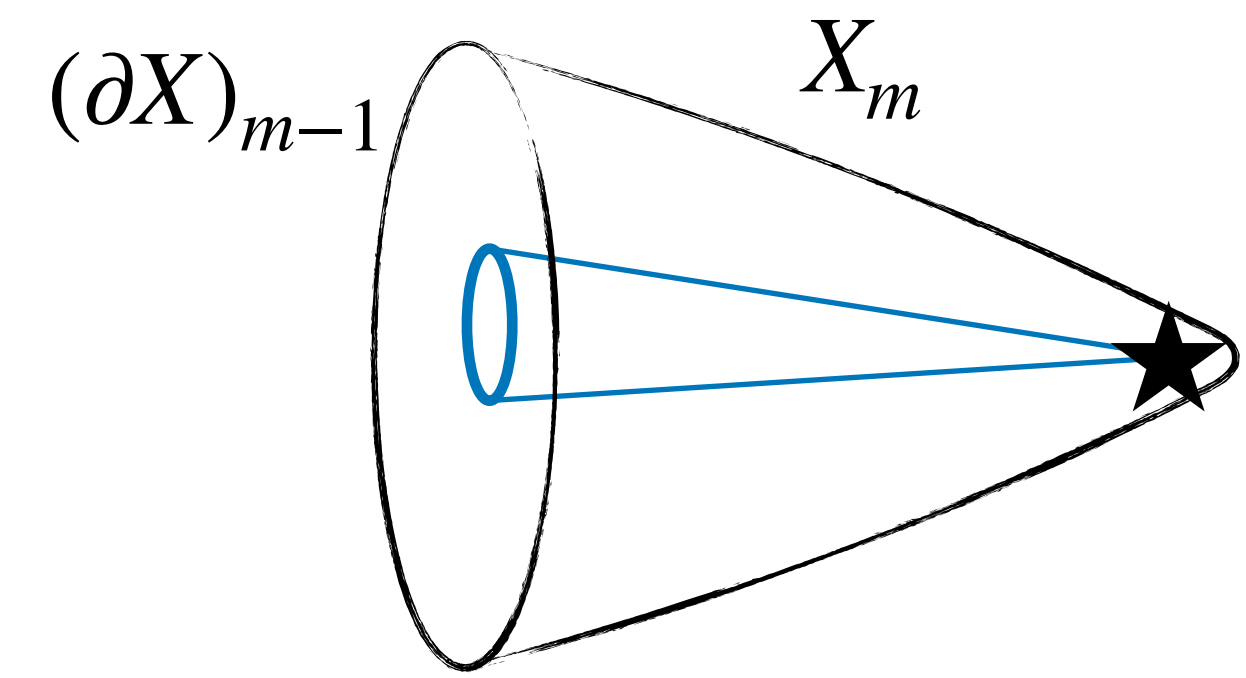
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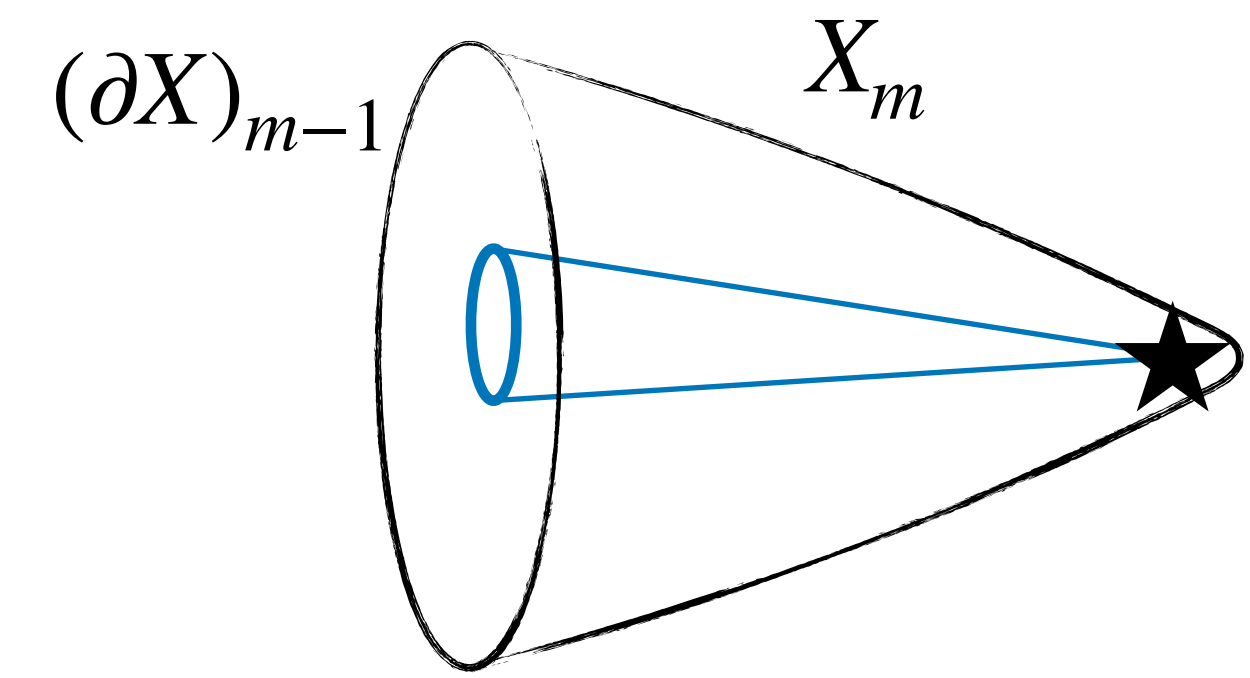
$$\int G_4 G_7 \rightsquigarrow \ell(e^{(2)}, e^{(m-2)}) \int_{W_{12-m}} B_2 \delta B_{9-m} = \frac{1}{N} \int_W B_2 \delta B_{9-m}$$

## Coupling to physical boundary...



- Dynamical objects in  $\mathfrak{B}_{\text{phys}}$  by wrapping  $H_p(X)$  (compact cycles)
- Probe objects from  $H_p(X, \partial X)/H_p(X) \equiv H_p(X, \partial X)/\text{im}(j) = \text{coker}(j)$

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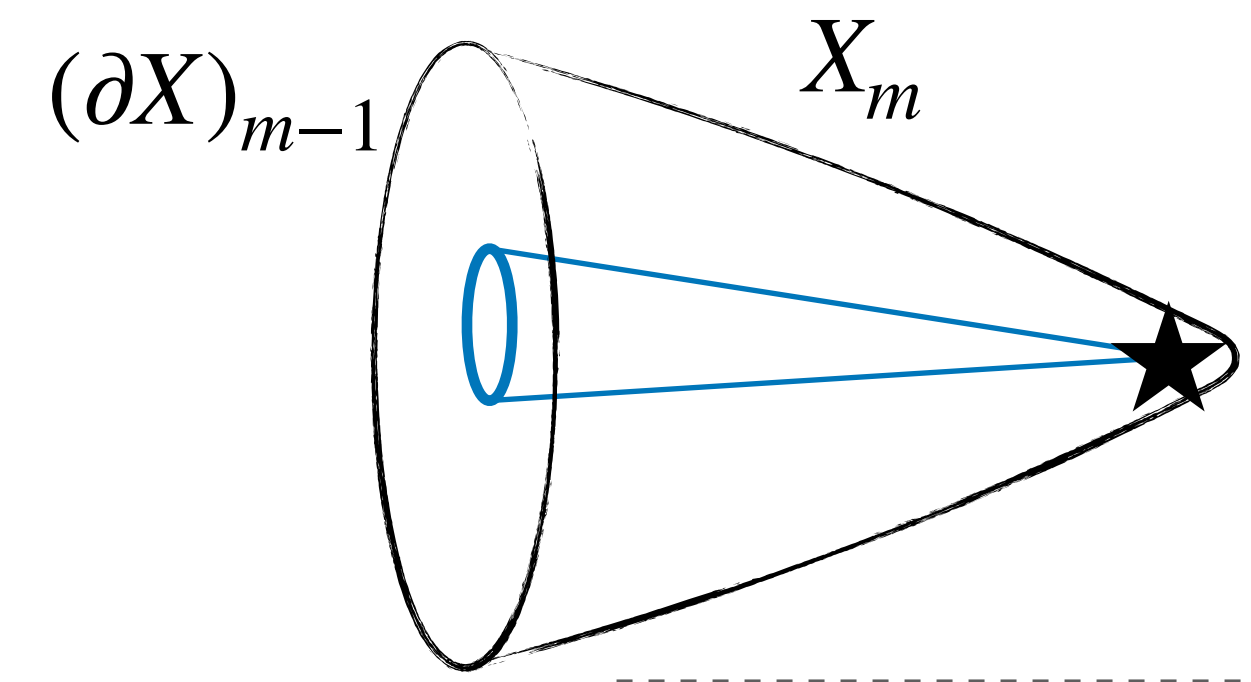
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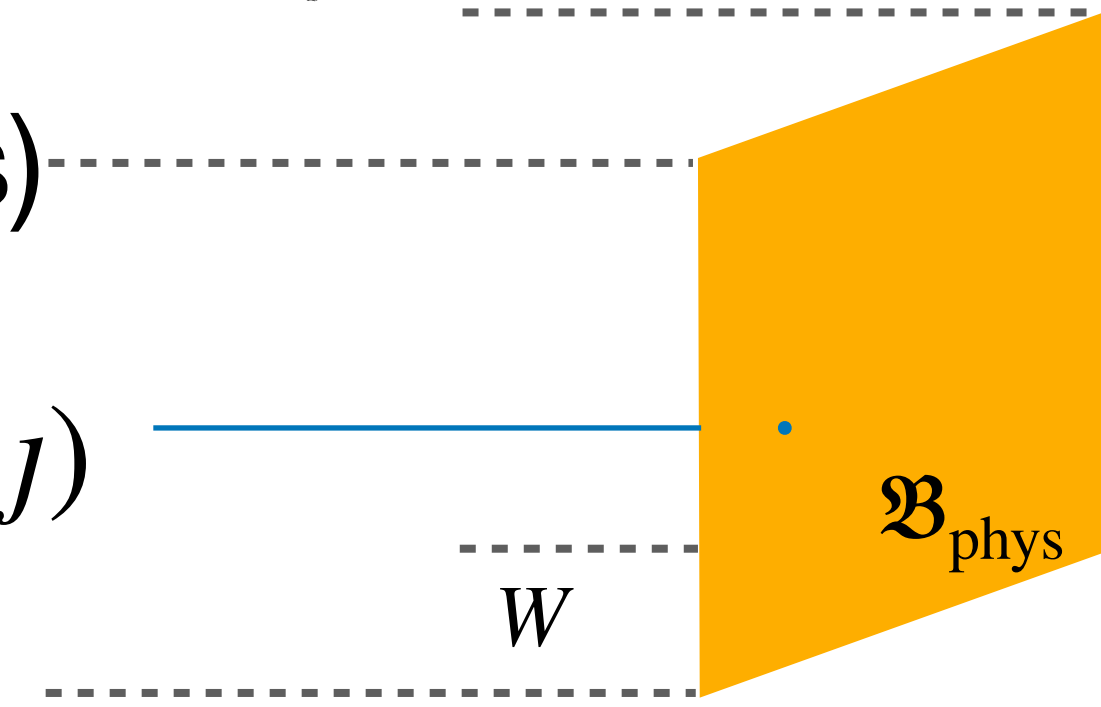


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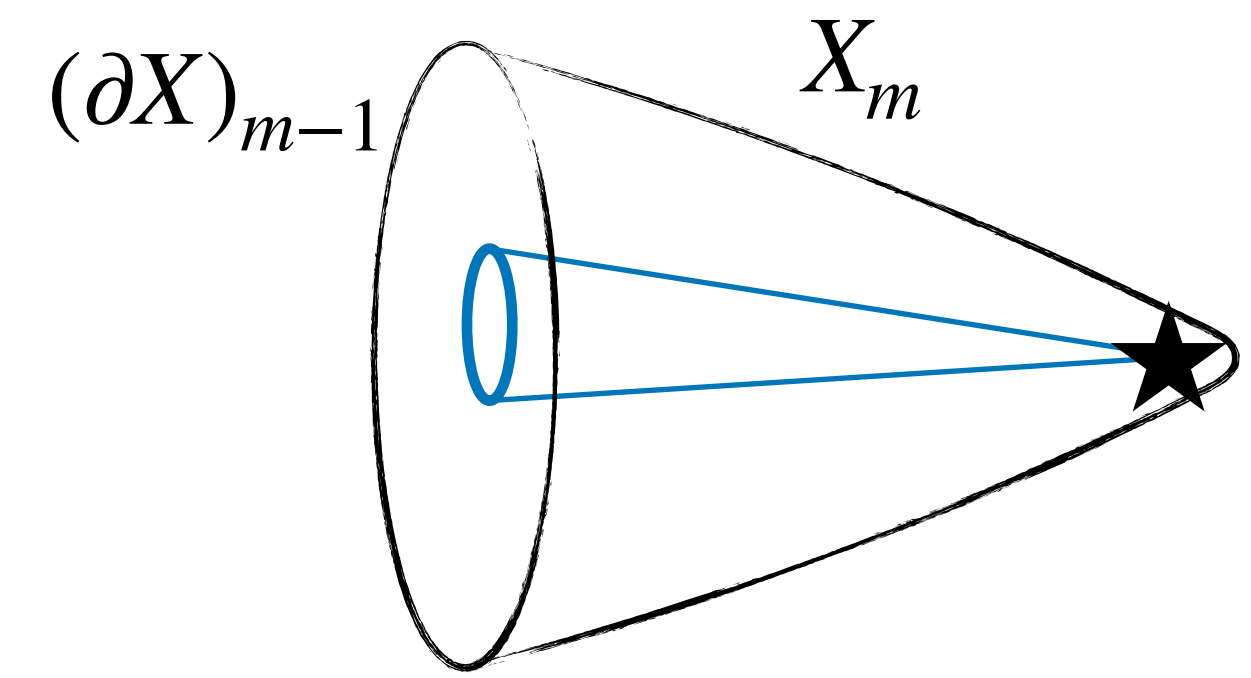
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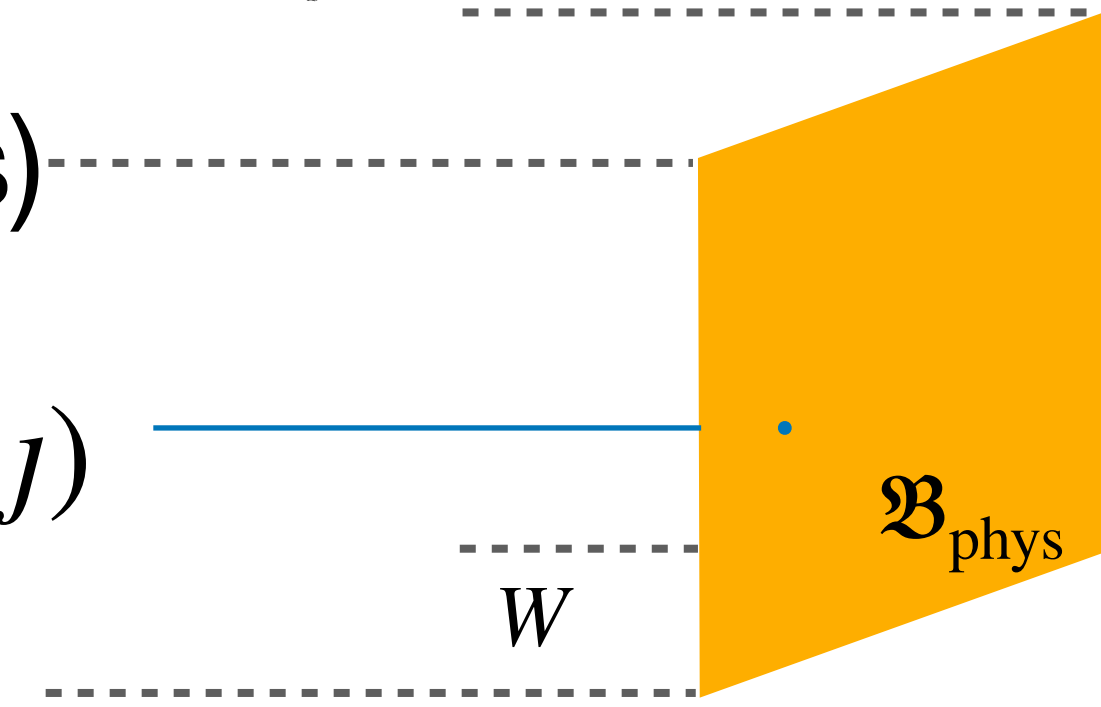


