



Synergy and complementarity between neutrino physics and high-intensity frontiers

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Neutrino oscillations: gateway to new physics

- **Neutrino oscillations** provided the 1st laboratory **evidence of New Physics**
⇒ **SM must be clearly extended** (or embedded in a larger framework)!

Several **possible models** successfully account for ν data
such extensions might even allow **to address SM caveats**

[↵ presentations by A. De Gouvea, P. Hernandez, S. Antusch, A. Boyarski, ...]

- Extend the SM: but how? **Hundredths of (motivated) theoretical constructions!!**

André de Gouvêa

Northwestern

Fork on the Road: Are Neutrinos Majorana or Dirac Fermions?



September 25, 2017

Overview

Neutrino oscillations: gateway to new physics

- **Neutrino oscillations** provided the 1st laboratory **evidence of New Physics**
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[↵ presentations by A. De Gouvea, P. Hernandez, S. Antusch, A. Boyarski, P. Paradisi, ...]

- Gateway to new **experimental signals** (deviation from SM) in the **lepton sector**:
Lepton Number Violation (if Majorana) - **$0\nu2\beta$, meson decays, colliders, ...**
Lepton flavour universality violation - weak boson and meson decays.. (e.g. R_K)
Electric dipole moments and **Anomalous magnetic moments**
Charged lepton flavour violation



- Rare processes searched for at **high-intensity facilities**
⇒ Learn about **neutrino mass models**! (At least **probe and disfavour...**)

Brief summary

- ▶ Leptonic high-intensity observables: signs of New Physics
- ▶ Observables and experimental status
 - Lepton number violation (observables at high and low energies)
 - Charged lepton flavour violation
 - CP violation: Electric dipole moments
 - Further observables
- ▶ Model-independent approaches to New Physics
- ▶ Models of neutrino mass generation: signals at high-intensities
 - Ad-hoc extensions
 - Seesaw realisations
 - Larger frameworks
- ▶ Overview & discussion

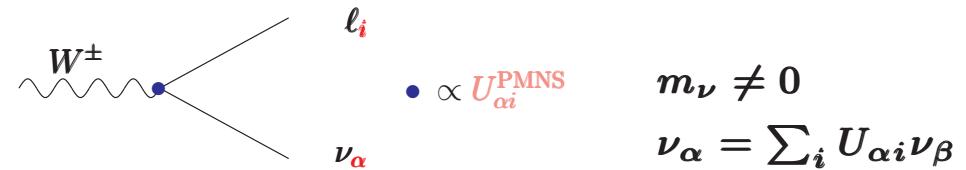
Leptonic observables: signs of New Physics

- In the **Standard Model**: (strictly) **massless neutrinos**
conservation of total lepton number & lepton flavours
tiny leptonic EDMs (at 4-loop level.. $d_e^{\text{CKM}} \leq 10^{-38} e \text{ cm}$)

- Extend the SM to accommodate $\nu_\alpha \leftrightarrow \nu_\beta$

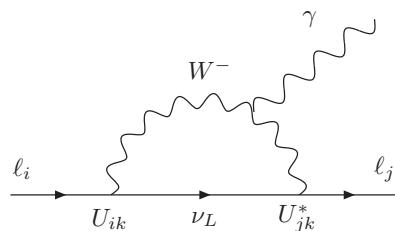
Assume **most minimal** extension **SM_{m_ν}**

[**SM_{m_ν}** = “ad-hoc” m_ν (Dirac), U_{PMNS}]



- In the **SM_{m_ν}**: (**total**) Lepton number conserved; what about lepton flavours? And CP?

- **SM_{m_ν} - cLFV possible??**



$$\text{BR}(\mu \rightarrow e\gamma) \propto \left| \sum U_{\mu i}^* U_{ei} \frac{m_{\nu_i}^2}{M_W^2} \right|^2 \sim 10^{-54}$$

[Petcov, '77]

Possible - yes... but **not observable!!**

- **SM_{m_ν} - observable EDMs?** Contributions from δ_{CP} (2-loop)... still $d_e^{\text{lep}} \leq 10^{-35} e \text{ cm}$

Leptonic observables: signs of New Physics

- ▶ Explore the underlying **synergy** between
ν **physics** and **high-intensity** observables
to constrain the **New Physics** model
at the origin of neutrino phenomena

- ▶ And keep an open eye on **collider** searches and **new oscillation** phenomena !
(not addressed here...)



 **Leptonic observables: current status**

[\rightsquigarrow WG4 presentations and reviews, Monday-Friday!]

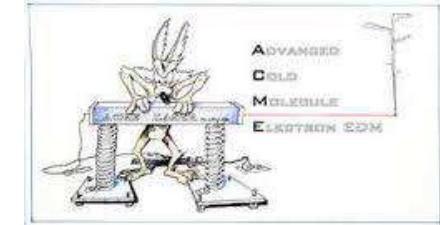
Leptonic dipole moments

► Electric dipole moments of charged leptons

$$\mathcal{L}_{\text{EDM}} = -i/2 \mathbf{d}_\ell \bar{\ell} \sigma^{\mu\nu} \gamma_5 \ell F_{\mu\nu}$$



EDM ($e \text{ cm}$)	Current bounds	Future sensitivity
$ d_e $	8.7×10^{-29} [ACME]	$\mathcal{O}(10^{-30})$ [ACME]
$ d_\mu $	1.9×10^{-19} [Muon g-2]	$\mathcal{O}(10^{-21})$ [g-2/EDM Coll.]
$ \text{Re}(d_\tau) $	4.5×10^{-17} [Belle]	-
$ \text{Im}(d_\tau) $	2.5×10^{-17} [Belle]	-



► (Anomalous) magnetic moments of charged leptons

$$\vec{\mu} = \mathbf{g}_\ell \frac{e}{2m_\ell} \vec{S} \Rightarrow \mathbf{a}_\ell = \frac{1}{2} (\mathbf{g}_\ell - 2)$$

a_e : Best determination of α_{em} ...

$$a_e^{\text{the}} = 0.001159652181643(764) \leftrightarrow 5^{\text{th}} \text{ order in QED (12,672 diags)}$$

$$a_e^{\text{exp}} = 0.00115965218073(28)$$

a_μ : Current tension between theory and experiment →

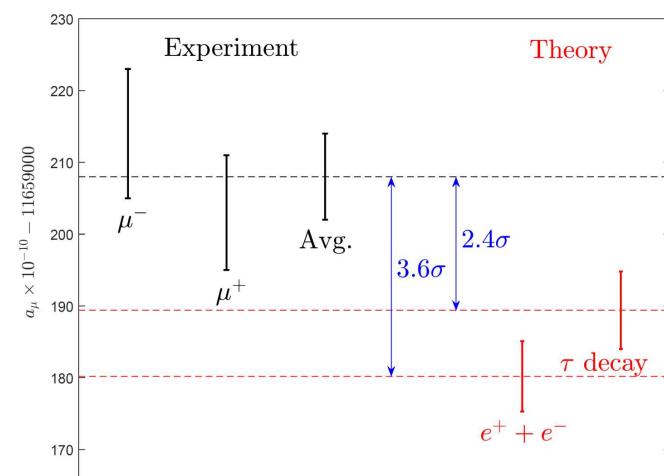
Very sensitive probe of New Physics close to Λ_{EW}

If δa_μ confirmed \rightsquigarrow discrepancies for $a_{e,\tau}$ and d_ℓ !

a_τ : Short tau lifetime...

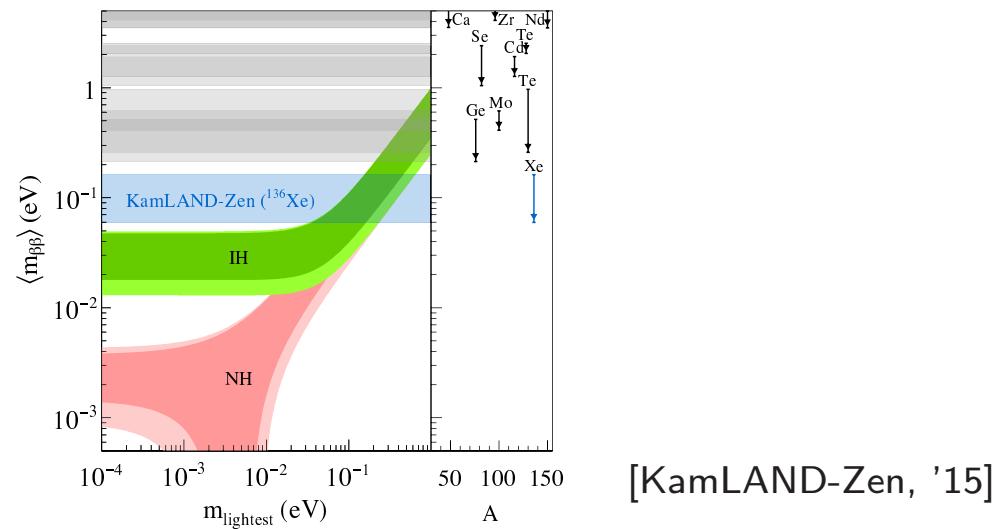
$$a_\tau^{\text{the}} = 0.00117721(5) [0701260]$$

$$-0.007 < a_\tau^{\text{exp}} < 0.005 [1601.07987]$$



Lepton number violation: $\Delta L = 2$ observables and searches

► Neutrinoless double beta decays

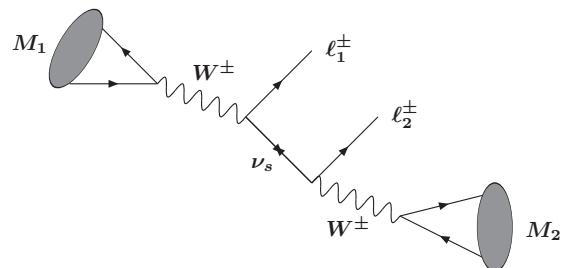


Experiment	$ m_{ee} $ (eV)
EXO-200 (4 yr)	0.075 - 0.2
nEXO (5 yr)	0.012 - 0.029
nEXO (5 yr + 5 yr w/ Ba tagging)	0.005 - 0.011
KamLAND-Zen (300 kg, 3 yr)	0.045 - 0.11
GERDA phase II	0.09 - 0.29
CUORE (5 yr)	0.051 - 0.133
SNO+	0.07 - 0.14
SuperNEMO	0.05 - 0.15
...	...