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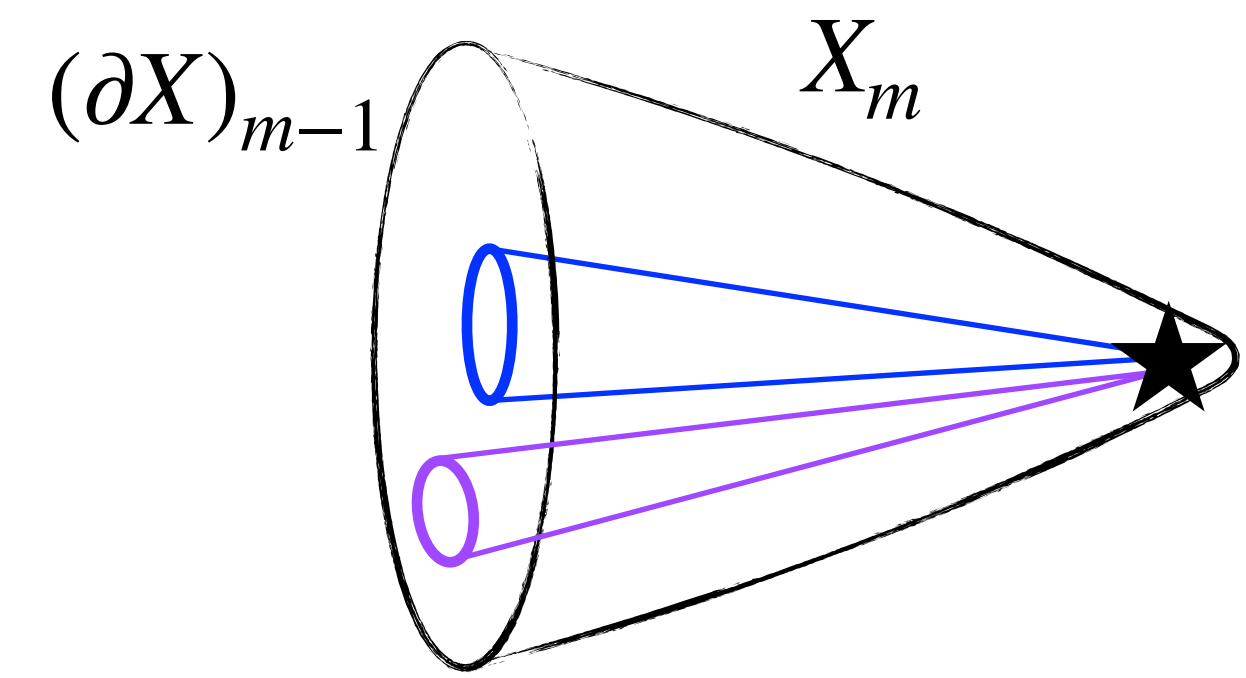
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- In cohomology:  $H^2(X) \ni \eta \xrightarrow{\partial^*} e \in H^2(\partial X)$ , with  $G_4 = B_2 e^{(2)}$  resp.  $G_4 = E_2 \eta^{(2)}$ , then  $E_2 = B_2|_{\mathfrak{B}}$  has Dirichlet b.c.



# Coupling to physical boundary... without torsion

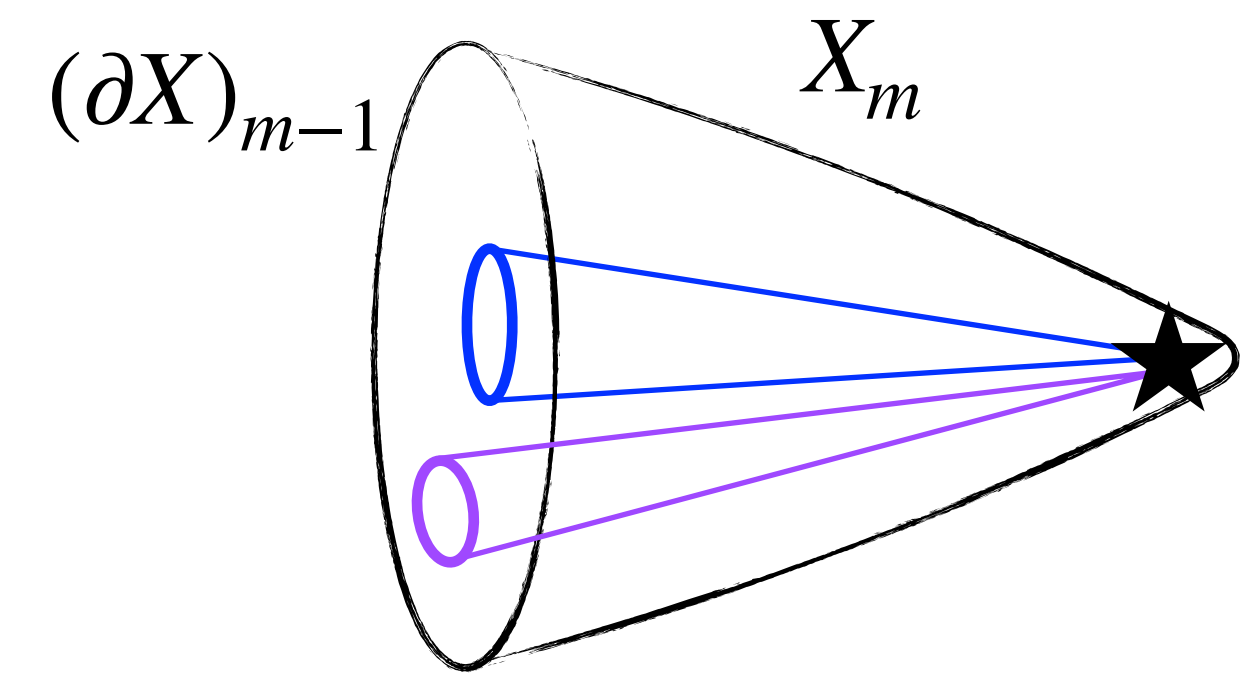


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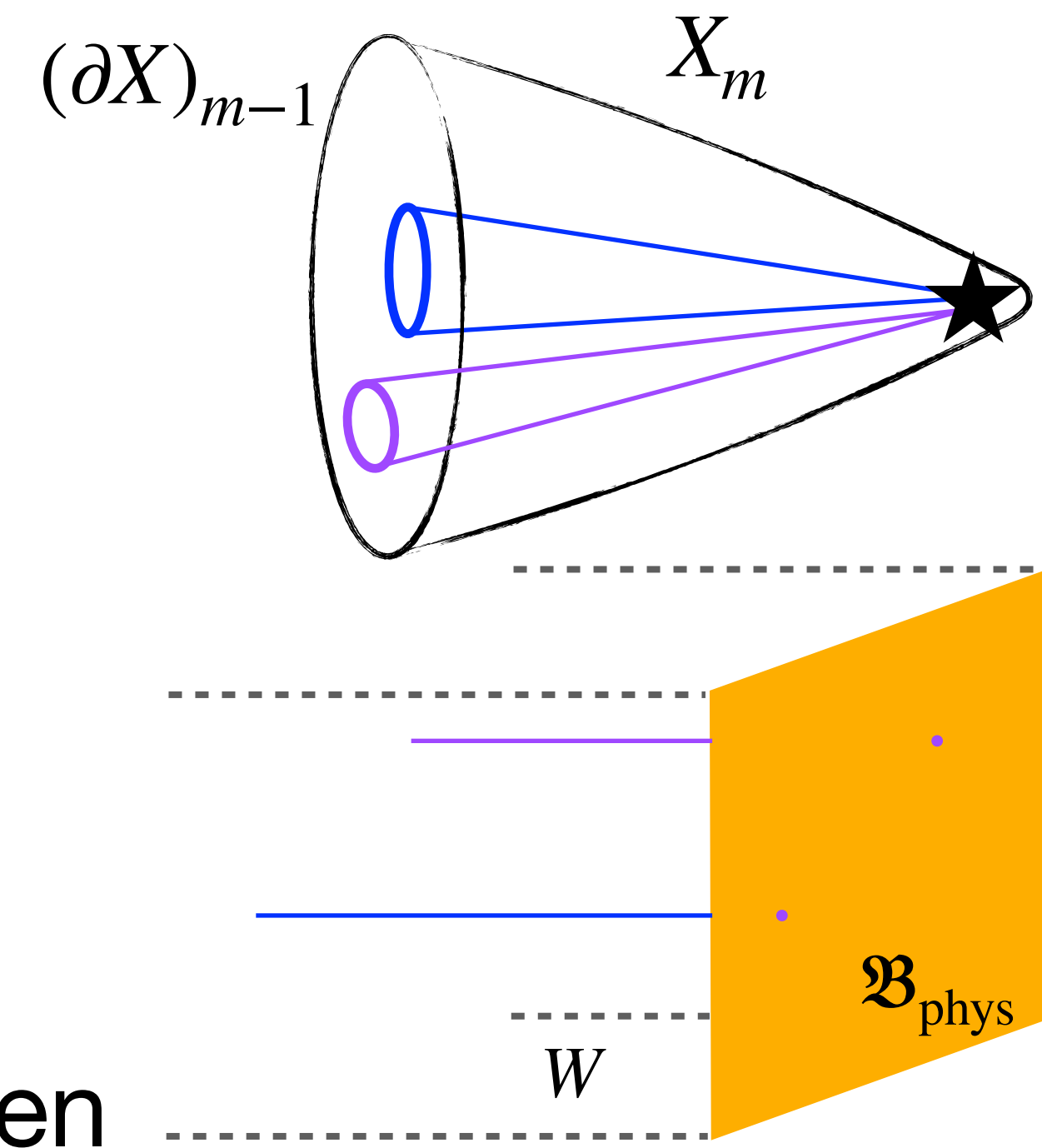
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$$1) H_p^{\text{free}}(X, \partial X) = H_{\text{free}}^{m-p}(X) = \text{Hom}(H_{m-p}^{\text{free}}(X), \mathbb{Z}), \quad 2) \text{im}(\partial) \supset H_{p-1}^{\text{tor}}(\partial X) = \mathbb{Z}_N$$

$$\text{with } \eta_2 \cdot \eta_{m-2} = \ell(r(\eta_2), r(\eta_{m-2})) \pmod{\mathbb{Z}}.$$

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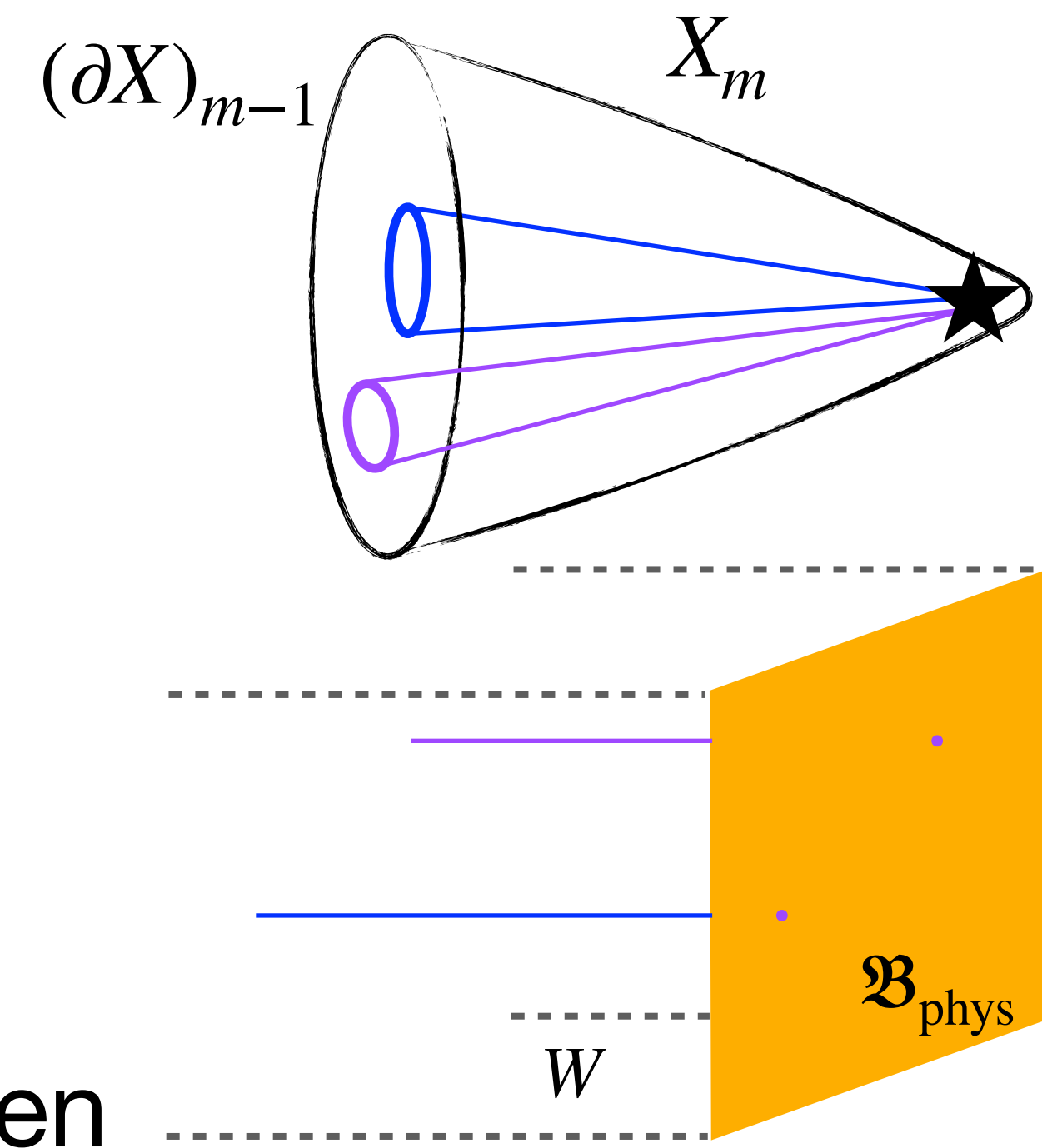


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➔ All SymTFT operators end on charged defects on  $\mathfrak{B}$ .