► LNV in semileptonic tau and/or meson decays

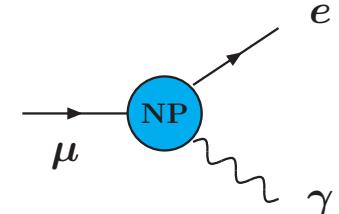
LNV decay	Current Bound	
	$\ell = e, \ell' = e$	$\ell = \mu, \ell' = \mu$
$K^- \rightarrow \ell^- \ell'^- \pi^+$	6.4×10^{-10}	1.1×10^{-9}
$D^- \rightarrow \ell^- \ell'^- \pi^+$	1.1×10^{-6}	2.2×10^{-8}
$D^- \rightarrow \ell^- \ell'^- K^+$	9.0×10^{-7}	1.0×10^{-5}
$B^- \rightarrow \ell^- \ell'^- \pi^+$	2.3×10^{-8}	4.0×10^{-9}
$B^- \rightarrow \ell^- \ell'^- K^+$	3.0×10^{-8}	4.1×10^{-8}
$B^- \rightarrow \ell^- \ell'^- \rho^+$	1.7×10^{-7}	4.2×10^{-7}
$B^- \rightarrow \ell^- \ell'^- D^+$	2.6×10^{-6}	6.9×10^{-7}

LNV decay	Current Bound	
	$\ell = e$	$\ell = \mu$
$\tau^- \rightarrow \ell^+ \pi^- \pi^-$	2.0×10^{-8}	3.9×10^{-8}
$\tau^- \rightarrow \ell^+ \pi^- K^-$	3.2×10^{-8}	4.8×10^{-8}
$\tau^- \rightarrow \ell^+ K^- K^-$	3.3×10^{-8}	4.7×10^{-8}



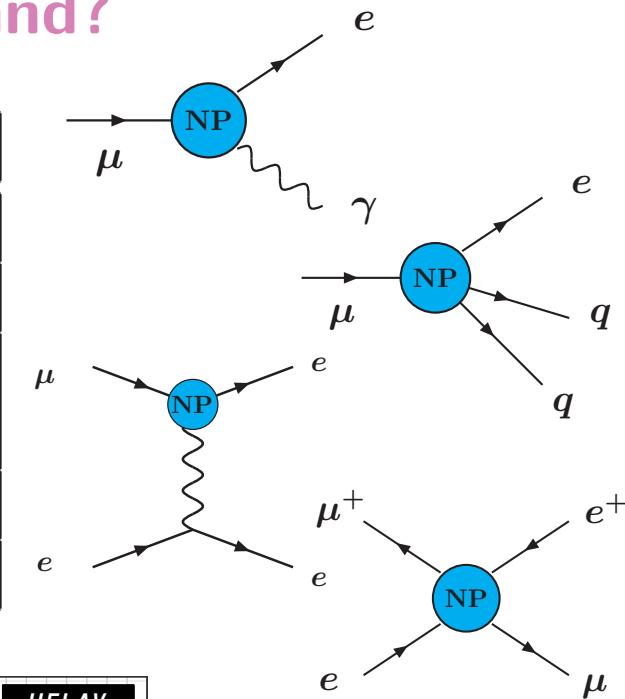
Signals of Lepton Flavour Violation

- ▶ Neutrino oscillations [ν-dedicated experiments]
- ▶ Rare leptonic decays and transitions [high-intensity facilities]
 - $\ell_i \rightarrow \ell_j \gamma, \ell_i \rightarrow 3\ell_j$, mesonic τ decays...
 - nucleus assisted $\mu - e$ transitions (also LNV!), Muonium channels...
- ▶ Meson decays: lepton flavour violating decays - $B \rightarrow \tau \mu, \dots$ [high-intensity; LHCb]
 - cLFV & Lepton Number violating decays - $B \rightarrow D e^- \mu^-, \dots$
 - violation of lepton flavour universality (e.g. R_K)
- ▶ Rare (new) heavy particle decays (typically model-dependent) [colliders]
 - SM boson decays: $H \rightarrow \tau \mu, Z \rightarrow \ell_i \ell_j$
 - SUSY $\tilde{\ell}_i \rightarrow \ell_j \chi^0$; FV KK-excitation decays; ...
 - LFV final states: for example, $e^\pm e^- \rightarrow e^\pm \mu^- + E_{\text{miss}}$
- ▶ And many others ... all absent in the SM!

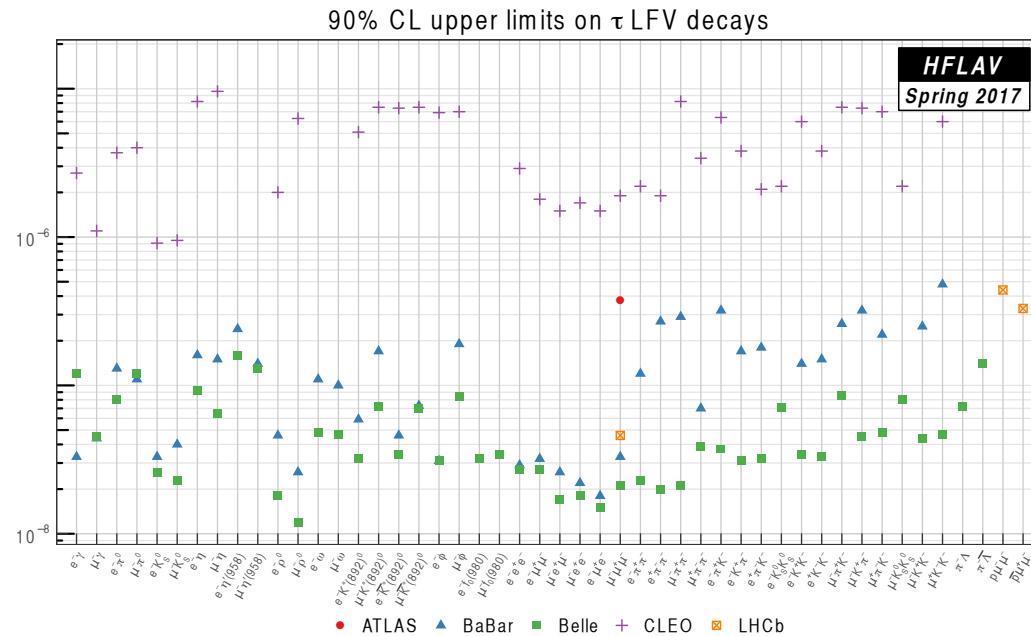


Searches for cLFV: where do we stand?

Observable	Bound (90% C.L.)	future sensitivity
$\text{BR}(\mu \rightarrow e\gamma)$	4.2×10^{-13}	4×10^{-14} [MEG II]
$\text{BR}(\mu \rightarrow 3e)$	1.0×10^{-12}	10^{-15} [Mu3e Phase I]
$\text{BR}(\mu - e, N)$	7×10^{-13} (Au)	10^{-14} [SiC, DeeMe] 10^{-17} [AI, Mu2e/COMET]
$\text{BR}(\mu^- e^- \rightarrow e^- e^-)$	—	10^{-17} [AI, COMET ?]
$P(\text{Mu} - \overline{\text{Mu}})$	8.3×10^{-11}	10^{-14} [FNAL ?]



[\rightsquigarrow Y. Ulrich (NLO!)]



[~> See WG4 contributions
Monday-Friday!]

Further bounds: $\text{BR}(Z \rightarrow \ell_i \ell_j)$, ..., $\text{BR}(K, D, B \rightarrow (h) \ell_i \ell_j)$, ... [\rightsquigarrow presentation by R. Bernstein!]

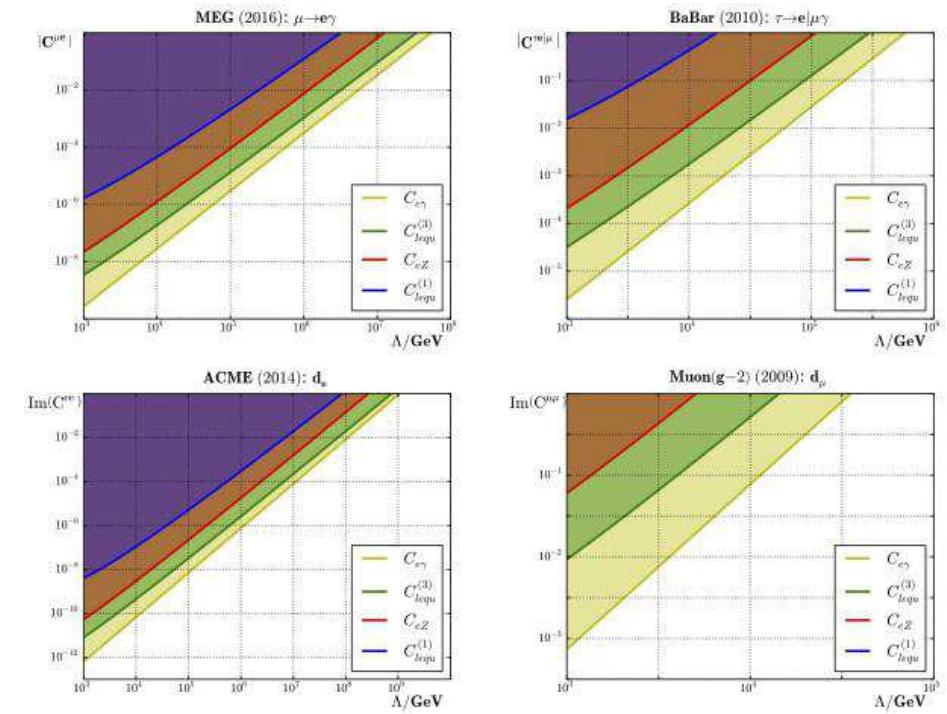
► Model independent approach

Neutrinoless radiative decay

$$\text{Br}(\mu \rightarrow e\gamma) = \frac{\alpha_e m_\mu^5}{\Lambda^4 \Gamma_\mu} \left(|C_L^D|^2 + |C_R^D|^2 \right).$$

Neutrinoless three-body decay

$$\begin{aligned} \text{Br}(\mu \rightarrow 3e) &= \frac{\alpha_e^2 m_\mu^5}{12\pi \Lambda^4 \Gamma_\mu} \left(|C_L^D|^2 + |C_R^D|^2 \right) \left(8 \log \left[\frac{m_\mu}{m_e} \right] - 11 \right) \\ &+ \frac{m_\mu^5}{3(16\pi)^3 \Lambda^4 \Gamma_\mu} \left(|C_{ee}^{S\,LL}|^2 + 16 |C_{ee}^{V\,LL}|^2 + 8 |C_{ee}^{V\,LR}|^2 \right. \\ &\quad \left. + |C_{ee}^{S\,RR}|^2 + 16 |C_{ee}^{V\,RR}|^2 + 8 |C_{ee}^{V\,RL}|^2 \right). \end{aligned}$$



[↵ presentation by G. M. Pruna]

- ▶ Accounting for neutrino masses and mixings:
SM extensions ...
... and high-intensity observables

Theoretical frameworks

- ▶ Simplified “toy models” for phenomenological analyses: $\text{SM} + \nu_s$
 - “ad-hoc” construction (no specific assumption on mechanism of mass generation)
 - encodes the effects of N additional sterile states (well-motivated NP candidates)
 - in a **single one** [... Not to be confused with oscillation anomaly solution!...]
- ▶ Complete SM extensions accounting for ν masses and mixings
 - Models of ν -mass generation - Standard seesaws [type I, type II, type III] & variants
 - Low-scale, νMSM , Inverse Seesaw (ISS), ...
 - Additional states: Multi-Higgs doublet models,
 - leptoquarks, Z' , vector-like, ...
 - Extended frameworks: extra dimensions, ...
 - SUSY seesaw,
 - Left-Right models, GUTs, ...
- ▶ High-intensity probes to distinguish between them!

- ▶ Minimal toy-model: SM + ν_s

Assuming that New Physics is encoded into such a simple model,
what can we expect and learn?

“Toy model” for phenomenological analyses: SM + ν_s

- ▶ Assumptions: 3 active neutrinos + 1 sterile state $n_L = (\nu_{Le}, \nu_{L\mu}, \nu_{L\tau}, \nu_s^c)^T$
 interaction basis \rightsquigarrow physical basis $n_L = U_{4 \times 4} \nu_i$
- $$U_{4 \times 4}^T M U_{4 \times 4} = \text{diag}(m_{\nu_1}, \dots, m_{\nu_4}) \quad \text{“Majorana mass”: } \mathcal{L}_{\text{toy}} \sim n_L^T C M n_L$$

▶ Active-sterile mixing $U_{\alpha i}$:

rectangular matrix $\leftarrow \mathbf{U} = U|_{3 \times 4}$

▶ Left-handed lepton mixing \tilde{U}_{PMNS} :

3×3 sub-block, non-unitary!

$$U_{4 \times 4} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} & U_{\mu 4} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} & U_{\tau 4} \\ U_{s1} & U_{s2} & U_{s3} & U_{s4} \end{pmatrix}$$

▶ Physical parameters: 4 masses [3 light (mostly active) + 1 heavier (mostly sterile) states]

6 mixing angles $[\theta_{12}, \theta_{23}, \theta_{13}, \& \theta_{i4}]$ and 6 phases [(3 Dirac and 3 Majorana)]

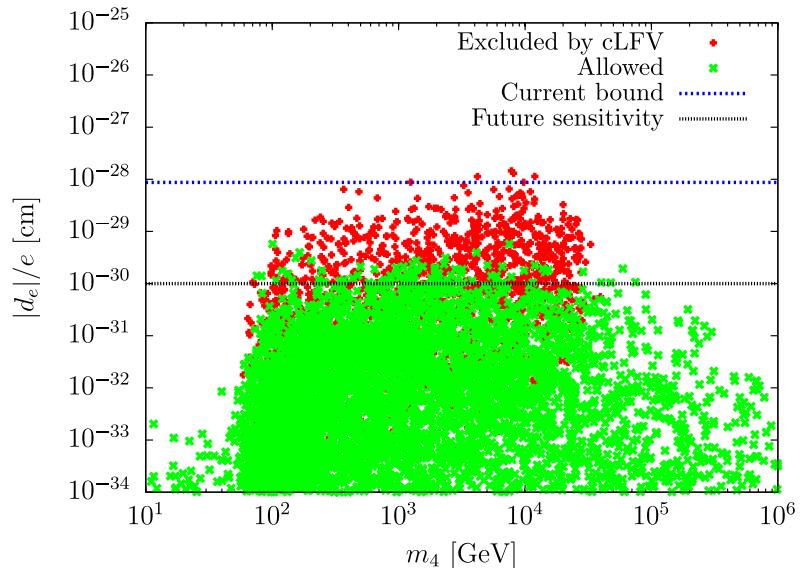
▶ Modified charged (W^\pm) and neutral (Z^0) current interactions:

$$\mathcal{L}_{W^\pm} \sim -\frac{g_w}{\sqrt{2}} W_\mu^- \sum_{\alpha=e,\mu,\tau} \sum_{i=1}^{3+n_S} \mathbf{U}_{\alpha i} \bar{\ell}_\alpha \gamma^\mu P_L \nu_i$$

$$\mathcal{L}_{Z^0} \sim -\frac{g_w}{2 \cos \theta_w} Z_\mu \sum_{i,j=1}^{3+n_S} \bar{\nu}_i \gamma^\mu \left[P_L (\mathbf{U}^\dagger \mathbf{U})_{ij} - P_R (\mathbf{U}^\dagger \mathbf{U})_{ij}^* \right] \nu_j$$

Sterile neutrinos: impact for lepton properties

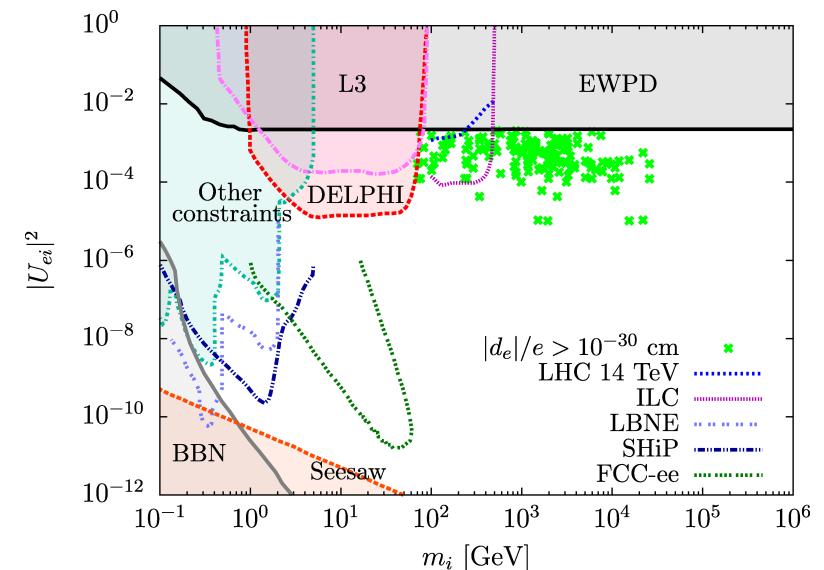
► Leptonic CP violation: electric dipole moments



- Majorana (and Dirac) phases \Rightarrow lepton EDMs
- Non-vanishing contributions: at least two sterile ν
- $|d_e|/e \geq 10^{-30}$ cm for $m_{\nu_{4,5}} \sim [100 \text{ GeV}, 100 \text{ TeV}]$

[Abada and Toma, '15]

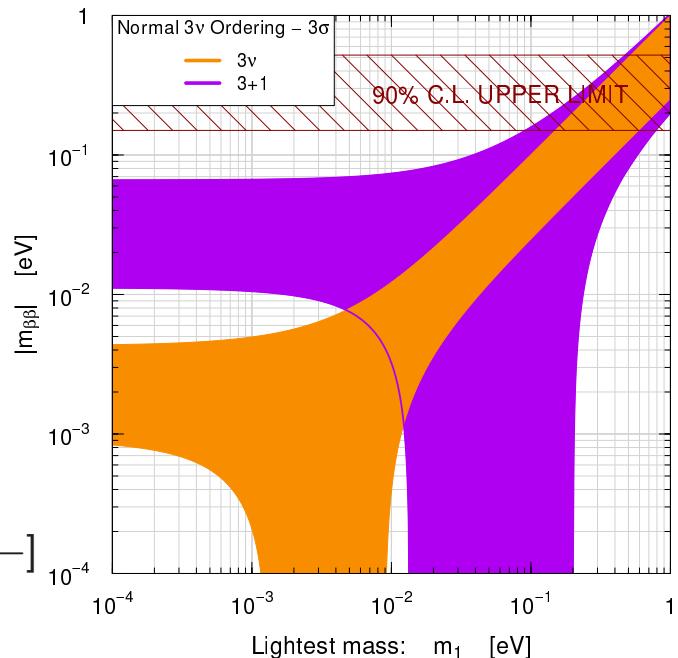
- Independent of active-sterile mixings
- Majorana contribution is dominant!
- EDM observation: suggest new sources of CPV
 \Rightarrow Majorana ν s? \rightsquigarrow Leptogenesis??
- Sterile states beyond (direct) collider reach...



Sterile neutrinos: impact for LNV observables

- ▶ Lepton number violation: $0\nu2\beta$ decays
- ▶ ν_s can strongly impact predictions for $|m_{ee}|$
⇒ augmented ranges for effective mass (*IO and NO*)
- ▶ Observation of $0\nu2\beta$ signal in future experiments
does not imply Inverted Ordering for light ν_s

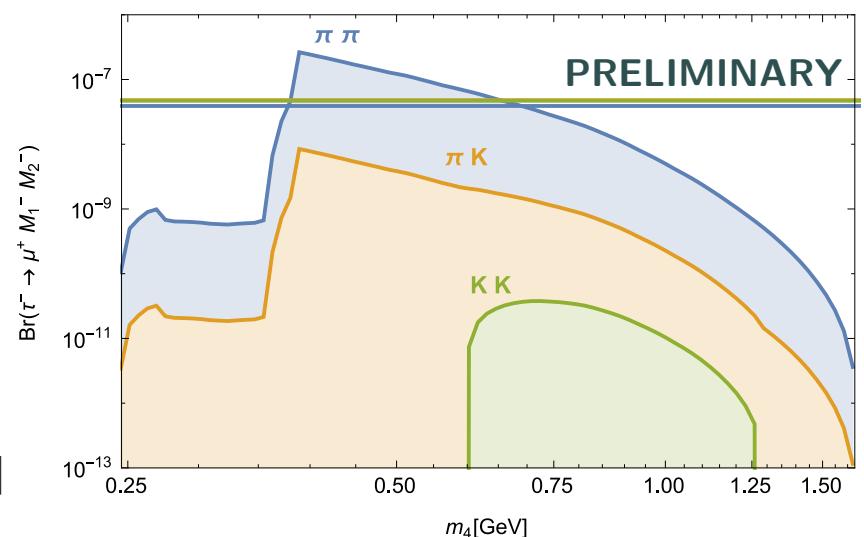
[Abada, De Romeri and AMT, '14; ...; Giunti et al, '15 ←]



- ▶ Lepton Number Violation in meson and τ decays

- ▶ If ν_s produced on-shell,
resonant enhancement of LNV decays
- $$M_1^- \rightarrow M_2^+ \ell^- \ell^- \text{ and } \tau^- \rightarrow \ell^+ M_1^- M_2^-$$