- Can show in cohomology:  $B_p|_{\mathfrak{B}} = \frac{1}{N} f_p = \frac{1}{N} da_{p-1}$  as eom;  
so defects are (non-top.) Wilson/'t Hooft operators of SYM sector.

# Coupling to physical boundary... with torsion

[see also Cvetič/Heckman/Hübner/Torres '23, Gould/LL/Sabag '23,  
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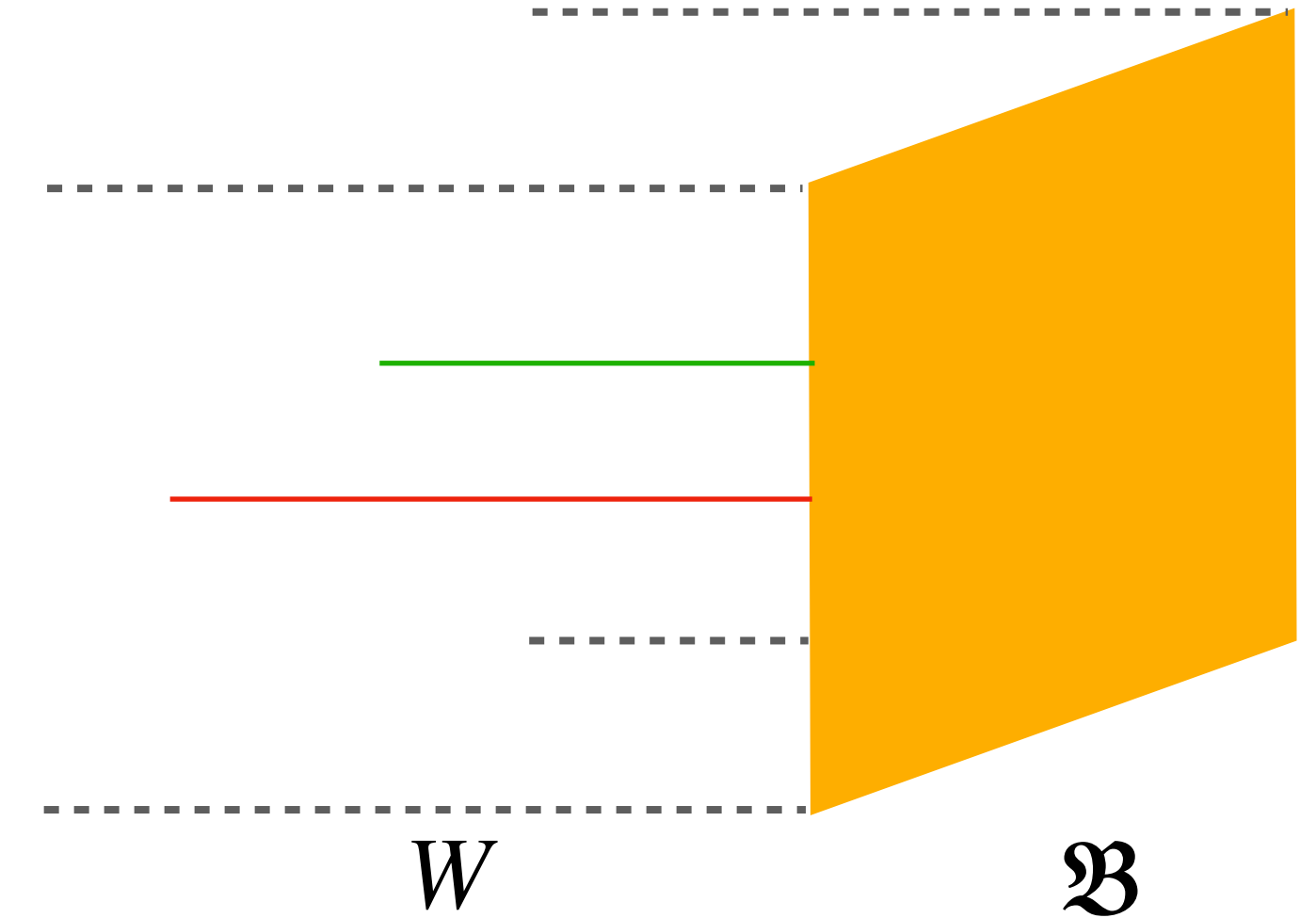
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- With torsion:  $H_p^{\text{tor}}(X, \partial X) = H_{\text{tor}}^{m-p}(X) = \text{Hom}(H_{m-p-1}^{\text{tor}}(X), \mathbb{Q}/\mathbb{Z}) \cong H_{m-p-1}^{\text{tor}}(X)$ .

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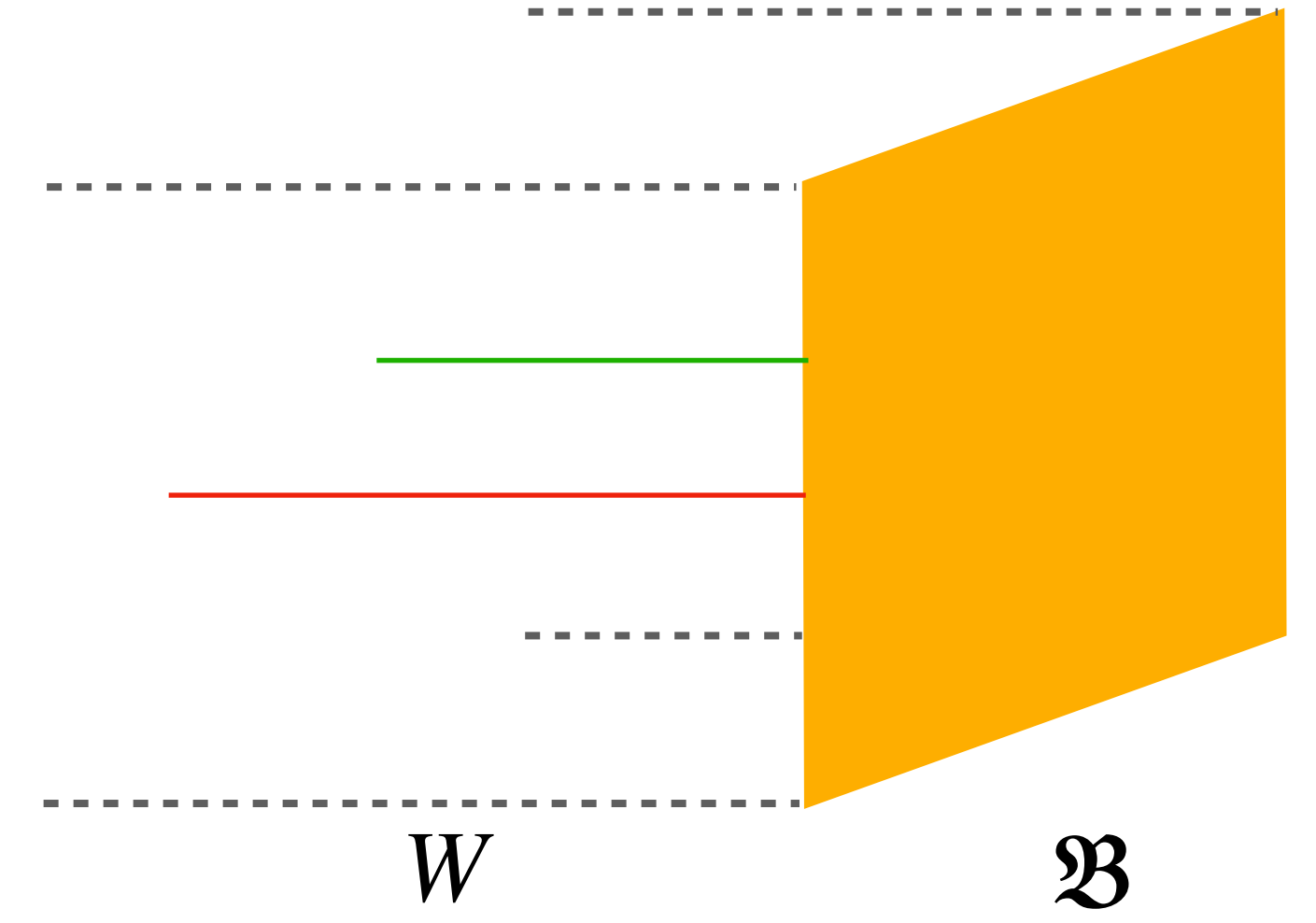
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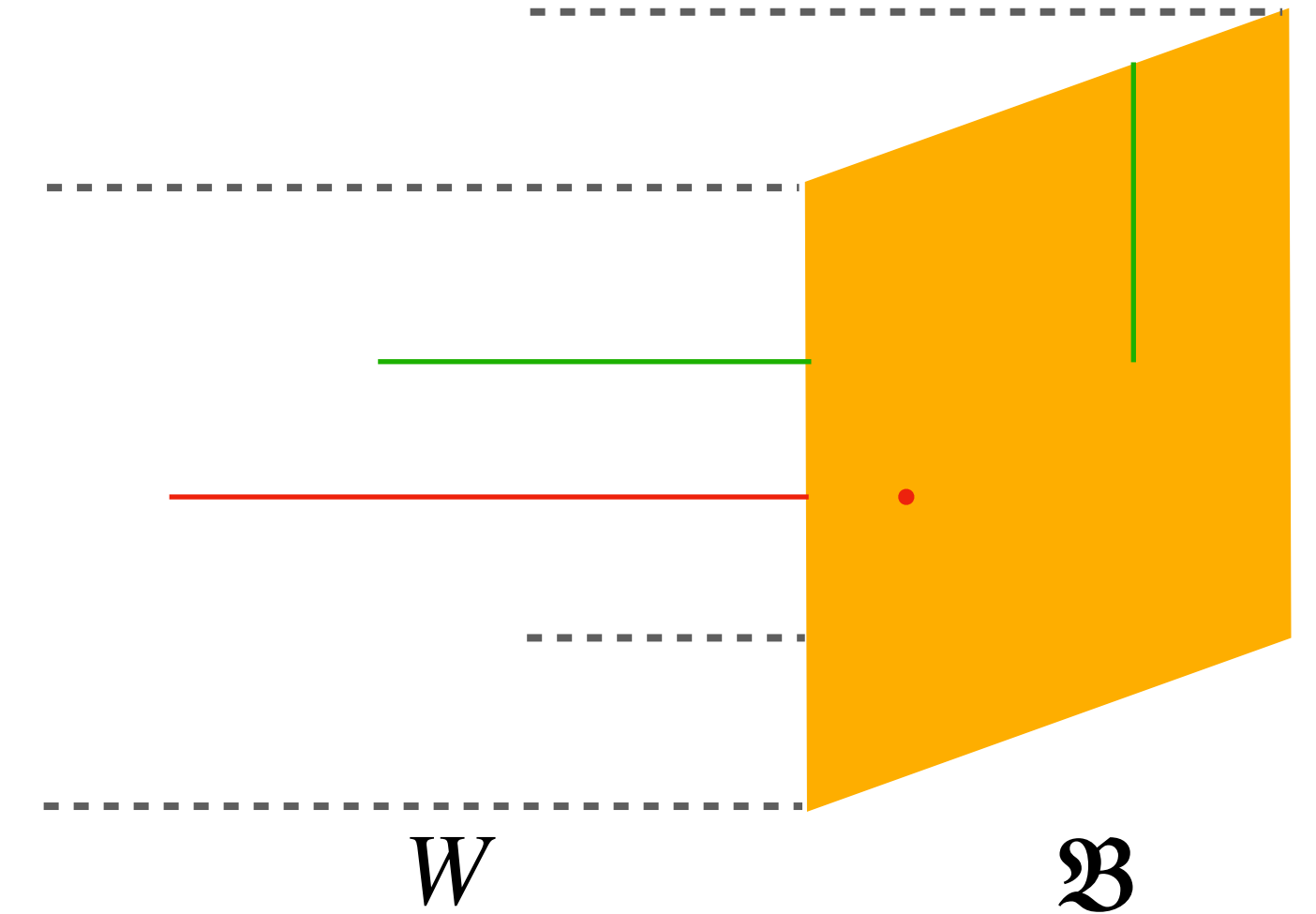
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- E.g.,  $0 \neq \tau_1 \in H_1^{\text{tor}}(X)$ ,  $\Sigma_{m-2} \in H_{m-2}^{\text{tor}}(X, \partial X)$  s.t.  $\partial(\Sigma_{m-2}) = s_{m-3}$ ,  $l(t_1) = \tau_1$  with  $(\tau_1, \Sigma_{m-2}) = \ell(t_1, s_{m-3})$ .

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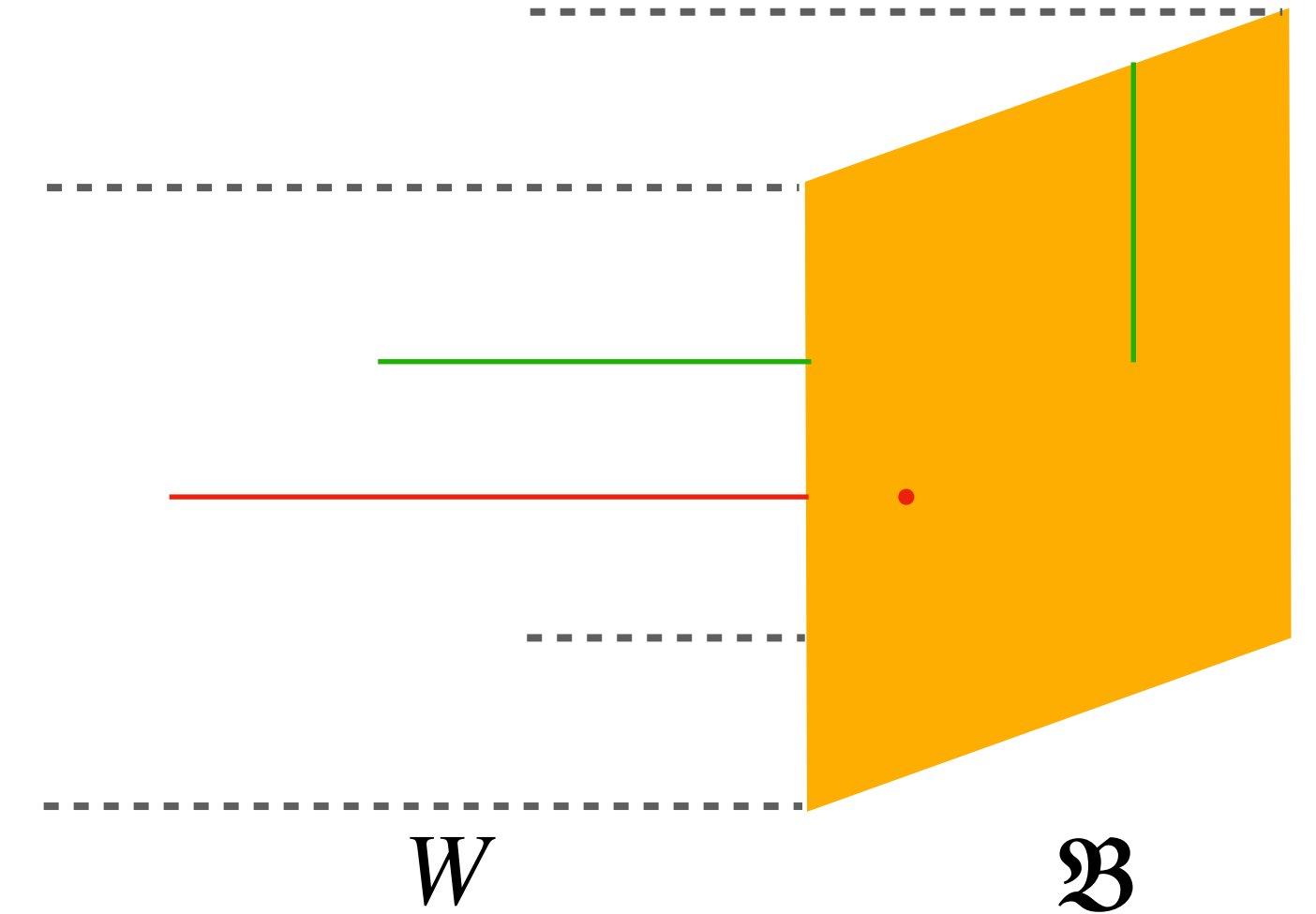
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  - E.g.,  $0 \neq \tau_1 \in H_1^{\text{tor}}(X)$ ,  $\Sigma_{m-2} \in H_{m-2}^{\text{tor}}(X, \partial X)$  s.t.  $\partial(\Sigma_{m-2}) = s_{m-3}$ ,  $l(t_1) = \tau_1$  with  $(\tau_1, \Sigma_{m-2}) = \ell(t_1, s_{m-3})$ .
- ➔  $M5[s_{m-3}]$  ends on  $M5[\Sigma_{m-2}]$ , but  $M2[t_1]$  continues as  $M2[\tau_1]$ !