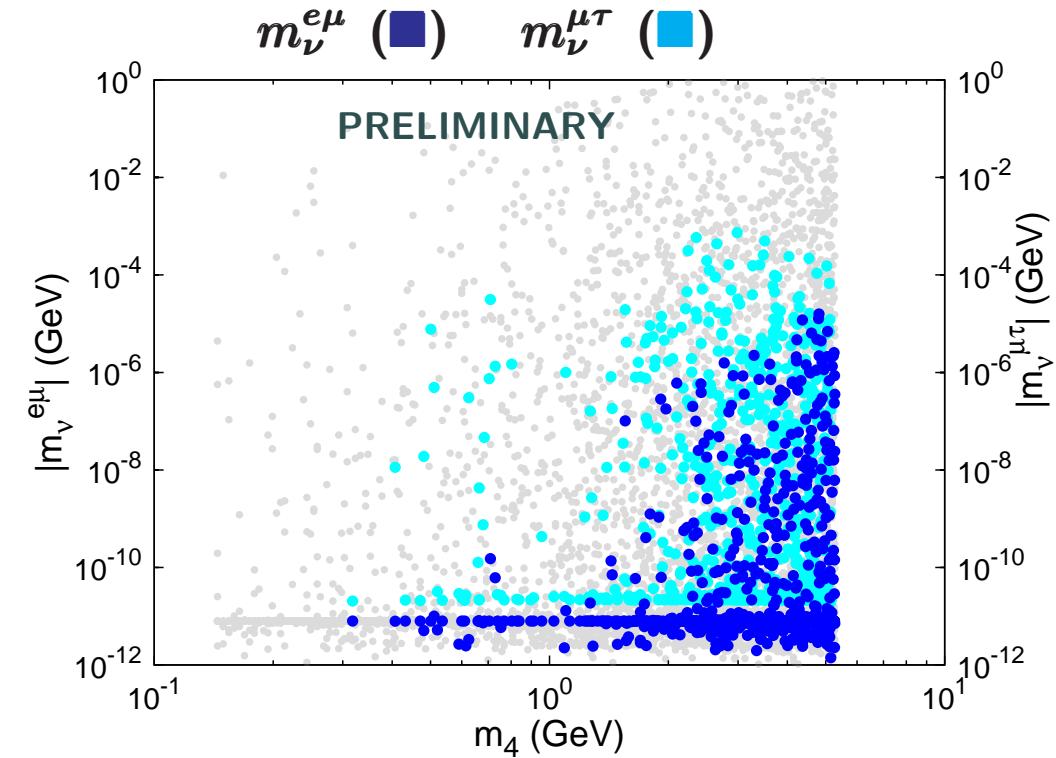
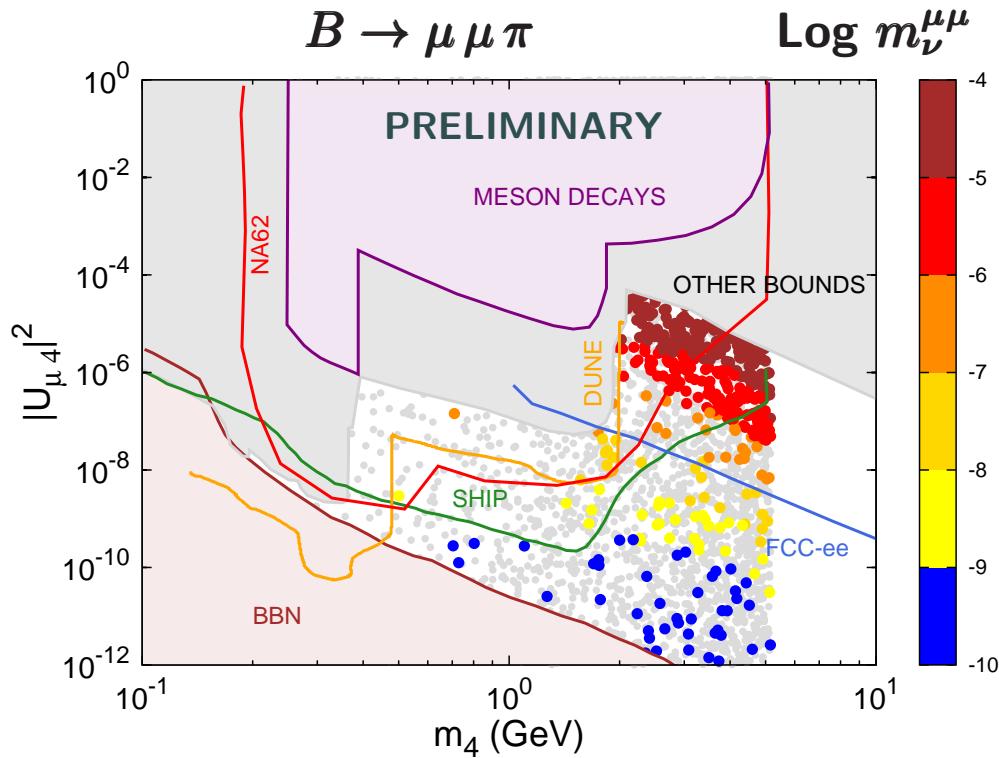
[Abada, De Romeri, Lucente, Toma, AMT, to appear]



Sterile neutrinos: impact for LNV meson and tau decays

- ▶ In addition to further constraining the active-sterile mixings [future sensitivities...]
- ▶ LNV meson and tau decays offer possibility to infer information on $m_\nu^{\ell_i \ell_j}$

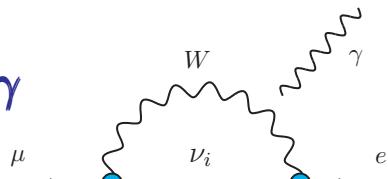
$$m_\nu^{\ell_\alpha \ell_\beta} \equiv \left| \sum_{i=1}^4 \frac{U_{\alpha i} m_i U_{\beta i}}{1 - m_i^2/p_{12}^2 + i m_i \Gamma_i / p_{12}^2} \right|$$



[Abada, De Romeri, Lucente, Toma, AMT, to appear]

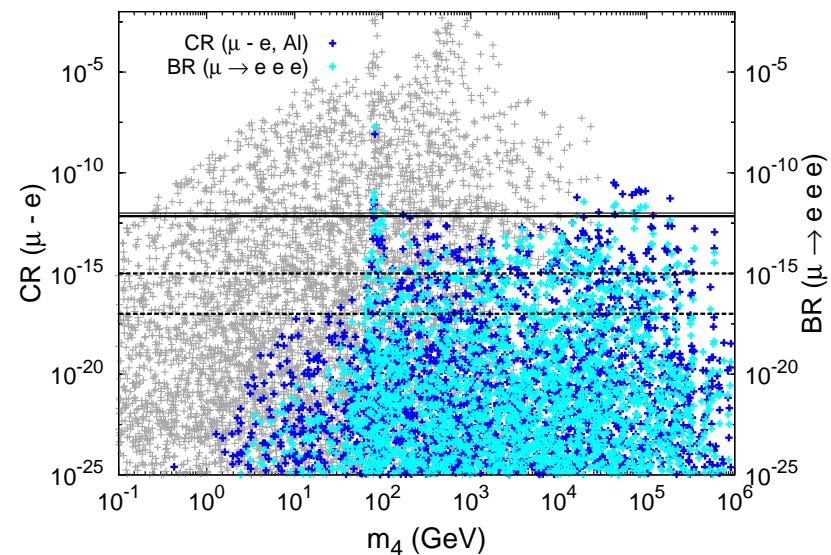
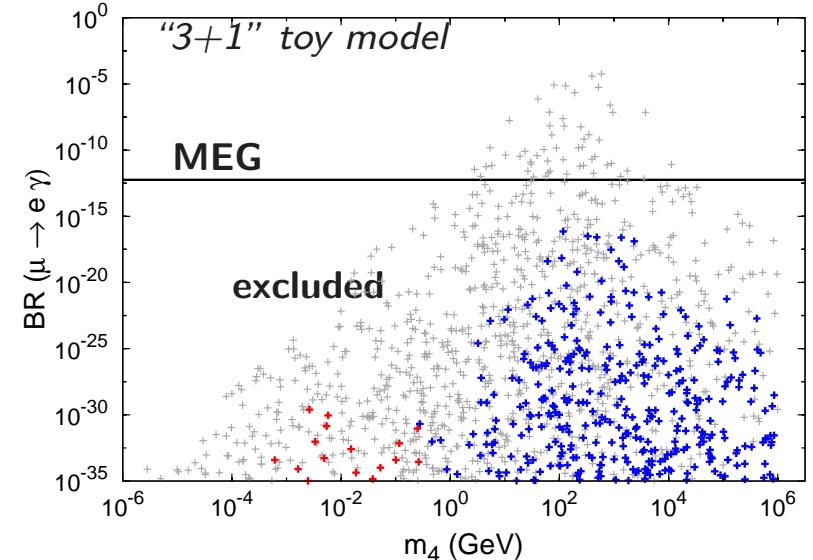
ν_s and cLFV: radiative and 3 body decays

- Radiative decays: $\ell_i \rightarrow \ell_j \gamma$

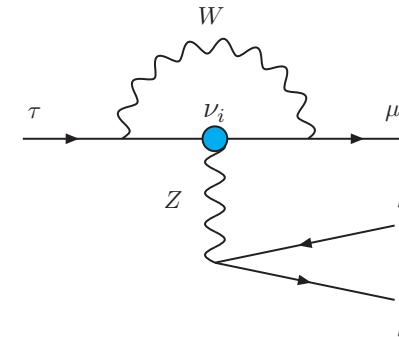


- Consider $\mu \rightarrow e \gamma$
- For $m_4 \gtrsim 10$ GeV sizable ν_s contributions
.. but precluded by other cLFV observables
as 3-body decays and nuclear conversion

- Three-body decays $\ell_i \rightarrow 3\ell_j$ (■) and conversion in Nuclei $\mu - e$ (■)



- For sterile states above EW scale, strongly dominated by Z penguin contributions



Sterile neutrinos: cLFV in “muonic atoms”

► cLFV $\mu^- - e^-$ conversion: $\mu^- + (A, Z) \rightarrow e^- + (A, Z)$

► Muonic atom decay: $\mu^- e^- \rightarrow e^- e^-$ [Koike et al, '10]

Coulomb interaction increases overlap between Ψ_{μ^-} and Ψ_{e^-}

Rate strongly enhanced in large Z atoms [Uesaka et al, '15-'16]

► cLFV in muonic atoms from ν_s :

$\mu^- e^- \rightarrow e^- e^-$ (■) vs

$\mu - e$ conversion (■) in Aluminium

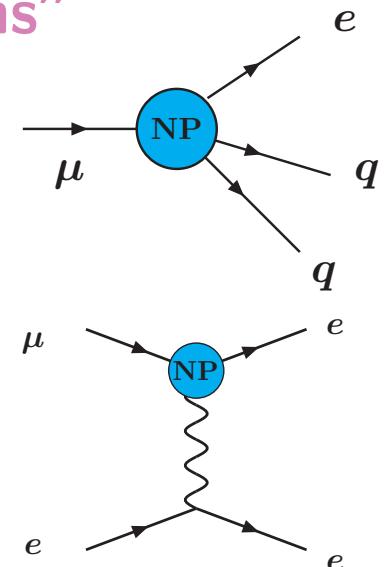
► For Aluminium, $CR(\mu - e)$ has

stronger experimental potential

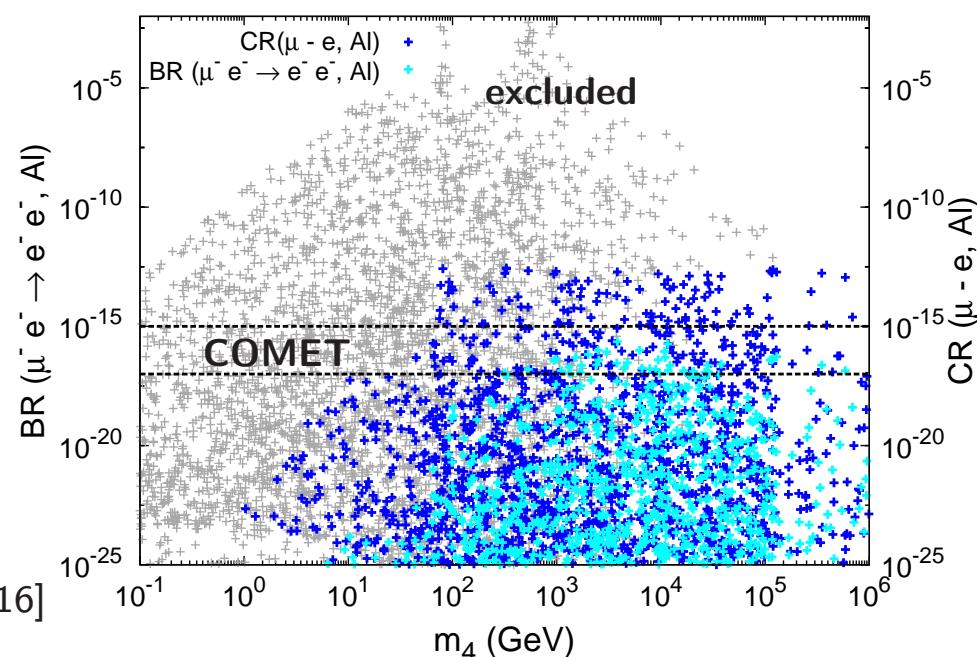
.. consider “heavy” targets to probe

$BR(\mu^- e^- \rightarrow e^- e^-)$

“3+1” toy model [Abada, De Romeri and AMT, '16]



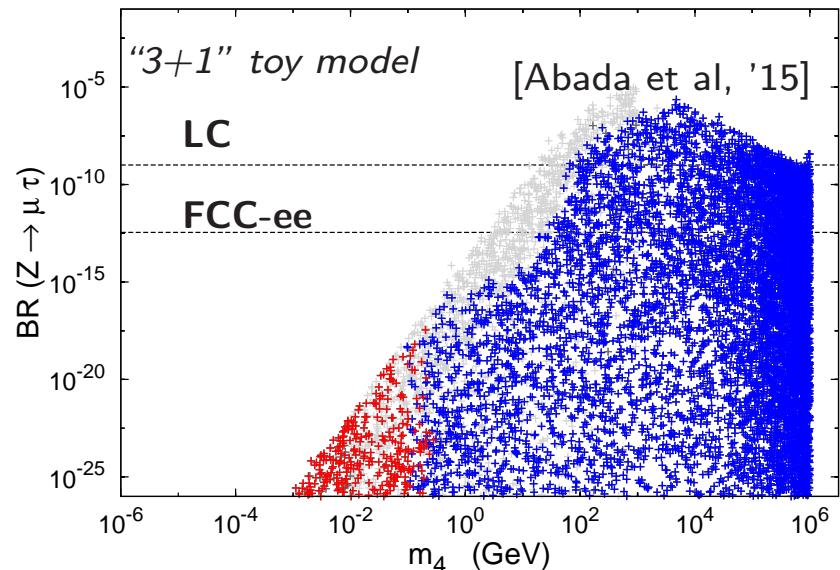
[↔ presentation by J. Sato]



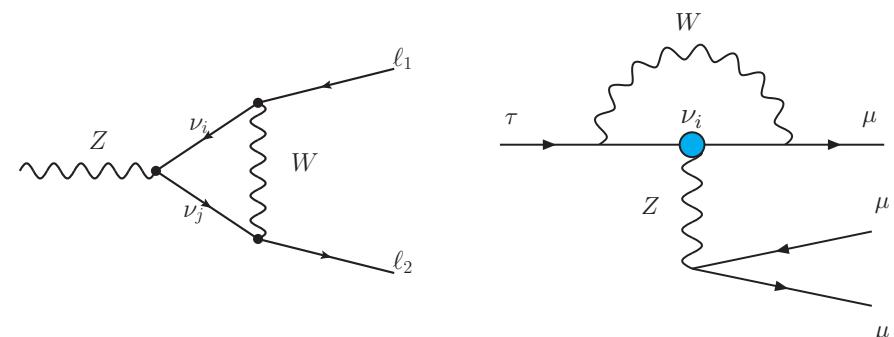
Sterile neutrinos and cLFV at higher energies

- cLFV Z decays at FCC-ee vs 3 body decays $\ell_i \rightarrow 3\ell_j$

[\rightsquigarrow presentation by S. Antusch]

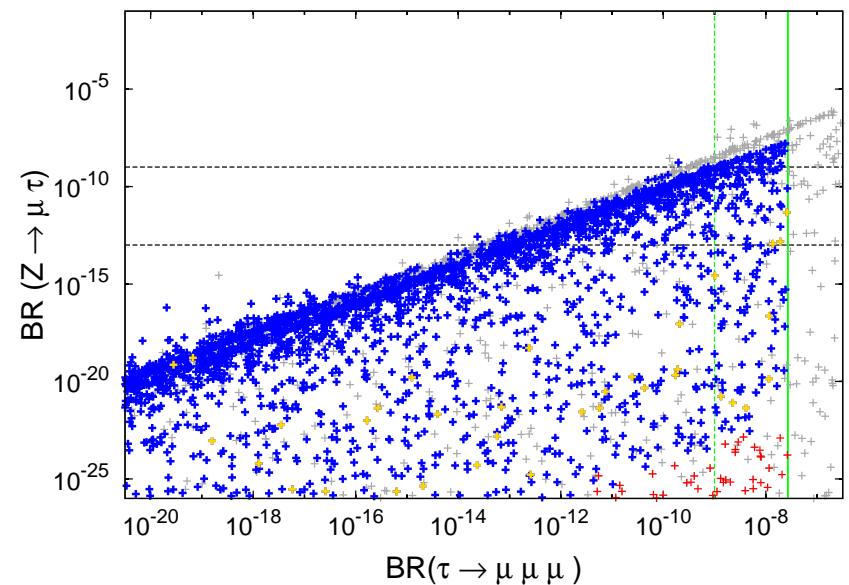


- Potentially **observable** at **Future Circular Collider**



- Recall: $\ell_i \rightarrow 3\ell_j$ dominated by Z penguins

- Strong correlation between $Z \rightarrow \mu\tau$ and $\tau \rightarrow 3\mu$
- **Probe $\mu - \tau$ cLFV beyond SuperB reach**
- **Complementarity probes** of ν_s cLFV
at **low- and high energies!**

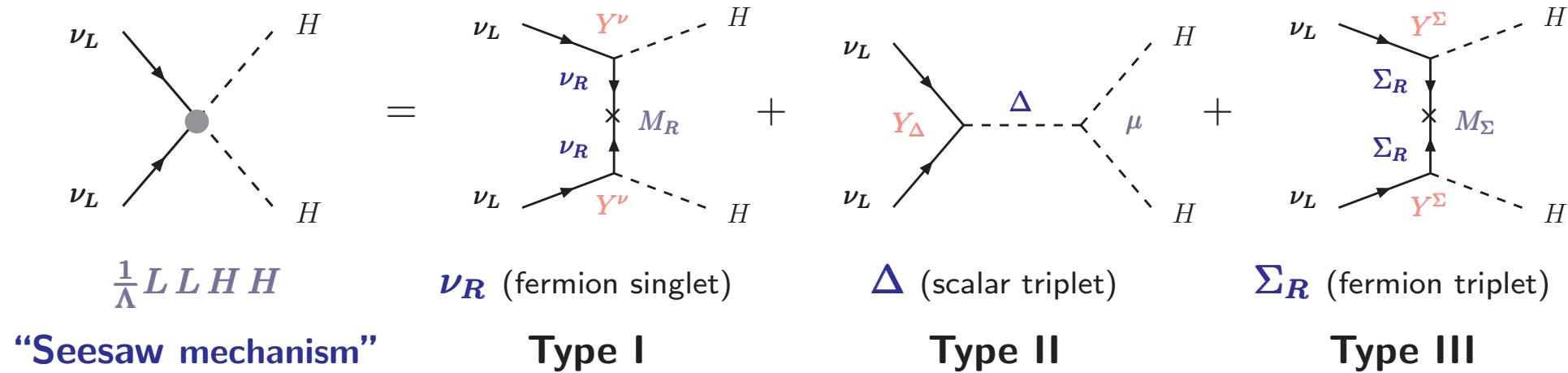
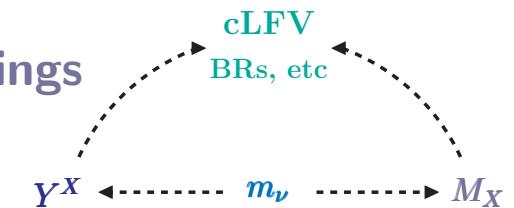




► Models of neutrino mass generation

The seesaw mechanism

★ **Seesaw mechanism:** explain **small ν masses** with “natural” couplings via new dynamics at “heavy” scale



- **Observables:** depend on **powers of Y^ν** \rightsquigarrow large rates \Rightarrow sizable Y^ν
and on the **mass of the (virtual) NP propagators**

- **Fermionic seesaws:** $Y^\nu \sim \mathcal{O}(1) \Rightarrow M_{\text{new}} \approx 10^{13-15} \text{ GeV!}$

Suppression of rates due to the **large mass of the mediators!**



- **Low scale seesaws:** rich phenomenology at **high-intensities!** (and also at LHC)

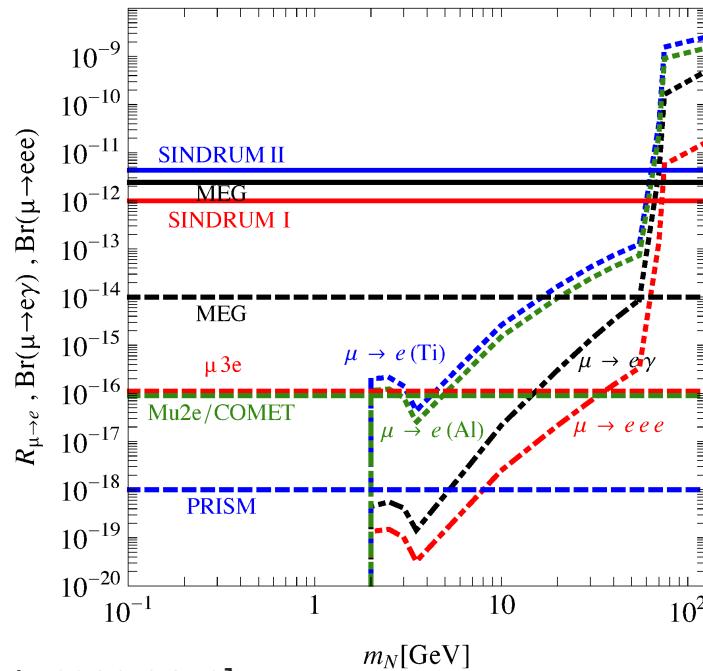
Low scale type I seesaw

- Addition of 3 “heavy” Majorana RH neutrinos to SM; $\text{MeV} \lesssim m_{N_i} \lesssim 10^{\text{few}} \text{TeV}$