# Coupling to physical boundary... with torsion

[see also Cvetič/Heckman/Hübner/Torres '23, Gould/LL/Sabag '23, Cvetič/Dierigl/LL/Torres/Zhang '24]



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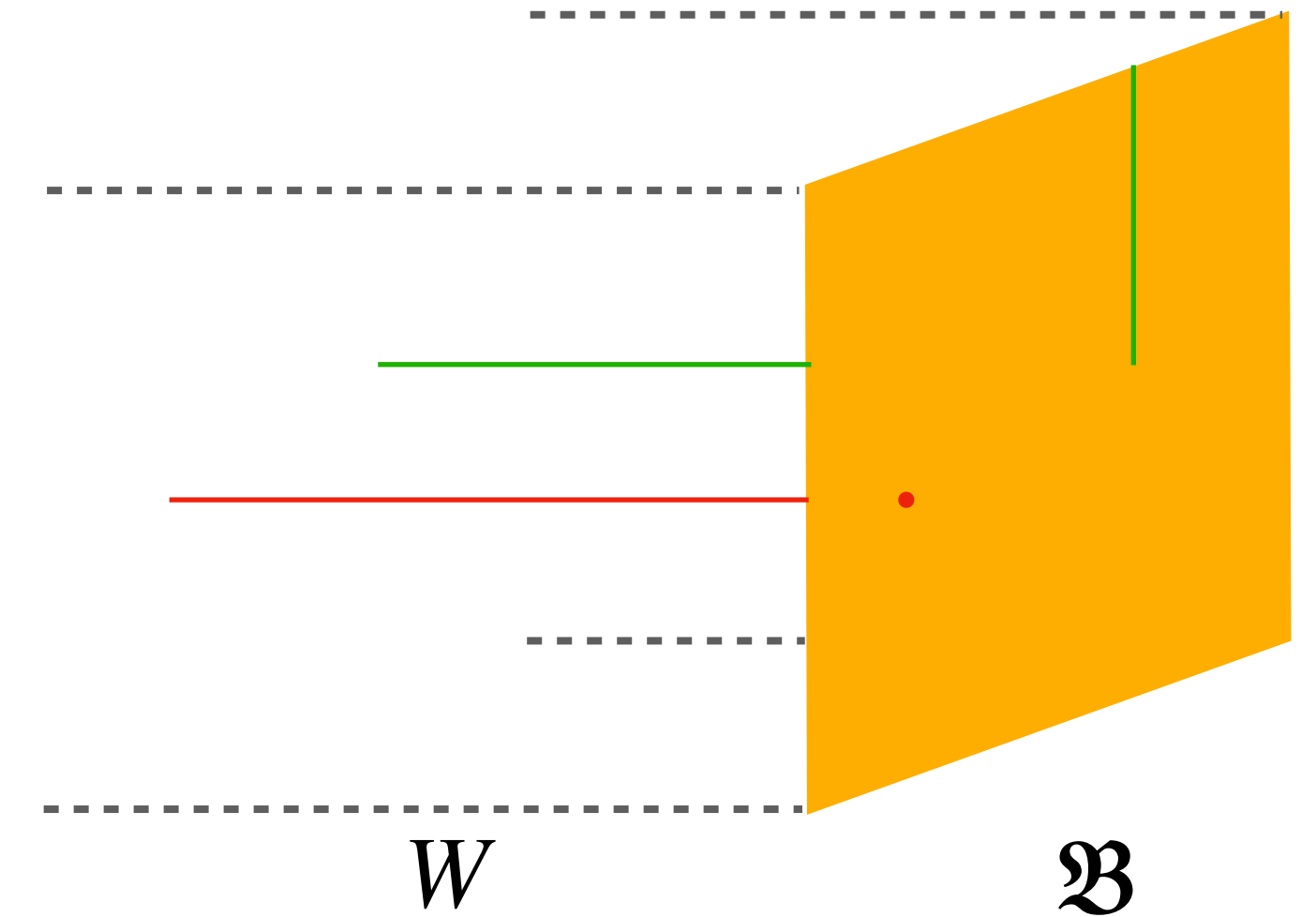
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- From cohomology:  $G_7 = c_{8-m} \omega^{(m-1)}$  with  $\omega^{(m-1)} \in H_{\text{tor}}^{m-1}(X, \partial X) = H_{m-2}^{\text{tor}}(X, \partial X) \rightsquigarrow B_{9-m}|_{\mathfrak{g}} = \delta c_{8-m}(\text{D})$ , but  $B_2|_{\mathfrak{g}}$  free (N) [for  $S_W \supset \frac{1}{N} \int_W B_2 \delta B_{9-m}$ ]

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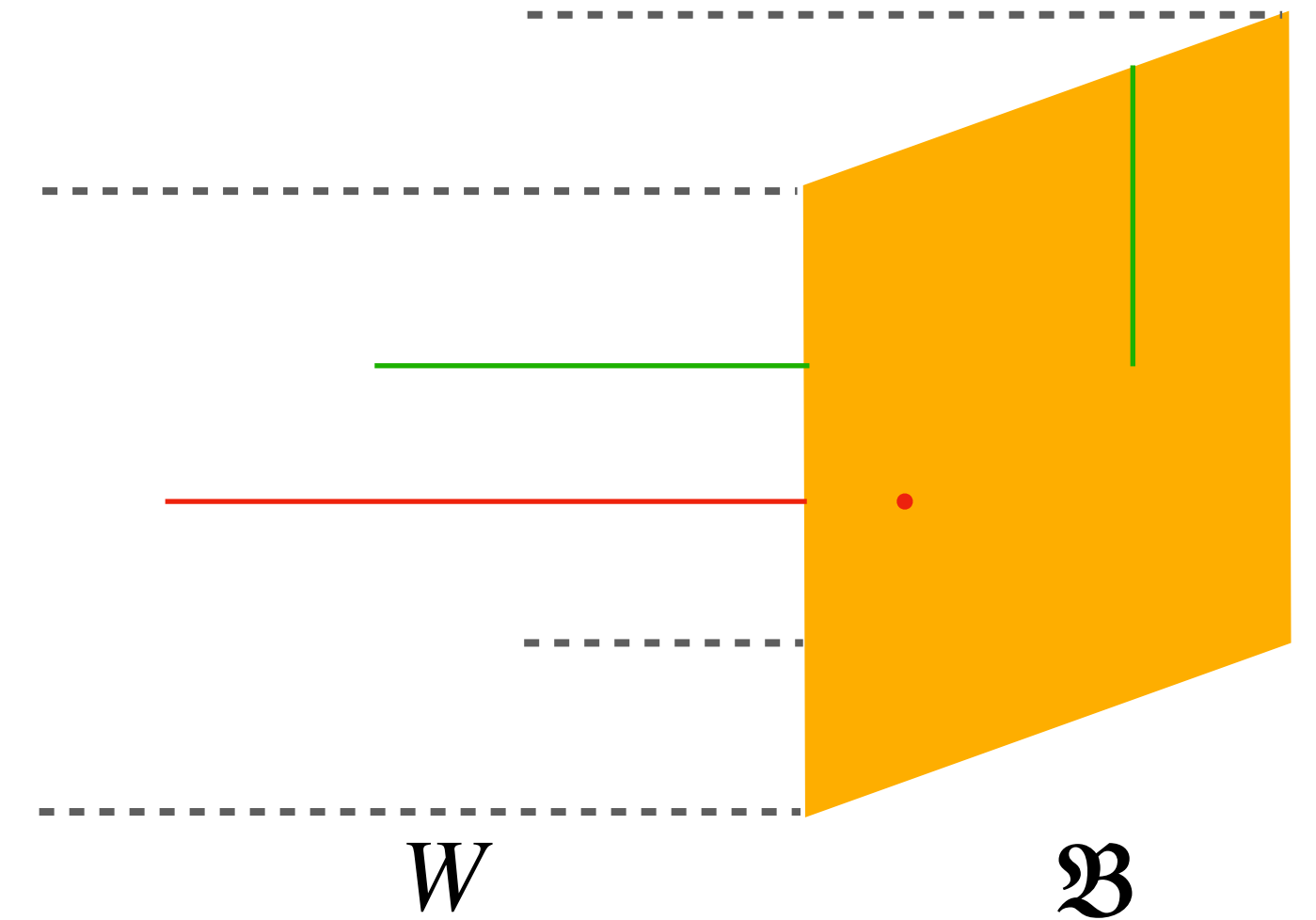


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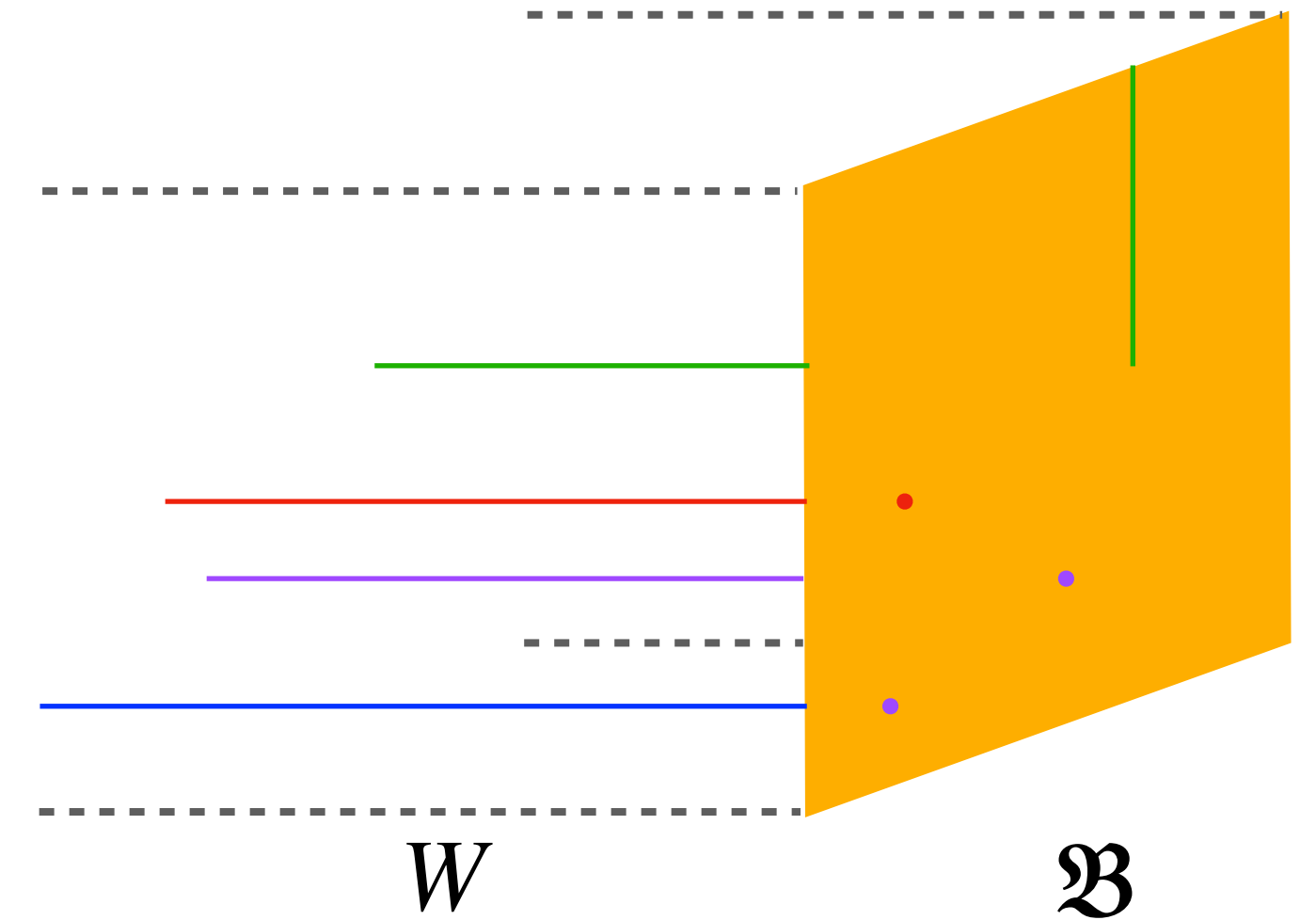