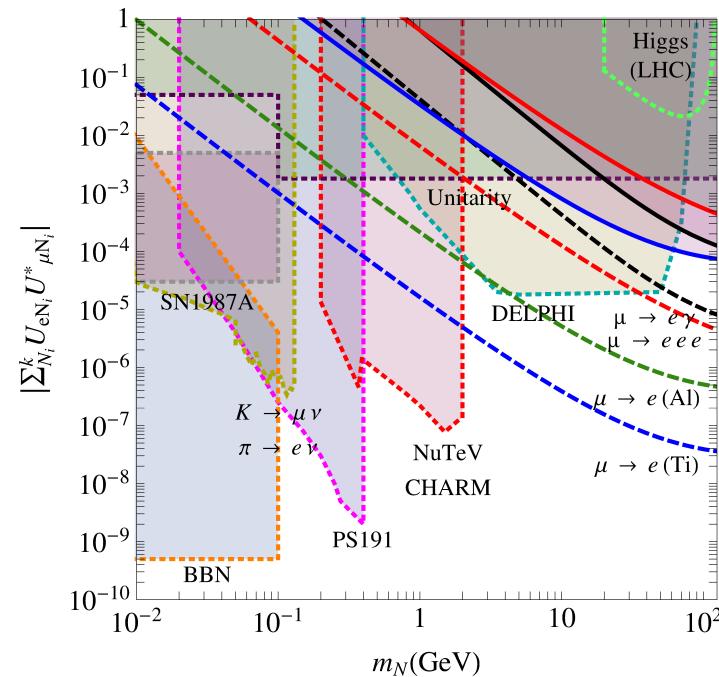
- Spectrum and mixings: $m_\nu \approx -v^2 Y_\nu^T M_N^{-1} Y_\nu$ $\mathbf{U}^T \mathcal{M}_\nu^{6 \times 6} \mathbf{U} = \text{diag}(m_i)$

$$\mathbf{U} = \begin{pmatrix} \mathbf{U}_{\nu\nu} & U_{\nu N} \\ U_{N\nu} & U_{NN} \end{pmatrix} \quad \mathbf{U}_{\nu\nu} \approx (1 - \varepsilon) \mathbf{U}_{\text{PMNS}} \quad \text{Non-unitary leptonic mixing } \tilde{\mathbf{U}}_{\text{PMNS}}!$$

- Heavy states do not decouple \Rightarrow modified neutral and charged leptonic currents
- Rich phenomenology at **high-intensity/low-energy** and at colliders!



[Alonso et al, 1209.2679]



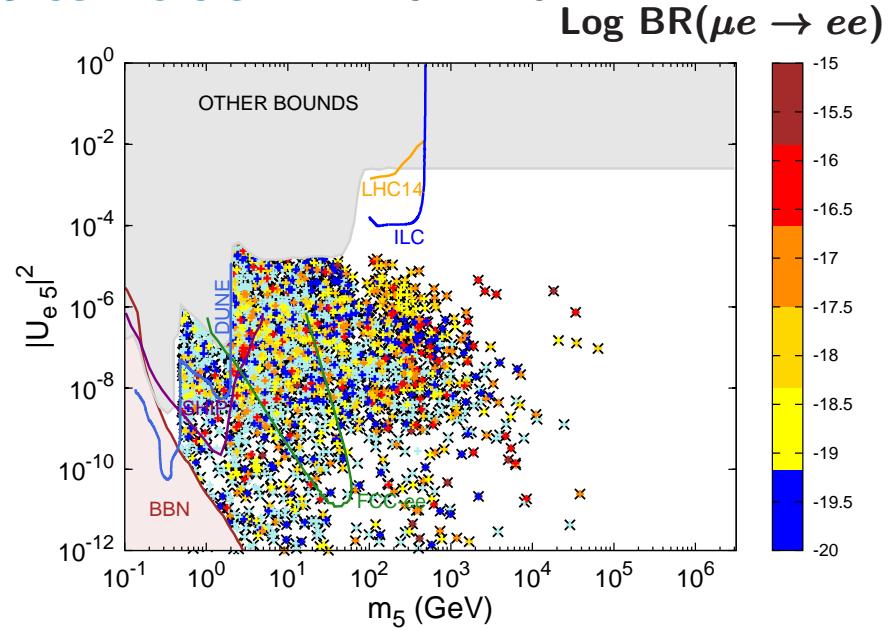
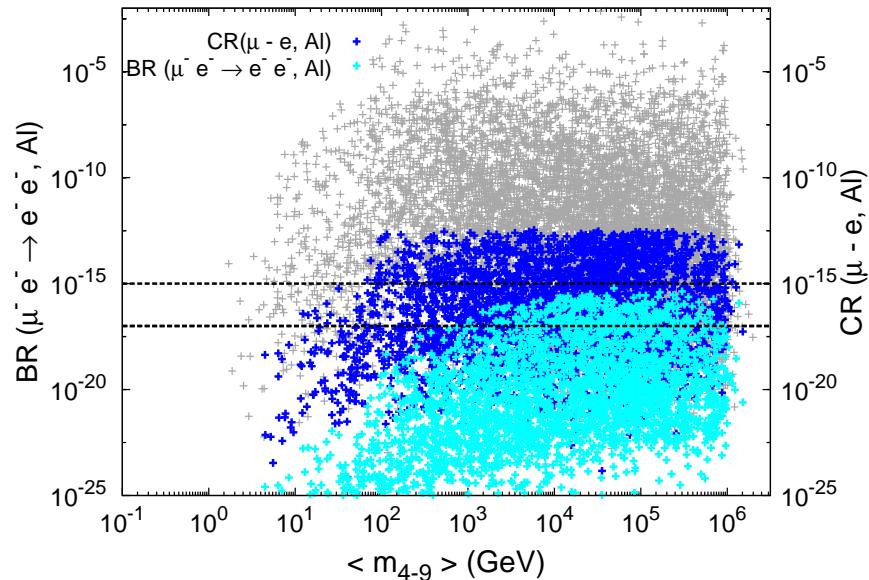
(see also Dinh et al, '12-'14)

Low scale: Inverse Seesaw (ISS)

- Addition of **3 “heavy” RH neutrinos** and **3 extra “sterile” fermions X** to the SM

$$\blacktriangleright \mathcal{M}_{\text{ISS}}^{9 \times 9} = \begin{pmatrix} 0 & Y_\nu v & 0 \\ Y_\nu^T v & 0 & M_R \\ 0 & M_R & \mu_X \end{pmatrix} \Rightarrow \begin{cases} \text{3 light } \nu : m_\nu \approx \frac{(Y_\nu v)^2}{(Y_\nu v)^2 + M_R^2} \mu_X \\ \text{3 pseudo-Dirac pairs} : m_{N^\pm} \approx M_R \pm \mu_X \end{cases}$$

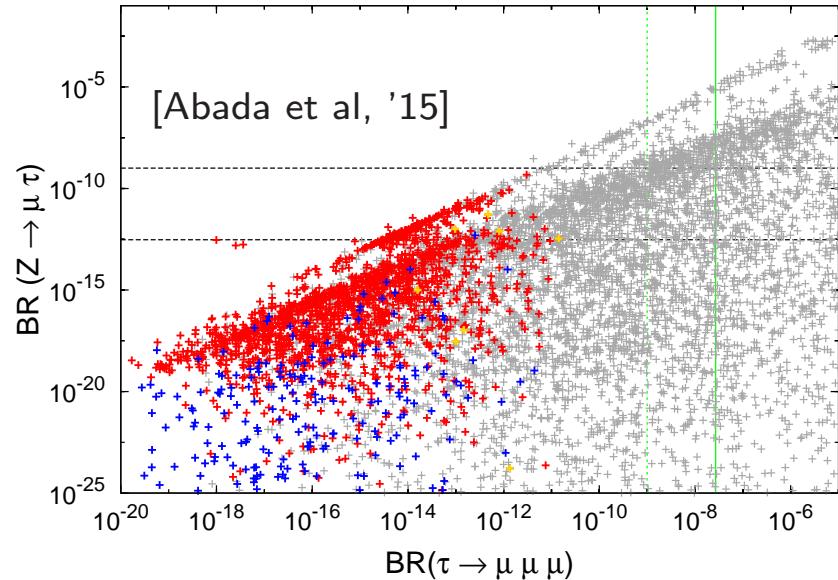
- Non-unitarity \tilde{U}_{PMNS} \Rightarrow modified neutral and charged leptonic currents
- New (virtual) states & modified couplings: **cLFV**, non-universality, signals at colliders!
- cLFV in muonic atoms: $\mu^- e^- \rightarrow e^- e^-$ vs $\mu - e$ conversion in Aluminium



[Abada, DeRomeri, AMT, '15]

Low scale: Inverse Seesaw (ISS)

► cLFV Z decays at FCC-ee vs 3 body decays $\ell_i \rightarrow 3\ell_j$

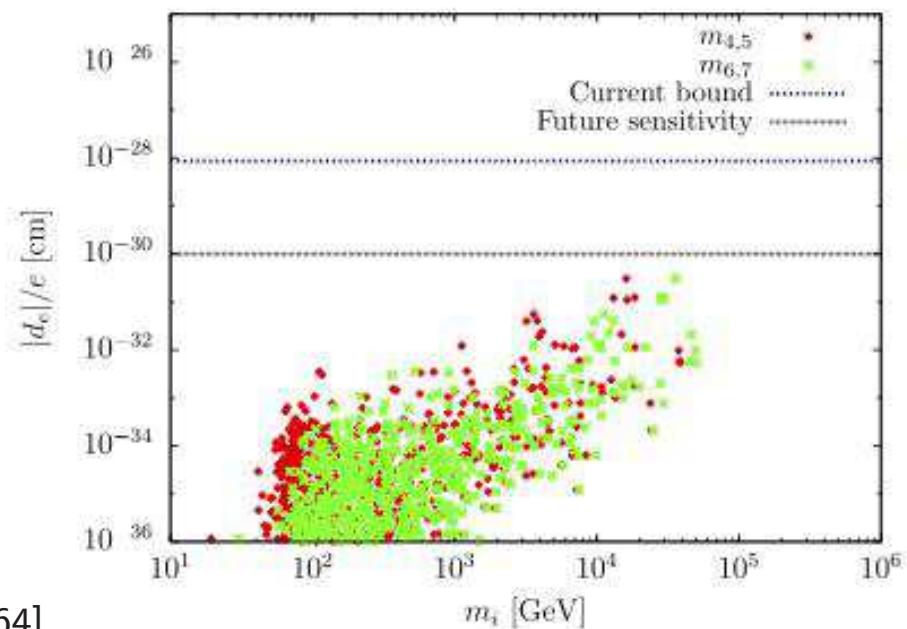


- ▶ Still dominated by Z penguin contributions
- ▶ Other cLFV bounds preclude large $\text{BR}(\tau \rightarrow 3\mu)$...
- Contrary to “3+1 toy model”, flavour textures & parameters constrained by ν data...
- ▶ Allows to probe $\mu - \tau$ cLFV beyond SuperB reach

► Leptonic CP violation: EDMs

- ▶ ISS contains additional sources of CPV!
- ▶ Majorana contributions nearly negligible
heavy steriles form pseudo-Dirac pairs
- ▶ Electron EDM beyond future sensitivity...