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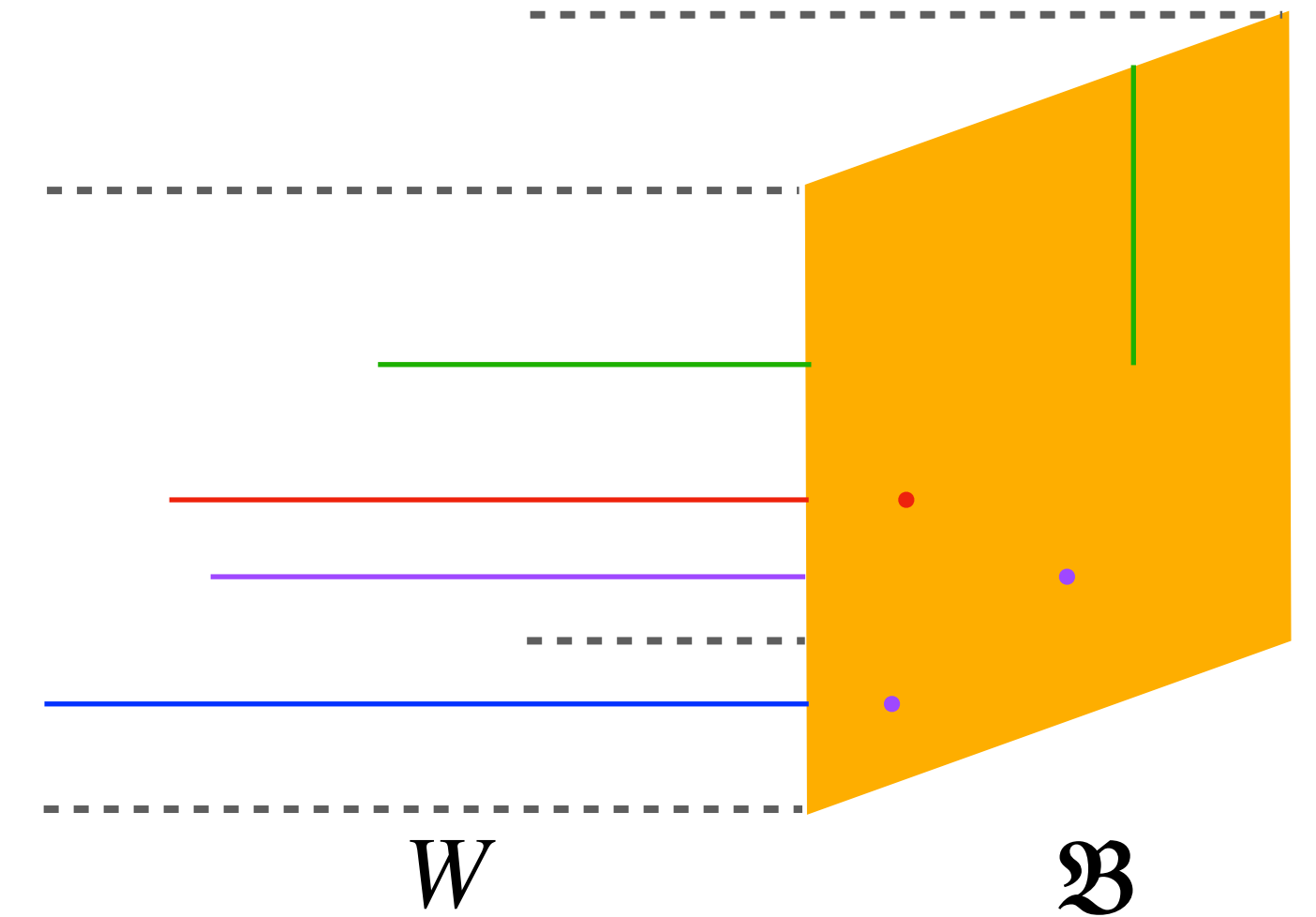
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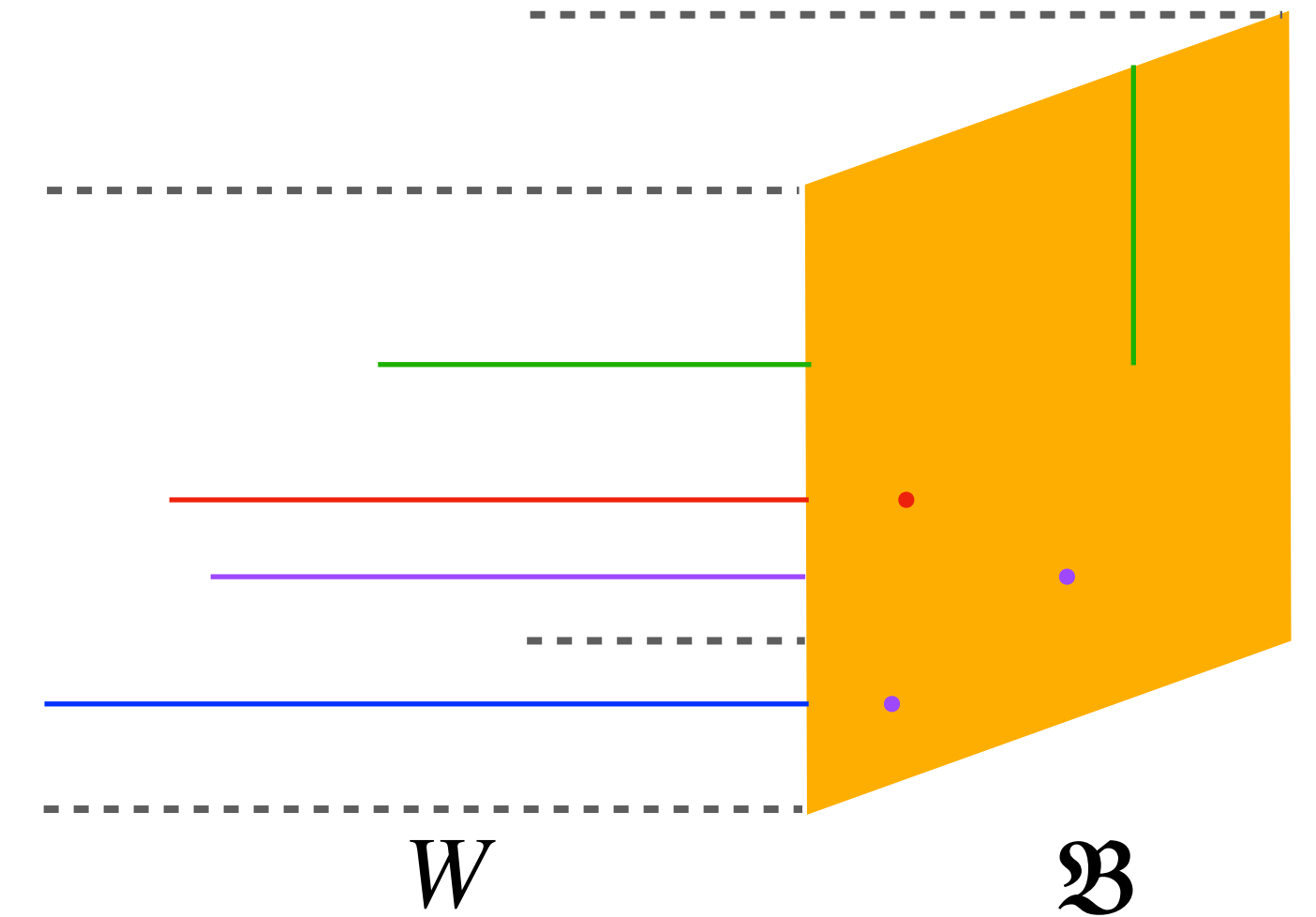


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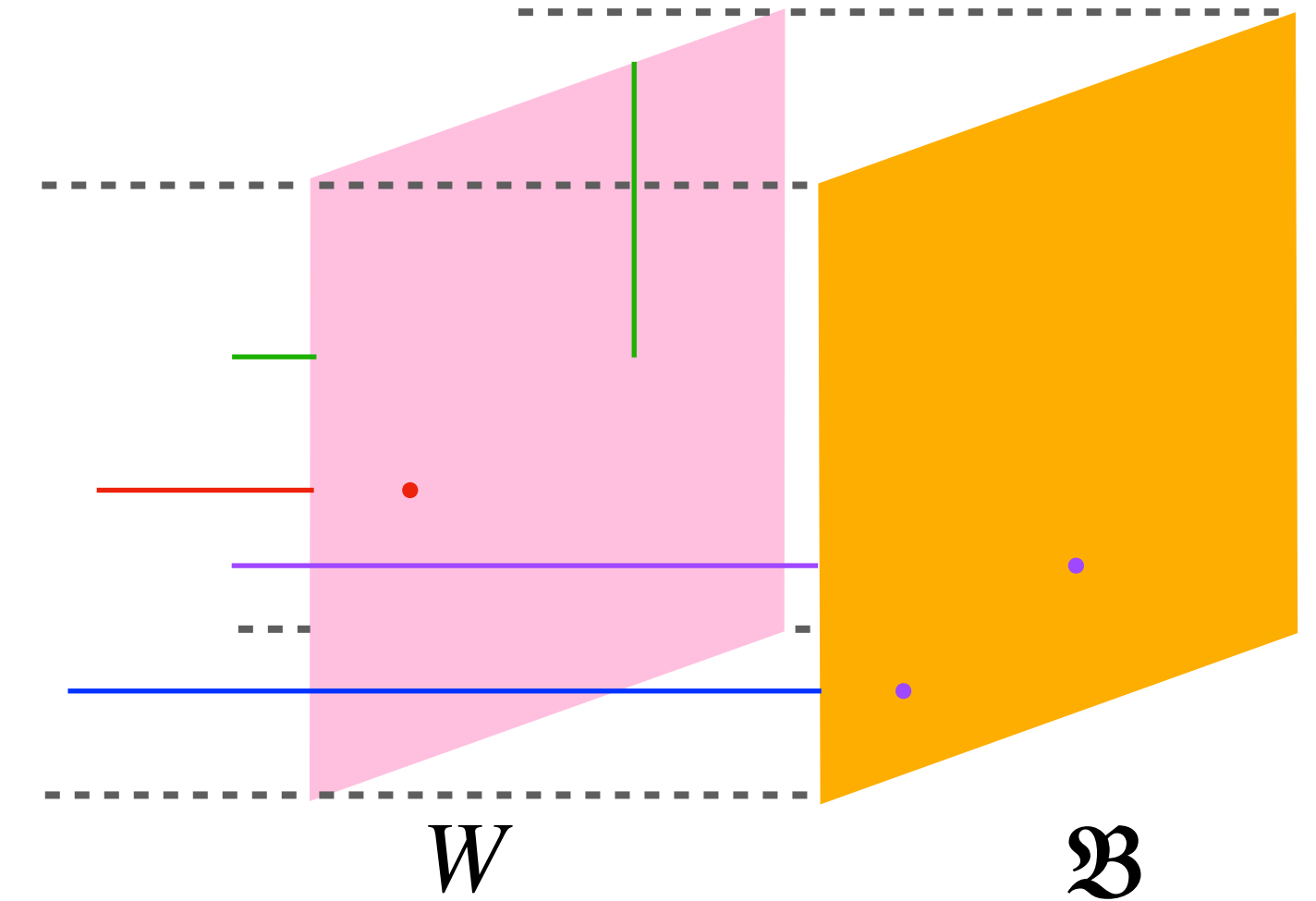


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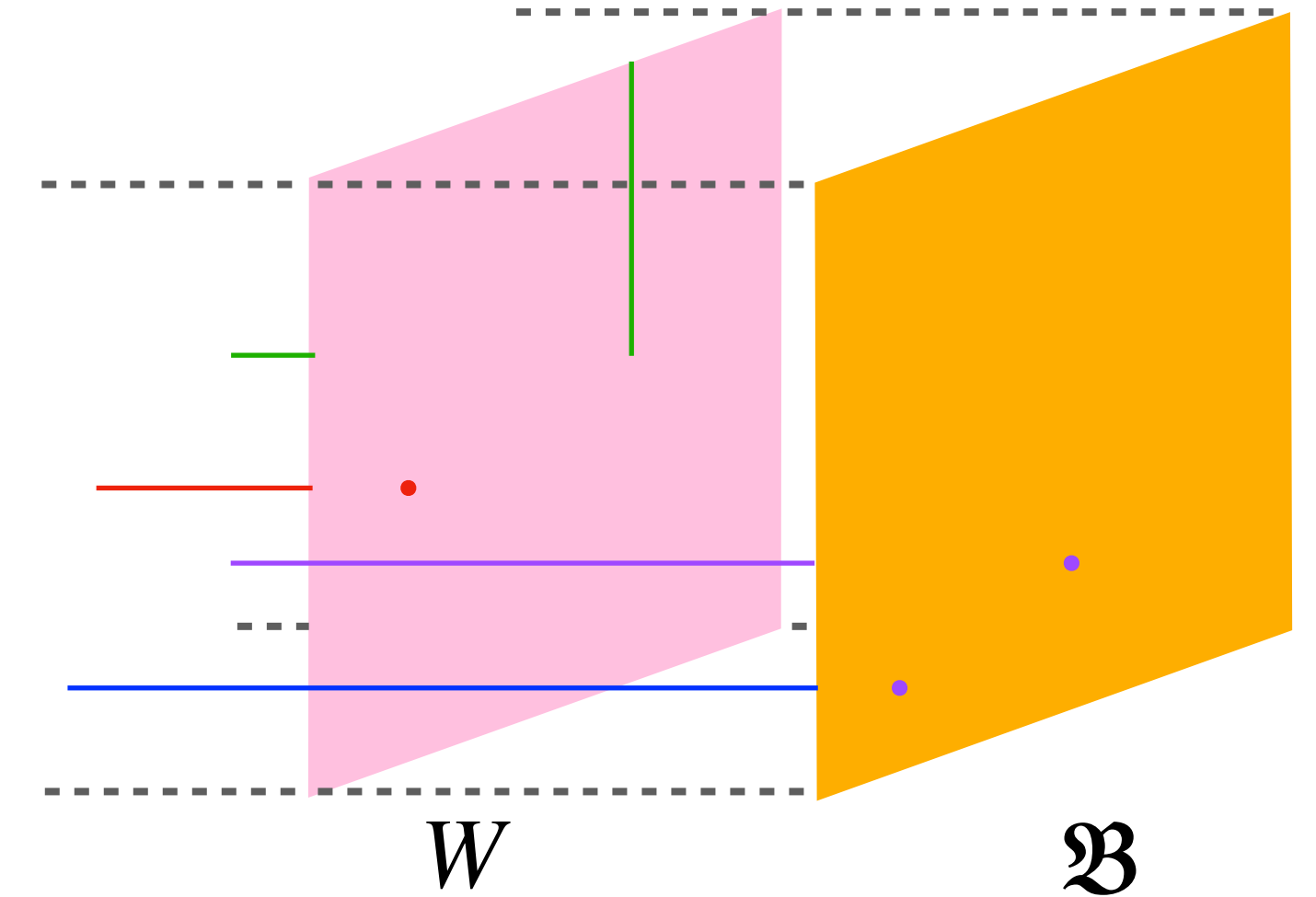