[Abada and Toma, 1611.03464]



The “triplet” seesaws

★ Weinberg operator realised via **triplet scalars Δ** (type II) or **fermions Σ** (type III)

► Very **distinctive signatures** for numerous **observables**: **cLFV example**

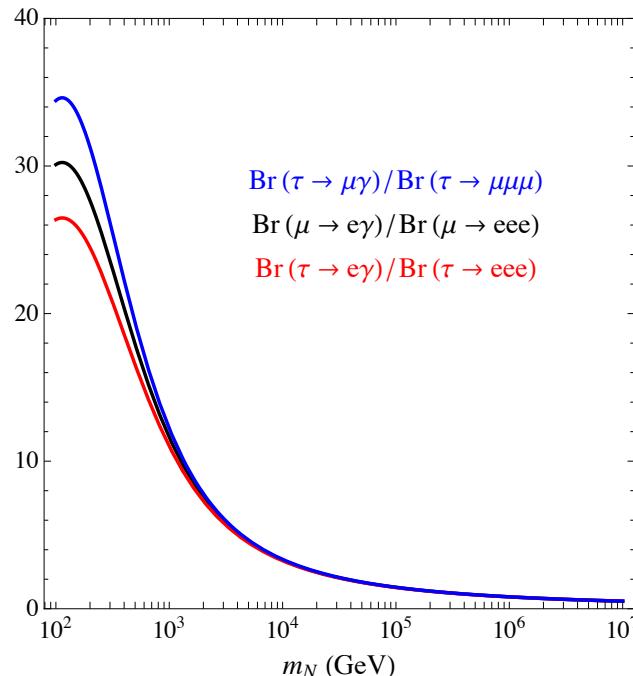
Type I: cLFV transitions at **loop level** (radiative, 3-body, conversion in Nuclei)

Type II: $\ell_i \rightarrow \ell_j \gamma$ & $\mu - e, N$ at loop level; 3-body decays $\ell_i \rightarrow 3\ell_j$ at **tree level!**

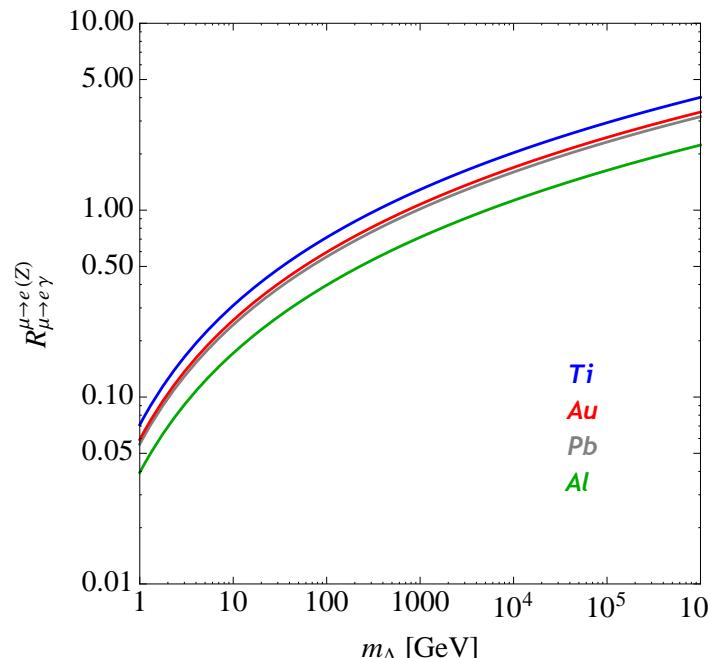
Type III: 3-body decays and coherent conversion at **tree-level!** $\ell_i \rightarrow \ell_j \gamma$ @ loop...

► Use **ratios of observables** to constrain and identify mediators!

Type I



Type II



Type III

$$\frac{\text{BR}(\mu \rightarrow e\gamma)}{\text{BR}(\mu \rightarrow 3e)} = 1.3 \times 10^{-3}$$

$$\frac{\text{BR}(\tau \rightarrow \mu\gamma)}{\text{BR}(\tau \rightarrow 3\mu)} = 1.3 \times 10^{-3}$$

$$\frac{\text{BR}(\mu \rightarrow e\gamma)}{\text{CR}(e - \mu, \text{Ti})} = 3.1 \times 10^{-4}$$

[Hambye, 2013]

The “triplet” seesaws

- **cLFV bounds** on the **seesaw mediators**: a comparative (“effective”) view

$$m_N \lesssim 100 \text{ TeV} \times \left(\frac{10^{-14}}{\text{BR}(\mu \rightarrow e\gamma)} \right)^{\frac{1}{4}} \times f(Y_{\ell_i \ell_j}^\nu)$$

- **Type I (singlet fermion):**

$$m_N \lesssim 300 \text{ TeV} \times \left(\frac{10^{-16}}{\text{BR}(\mu \rightarrow 3e)} \right)^{\frac{1}{4}} \times f(Y_{\ell_i \ell_j}^\nu)$$

$$m_N \lesssim 2000 \text{ TeV} \times \left(\frac{10^{-18}}{\text{CR}(\mu - e, \text{Ti})} \right)^{\frac{1}{4}} \times f(Y_{\ell_i \ell_j}^\nu)$$

$$m_\Delta \lesssim 70 \text{ TeV} \times \left(\frac{10^{-14}}{\text{BR}(\mu \rightarrow e\gamma)} \right)^{\frac{1}{4}} \times f(Y_{\ell_i \ell_j}^\Delta)$$

- **Type II (scalar triplet):**

$$m_\Delta \lesssim 2200 \text{ TeV} \times \left(\frac{10^{-16}}{\text{BR}(\mu \rightarrow 3e)} \right)^{\frac{1}{4}} \times f(Y_{\ell_i \ell_j}^\Delta)$$

$$m_\Delta \lesssim 600 \text{ TeV} \times \left(\frac{10^{-18}}{\text{CR}(\mu - e, \text{Ti})} \right)^{\frac{1}{4}} \times f(Y_{\ell_i \ell_j}^\Delta)$$

$$m_\Sigma \lesssim 100 \text{ TeV} \times \left(\frac{10^{-14}}{\text{BR}(\mu \rightarrow e\gamma)} \right)^{\frac{1}{4}} \times f(Y_{\ell_i \ell_j}^\Sigma)$$

- **Type III (fermion triplet):**

$$m_\Sigma \lesssim 1600 \text{ TeV} \times \left(\frac{10^{-16}}{\text{BR}(\mu \rightarrow 3e)} \right)^{\frac{1}{4}} \times f(Y_{\ell_i \ell_j}^\Sigma)$$

$$m_\Sigma \lesssim 20000 \text{ TeV} \times \left(\frac{10^{-18}}{\text{CR}(\mu - e, \text{Ti})} \right)^{\frac{1}{4}} \times f(Y_{\ell_i \ell_j}^\Sigma)$$

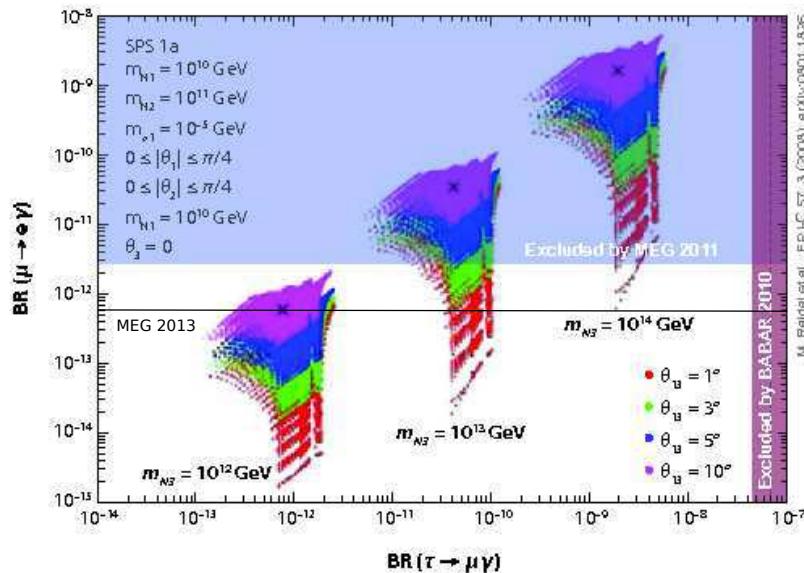
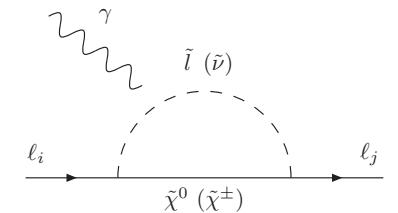
$f(Y_{\ell_i \ell_j}) \sim \text{combination of } \sqrt{Y} \sqrt{Y}$

► Embedding the seesaw in larger frameworks

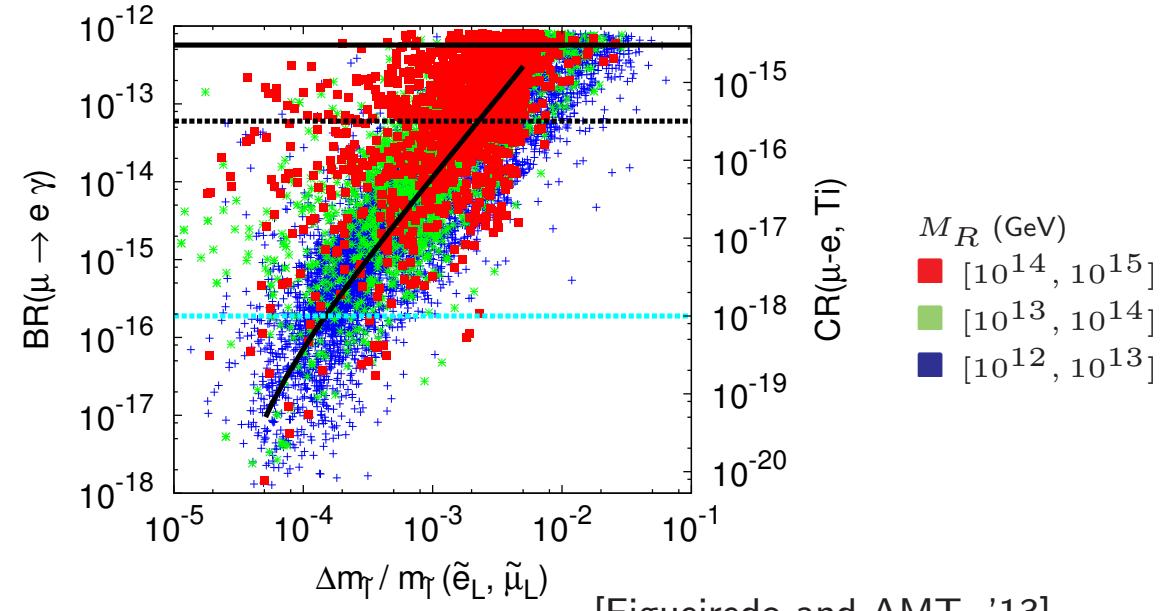
Supersymmetric type I seesaw

- Large Y^ν : sizable contributions to cLFV observables

cLFV driven by the exchange of *virtual SUSY particles*



[Antusch, Arganda, Herrero, AMT, '06 - '07]



[Figueiredo and AMT, '13]

- Y^ν unique source of LFV: synergy of high- and low-energy observables
- Isolated cLFV manifestations \Rightarrow disfavours SUSY seesaw hypothesis
- “Correlated” cLFV observations \Rightarrow strengthen SUSY seesaw hypothesis !