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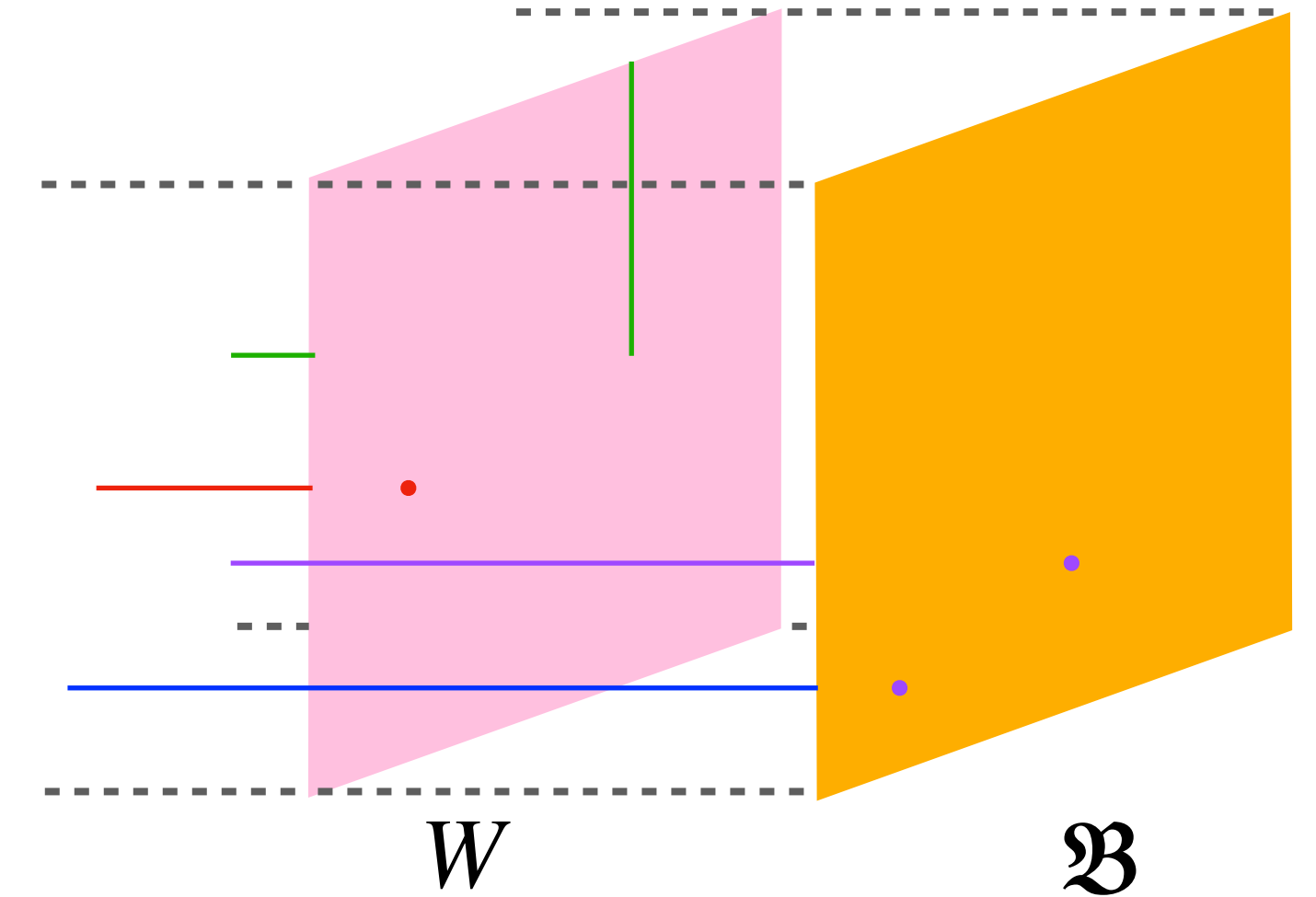


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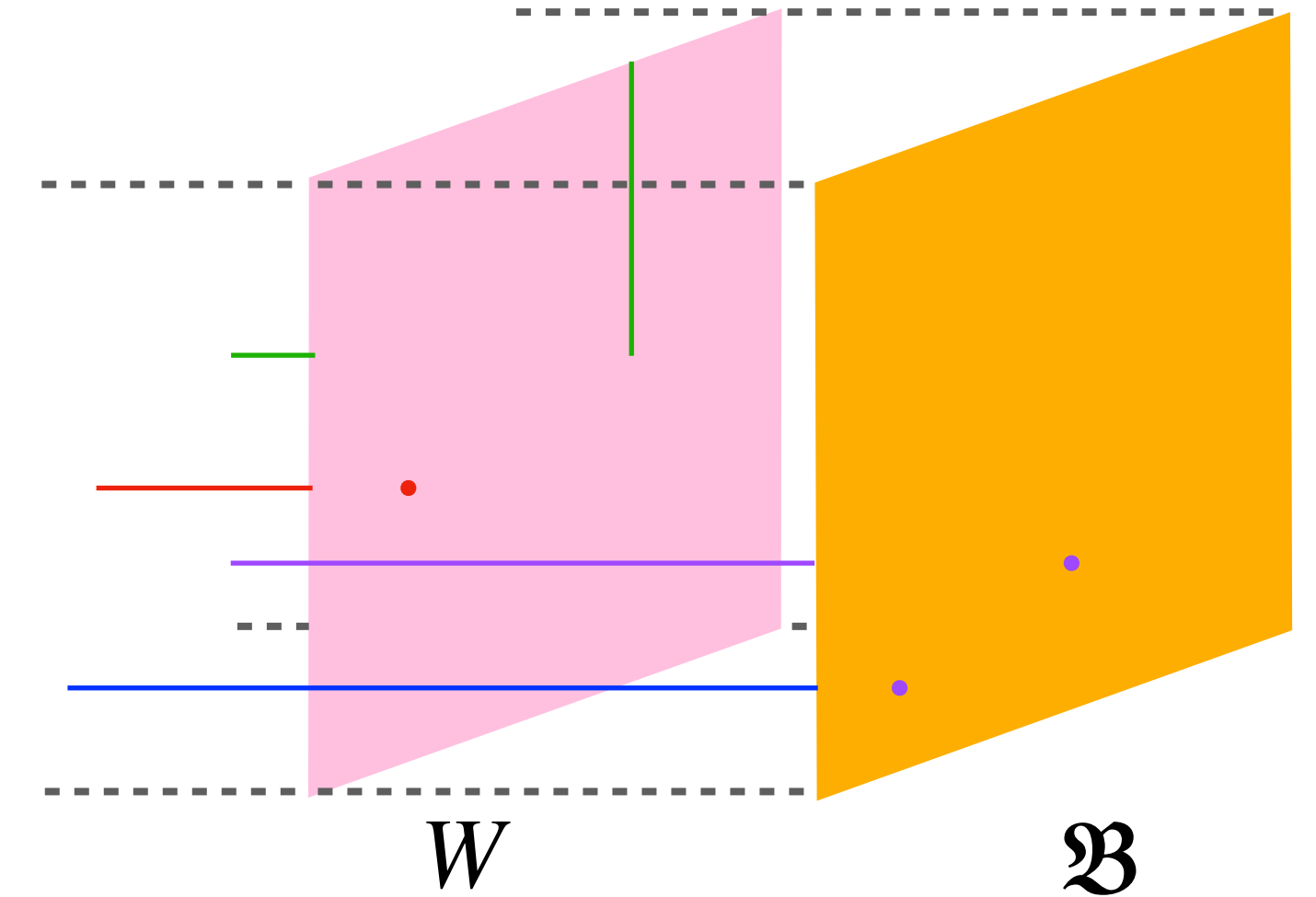


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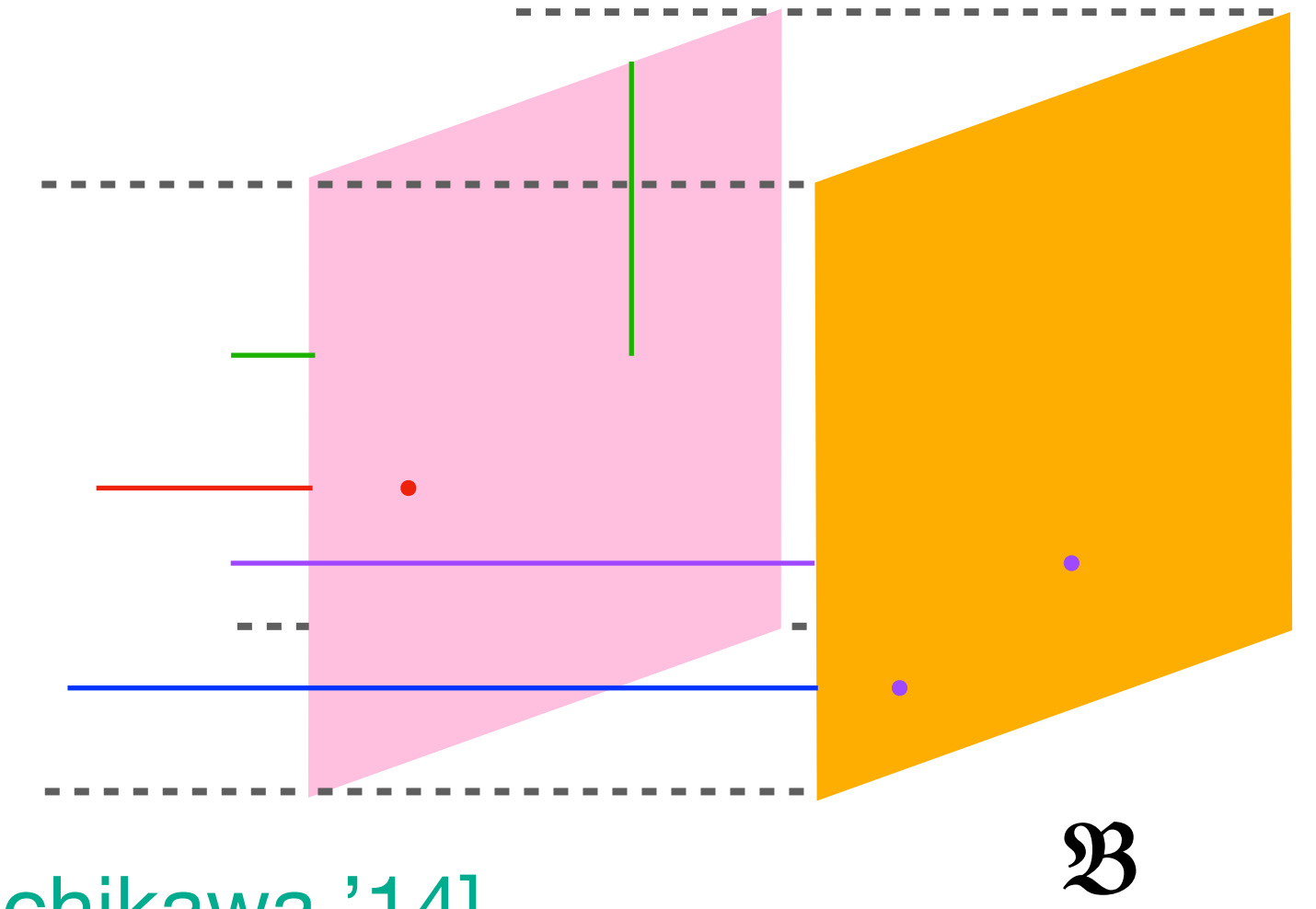
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- Note: geometry  $\rightsquigarrow$  interface on top of  $\mathfrak{B}$

# Plan of the talk

1. Topological SymTFT-interface in M-theory geometric engineering
2. Interfaces in SymTFT of “frozen”  $d \geq 6$   $\mathfrak{sp}(N)$  SYM
3. Summary & Outlook

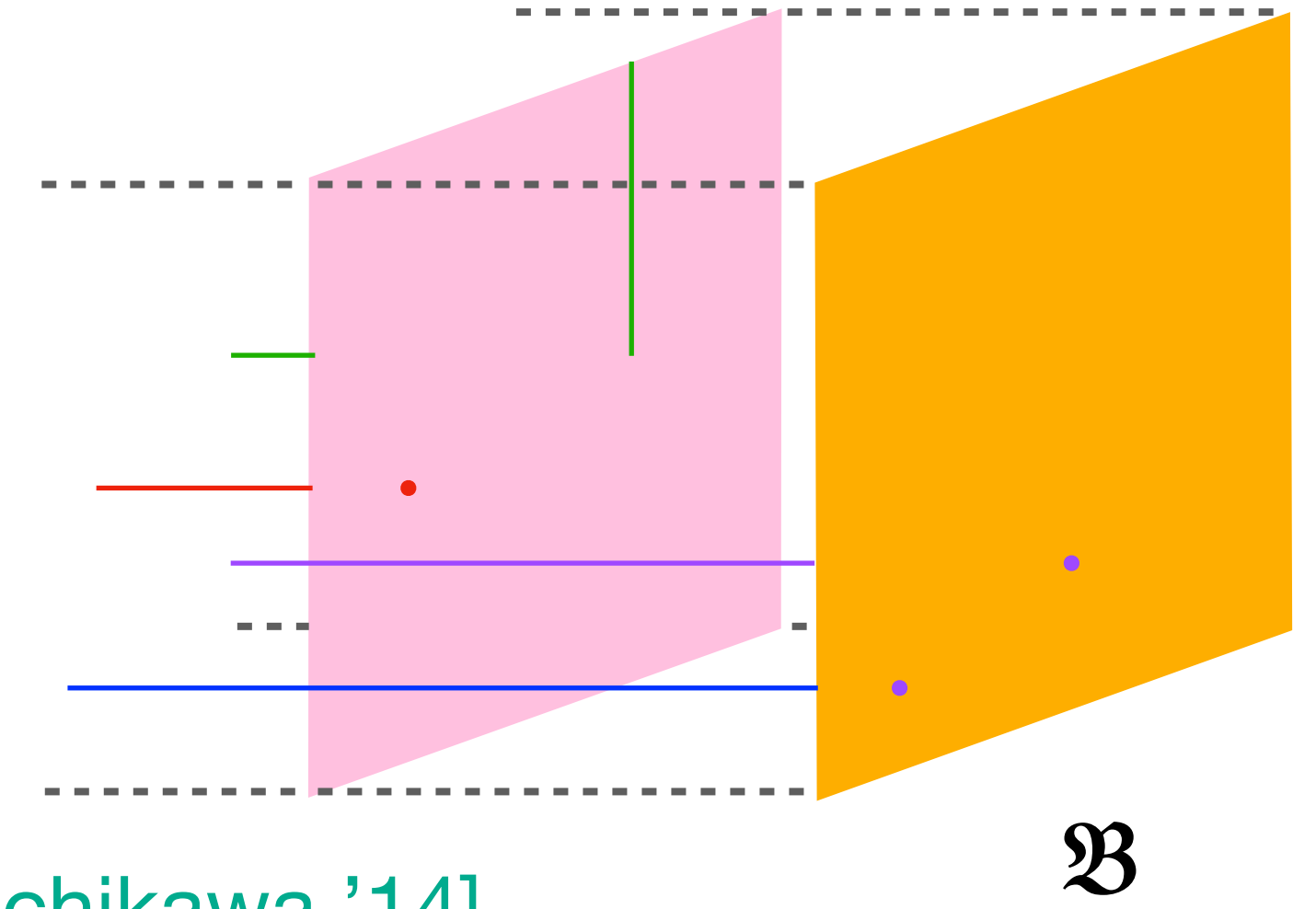
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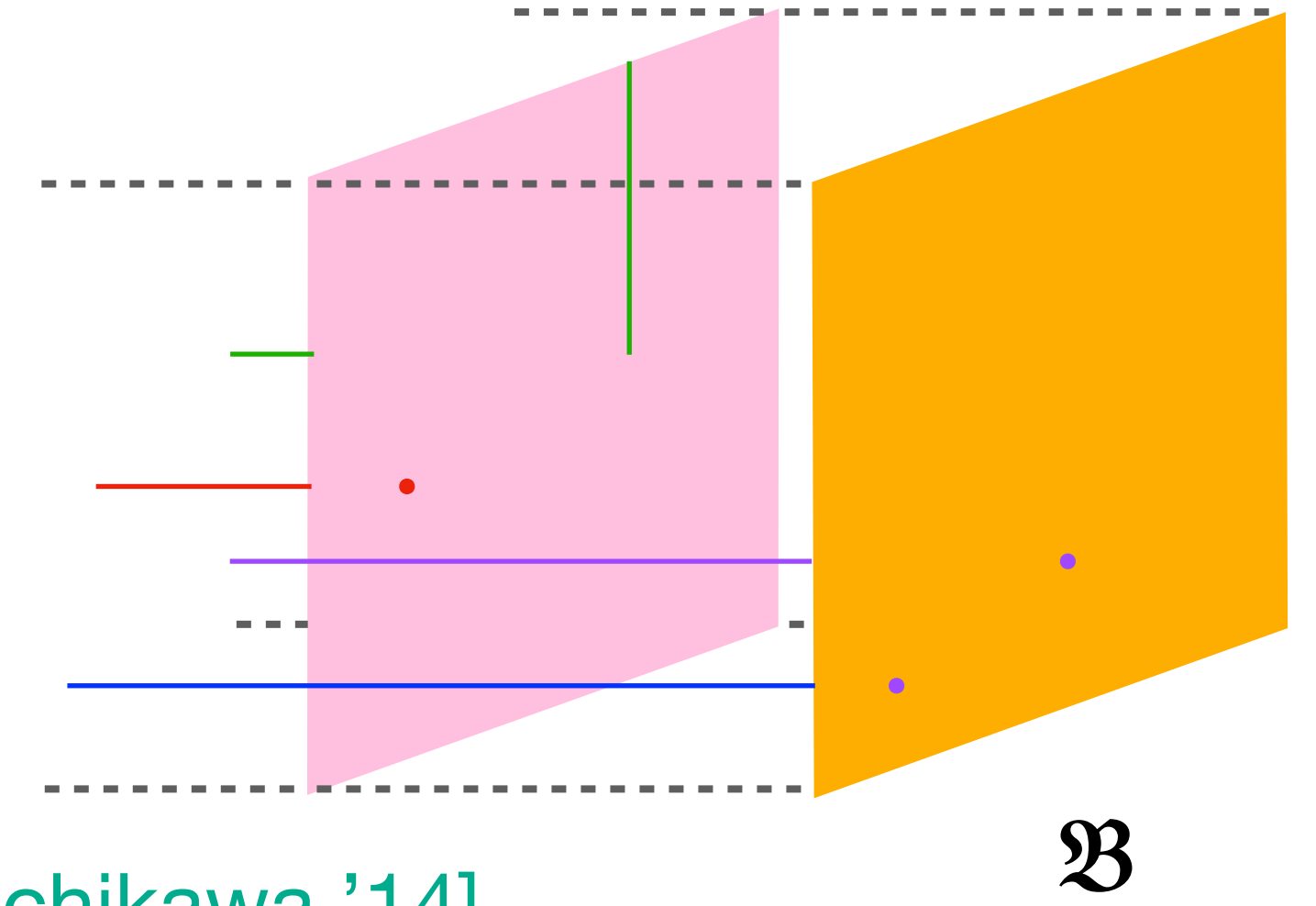


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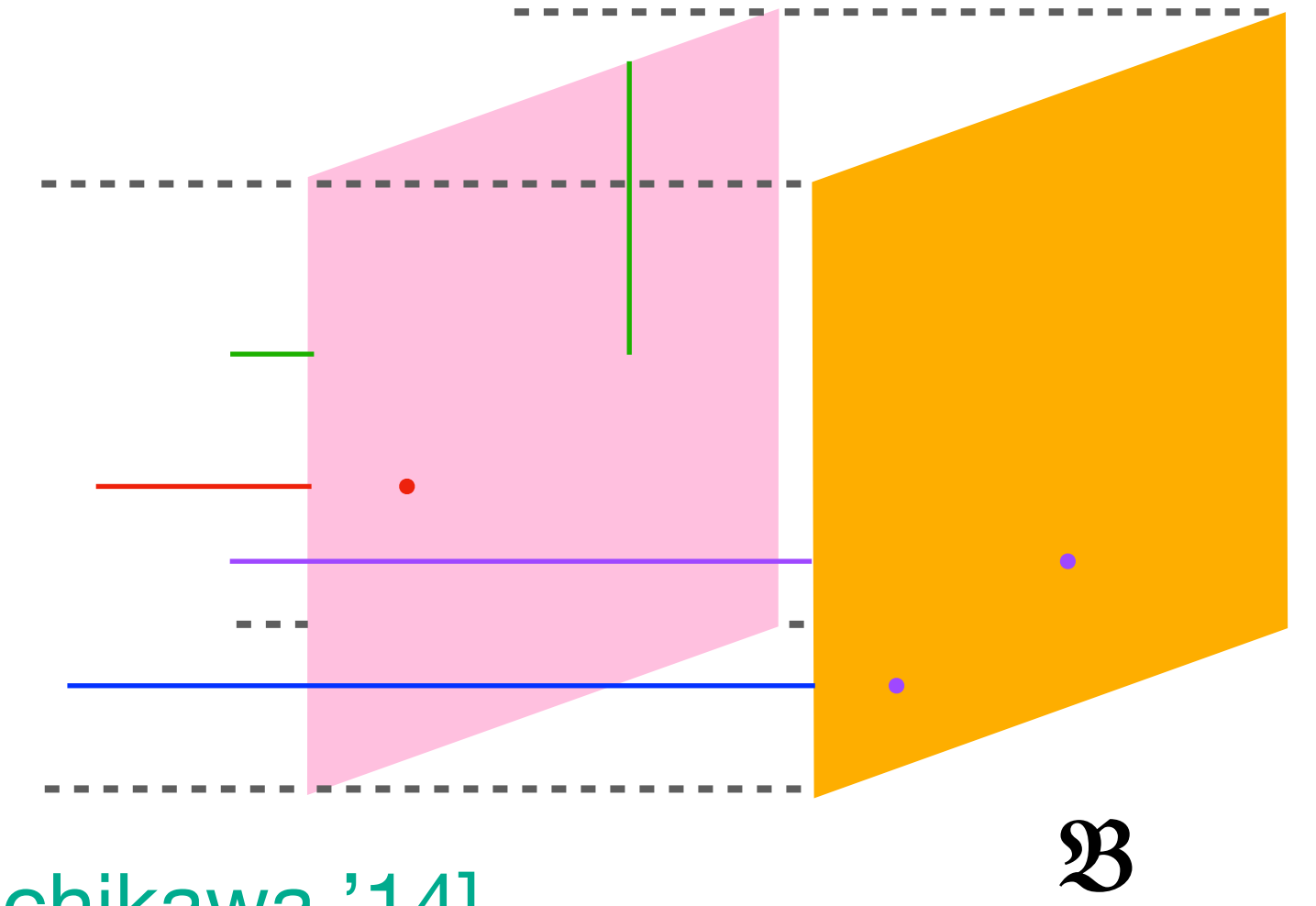
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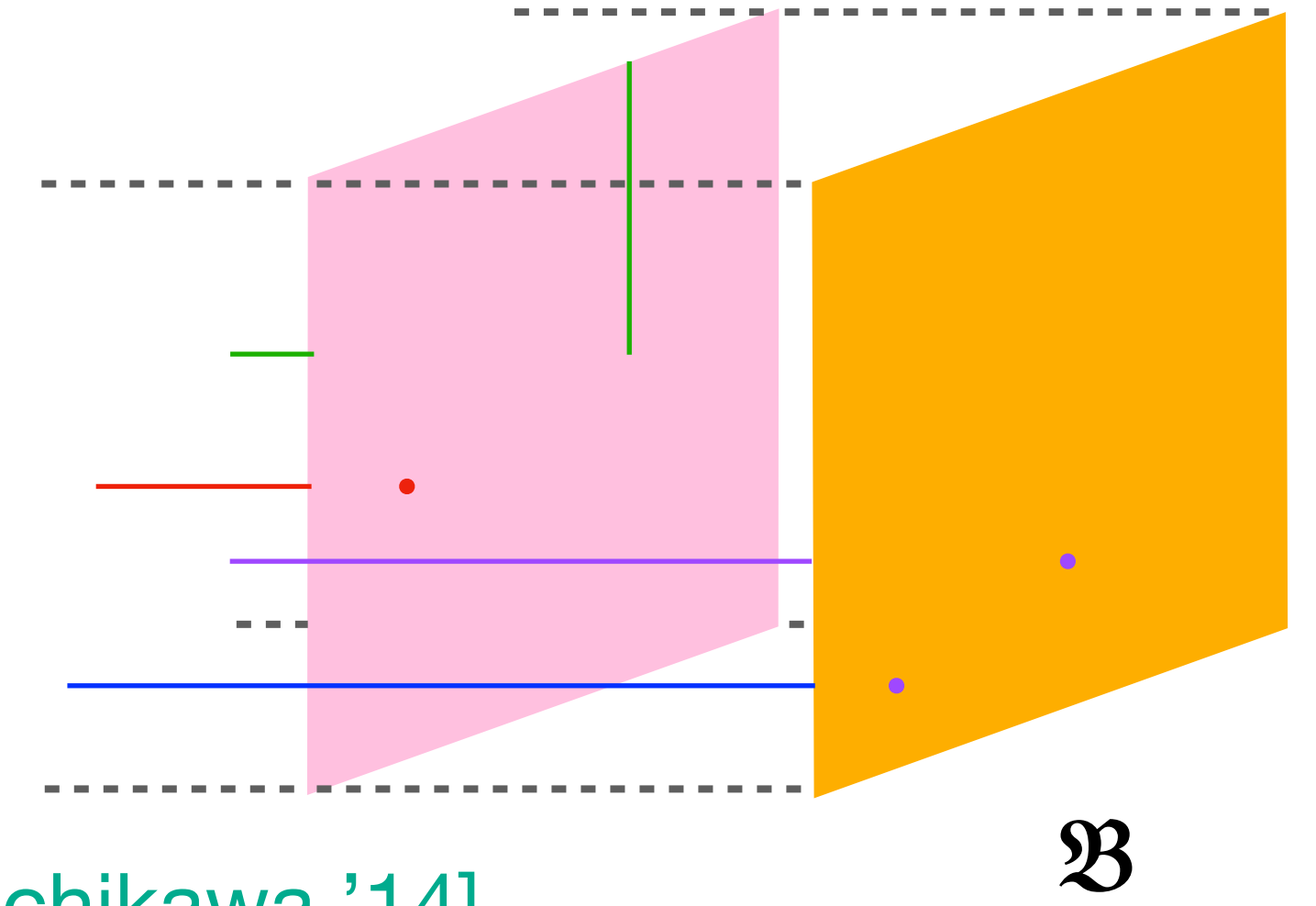
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- Stack  $\mathcal{F}$  &  $\mathcal{B} \rightsquigarrow$  charged defects for  $\mathfrak{sp}$ -center  $(\mathcal{W}^{(1)} / \mathcal{H}^{(d-3)})$  and  $\mathbb{Z}_2^{(d-3)}$  only, agrees with 7d/8d results [Cvetič/Dierigl/LL/Zhang '22, Cvetič/Dierigl/LL/Torres/Zhang '24].

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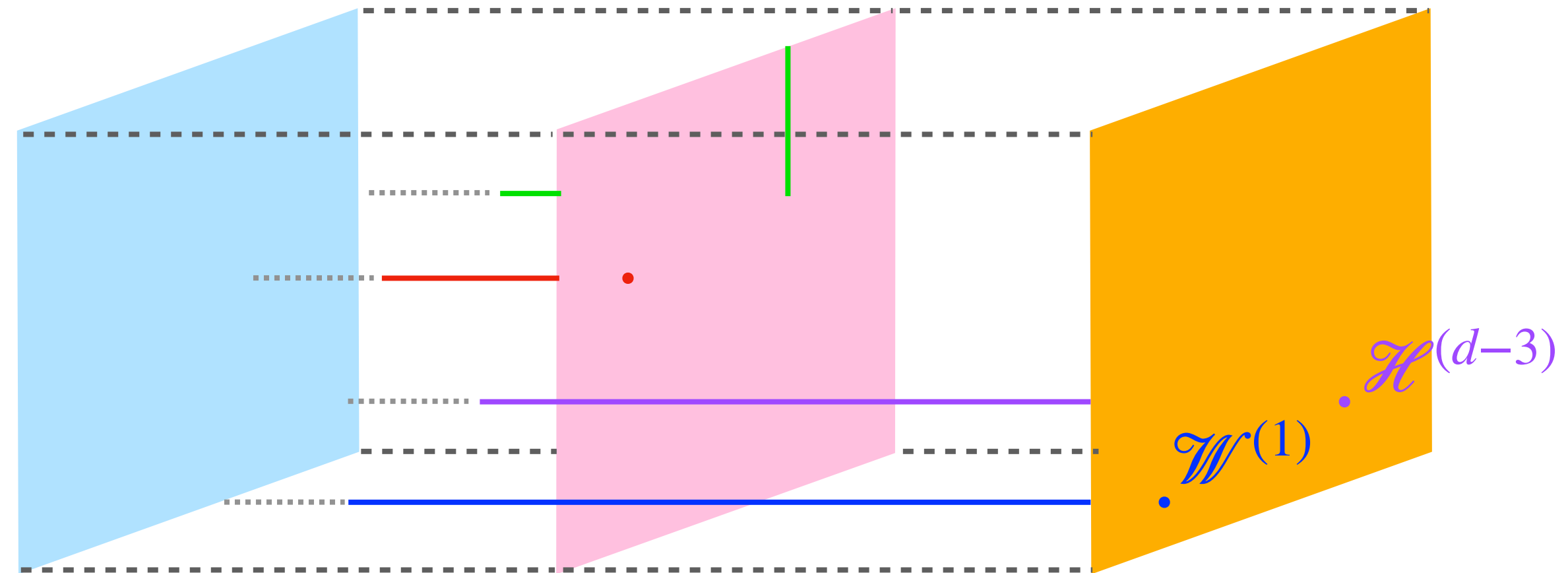
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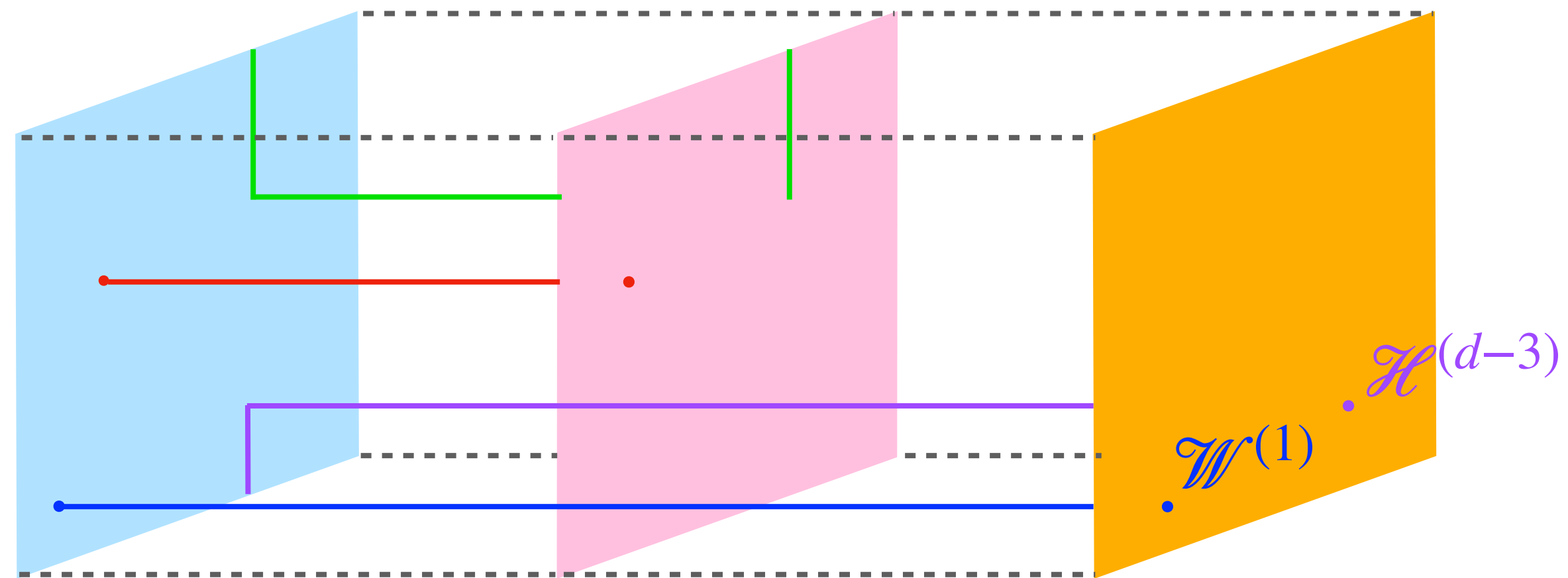
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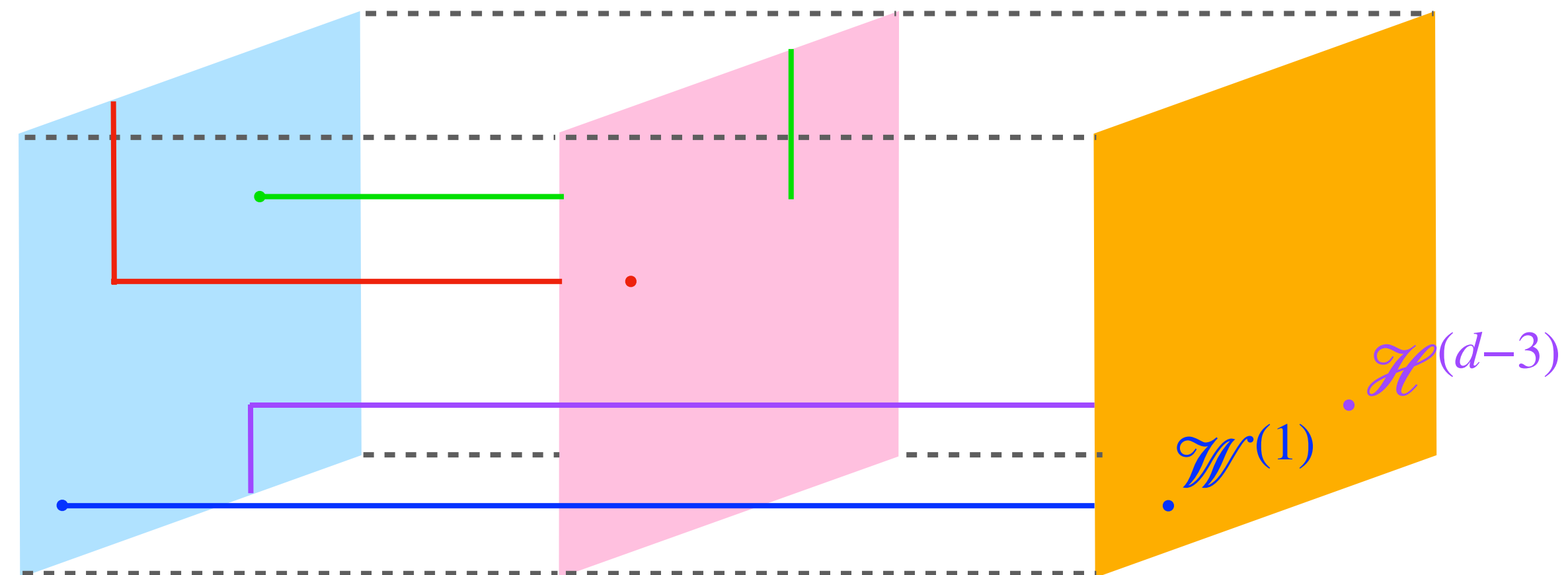


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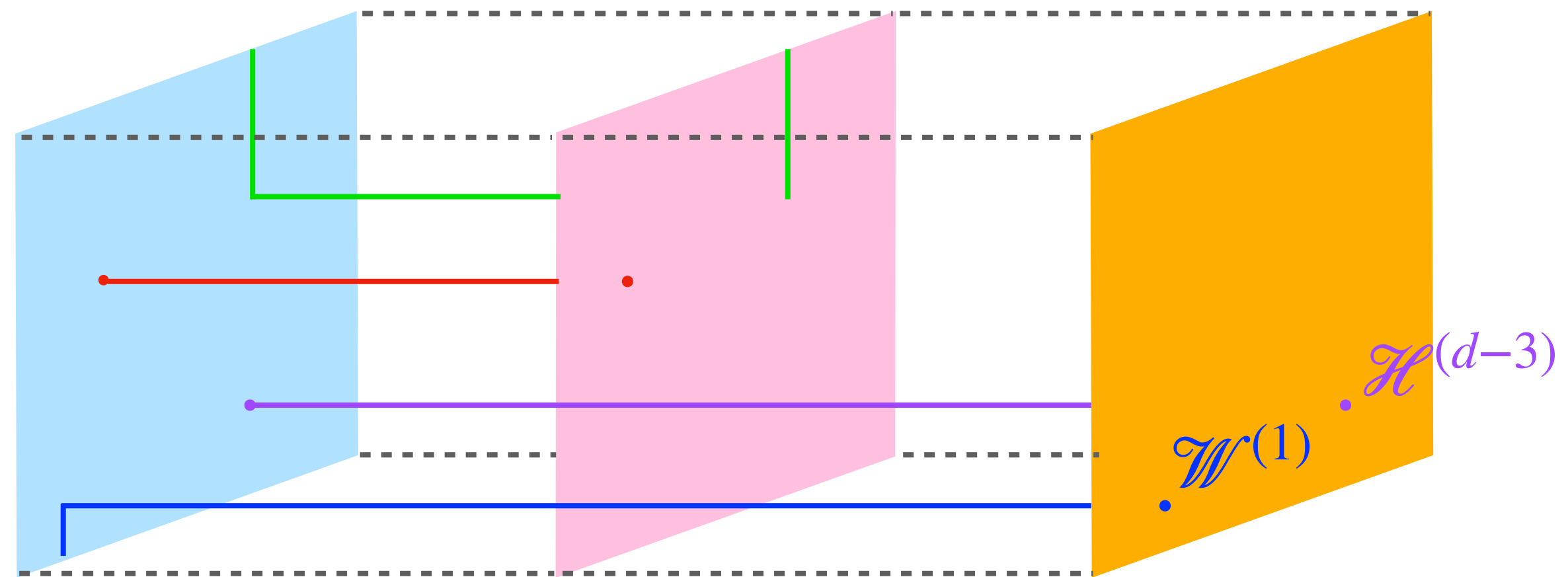
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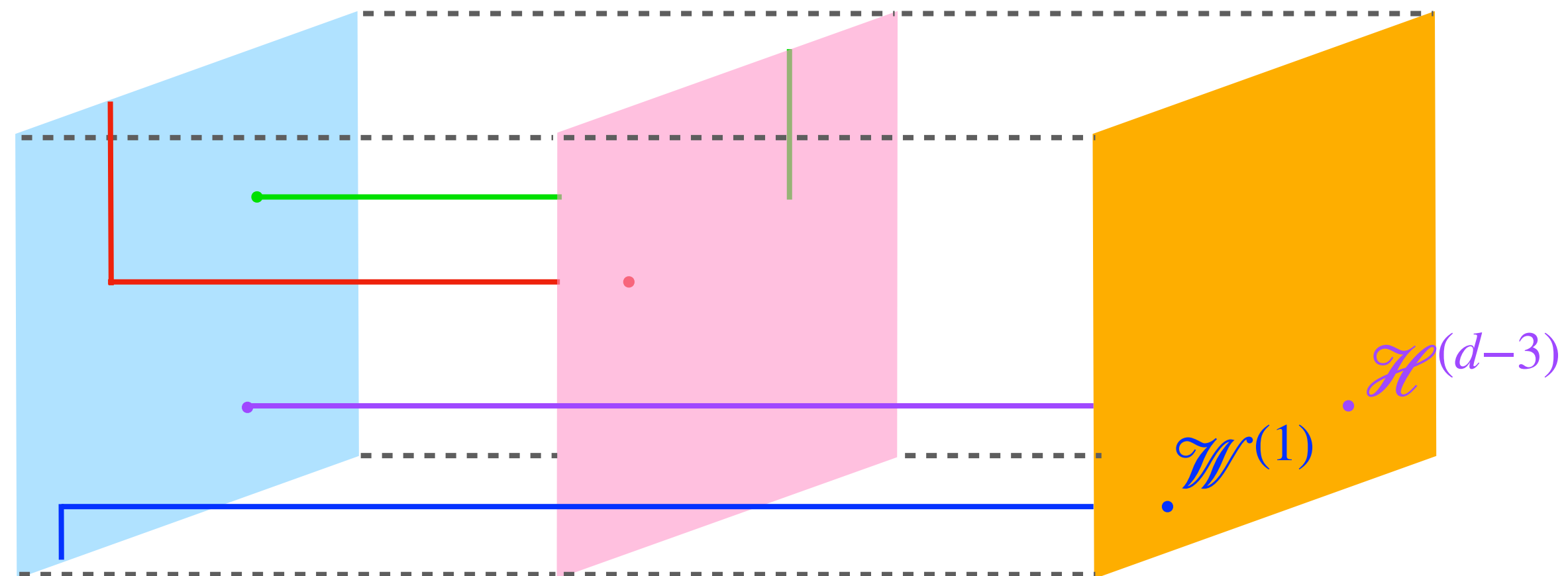
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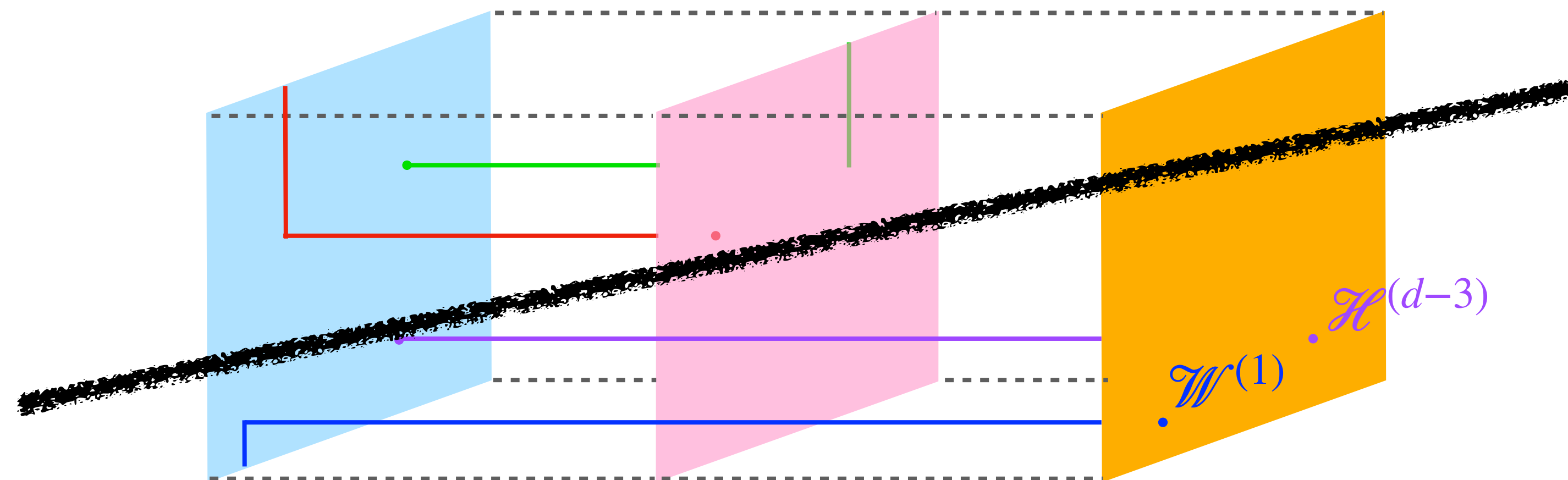
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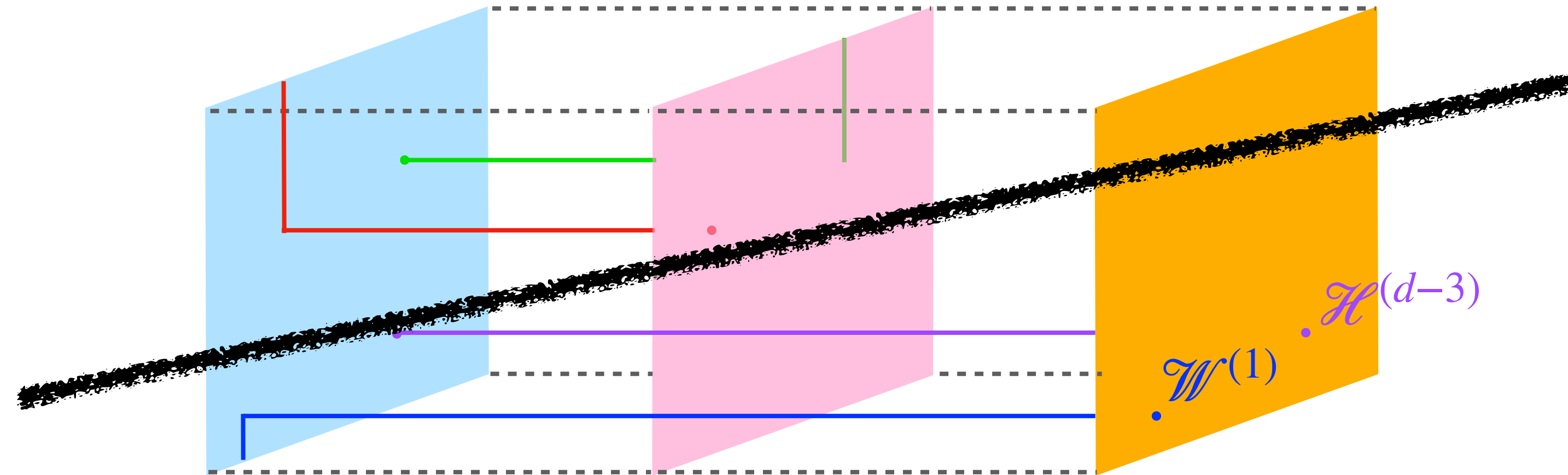
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## Constraints on IR phases

$$\mathcal{S}^{\text{Sym}} = \frac{1}{2} \int_{W_{D+1}} (C_2 \delta B_{d-2} + A_2 \delta A_{d-2} + \frac{1}{2} A_2 \delta B_{d-2})$$

- Anomaly-term in SymTFT distinguishes frozen  $\mathfrak{sp}(N)$  SYM from “gauge sector + decoupled defects”.
  - In particular: gives mixed anomaly for absolute  $Sp(N) + \mathbb{Z}^{(d-3)}$  theory.
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  - Trivial with SUSY, but: symmetry structure depends only on topology of  $X$
- ➔ Should hold with SUSY breaking but topology preserving deformations of geometric engineering setup!

# Summary & Outlook

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  - Constructions of  $\mathfrak{sp}$ -SYM in  $d \geq 6$  have interfaces for additional  $1/(d-3)$ -form symmetries.
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- Extension to 5d theories / CY3’s?
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*Thank you!*