$\frac{\Delta m_{\tilde{\ell}}}{m_{\tilde{\ell}}}(\tilde{e}_L, \tilde{\mu}_L) \gtrsim \mathcal{O}(0.5\%)$ and $\mu \rightarrow e\gamma|_{MEG}$ ✓ !! Hints on the seesaw scale: $M_R \sim 10^{14}$ GeV

Hints of an organising principle: SUSY seesaw and GUTs

★ Supersymmetric Grand Unified Theories

► Reduce arbitrariness of Y^q , Y^ℓ , Y^ν , ...: \Rightarrow increase predictivity and testability!

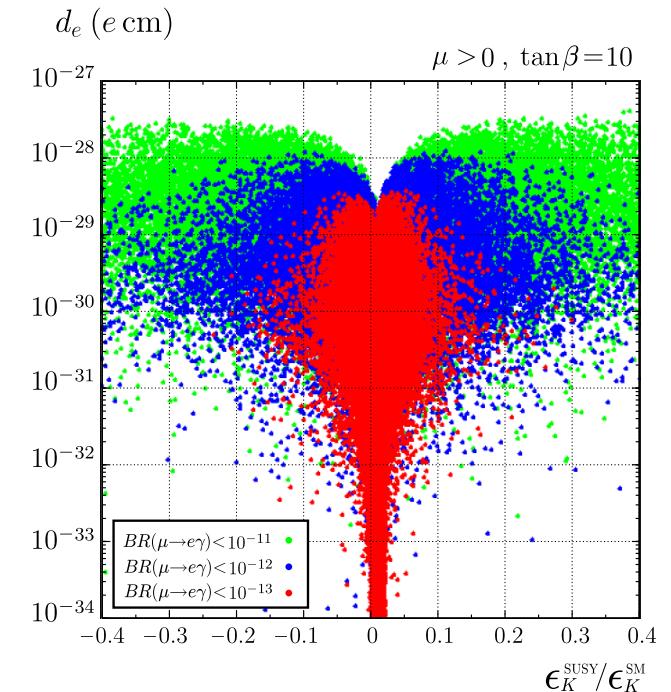
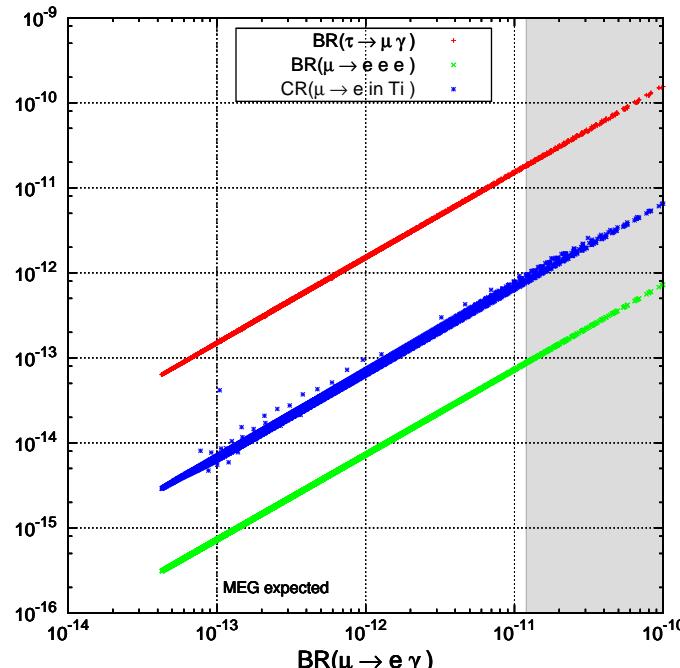
► SU(5) + RH neutrinos **SUSY GUTs**

Correlated CP violation and flavour observables

in lepton and hadron sectors

[Buras et al, 1011.4853]

► SO(10) type II SUSY seesaw



Leptogenesis motivated

highly correlated cLFV observables!

[Calibbi et al, 0910.0337]



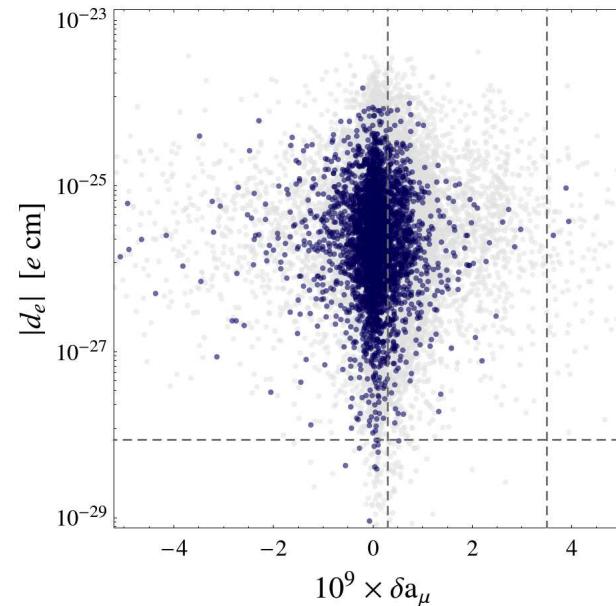
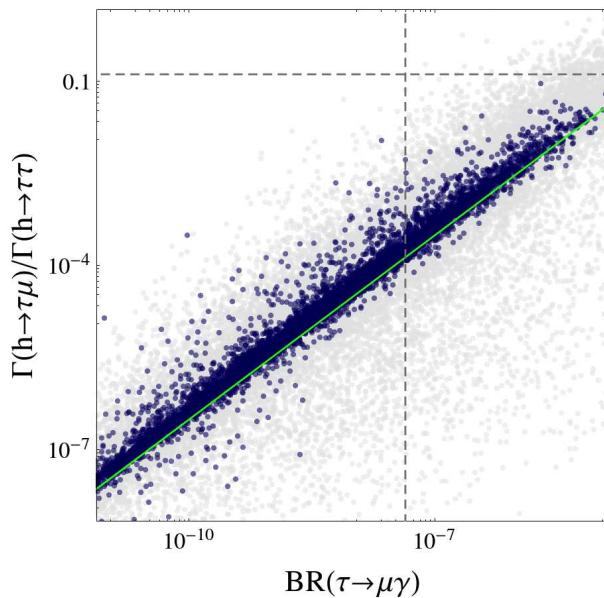


Further possibilities

Vector-like leptons: an example

- ▶ Massive vector-like fermions present in well-motivated SM extensions:
composite Higgs models, warped extra dimensions, ...
- ▶ Global view: generic set-up (composite Higgs inspired), 3 generations of L_i^V and E_i^V
massive neutrinos from additional ν_R and vector-like partners
- ▶ cLFV parametrised by small set of couplings
⇒ correlated observables!

$$\frac{\text{BR}(h \rightarrow \ell_i \ell_j)}{\text{BR}(\ell_i \rightarrow \ell_j \gamma)} \approx \frac{4\pi}{3\alpha} \frac{\text{BR}(h \rightarrow \ell_i \ell_i)|_{\text{SM}}}{\text{BR}(\ell_i \rightarrow \ell_j \nu_i \bar{\nu}_j)}$$



[Falkowski et al, '14]

- ▶ Synergy between FV Higgs decays and cLFV! Flavour conserving EDM and δa_μ as well!

The flavour puzzle: neutrino masses from flavour symmetries

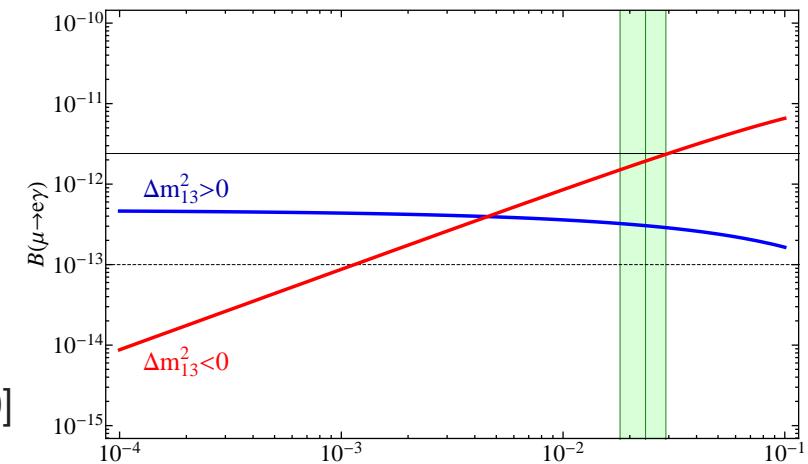
★ Texture of Y from spontaneous or explicit breaking of flavour symmetry G_f

► Continuous flavour symmetries: minimal Abelian case $\Rightarrow G_f = U(1)_{L_e+L_\mu} \times U(1)_{L_\tau}$

$$Y^\nu = \begin{pmatrix} \epsilon_e & ae^{-i\pi/4} & ae^{i\pi/4} \\ \epsilon_\mu & be^{-i\pi/4} & be^{i\pi/4} \\ \epsilon_\tau & \kappa_1 & \kappa_2 \end{pmatrix}$$

$$\text{BR}(\mu \rightarrow e\gamma) \approx 8.0 \times 10^{-4} \times a^2 b^2 \times \left(\frac{\Lambda_{\text{EW}}}{M_R}\right)^4$$

[Deppisch and Pilaftsis, '10]



► Discrete flavour symmetries: $G_f \sim \Delta(3n^2)$ type

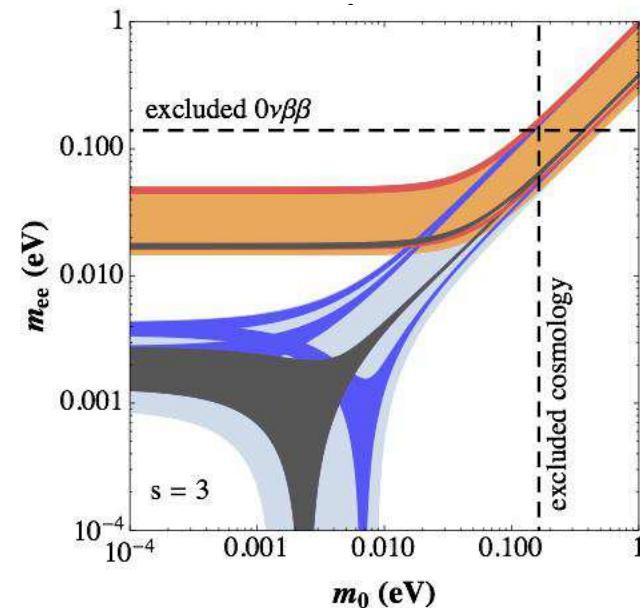
Flavour and CP symmetries \Rightarrow lepton mixings

& low/high energy CP phases

Strong predictions for $0\nu2\beta$ decay m_{ee}

Interplay of low-energy CP phases and BAU

[Hagedorn and Molinaro, '16]





Concluding remarks

Neutrino physics and high-intensity observables

- **Neutrinos** remain a very **open question** in **particle physics, astrophysics and cosmology**
- **Dedicated facilities** will provide crucial data ... but many questions (likely) remain!
- **Confirmed observations** and several “**tensions**” suggest the need to go **beyond the SM**
 - In the **lepton sector**, ν -masses provided the 1st laboratory **evidence of NP**
 - Many experimental “**tensions**” nested in **lepton-related observables**

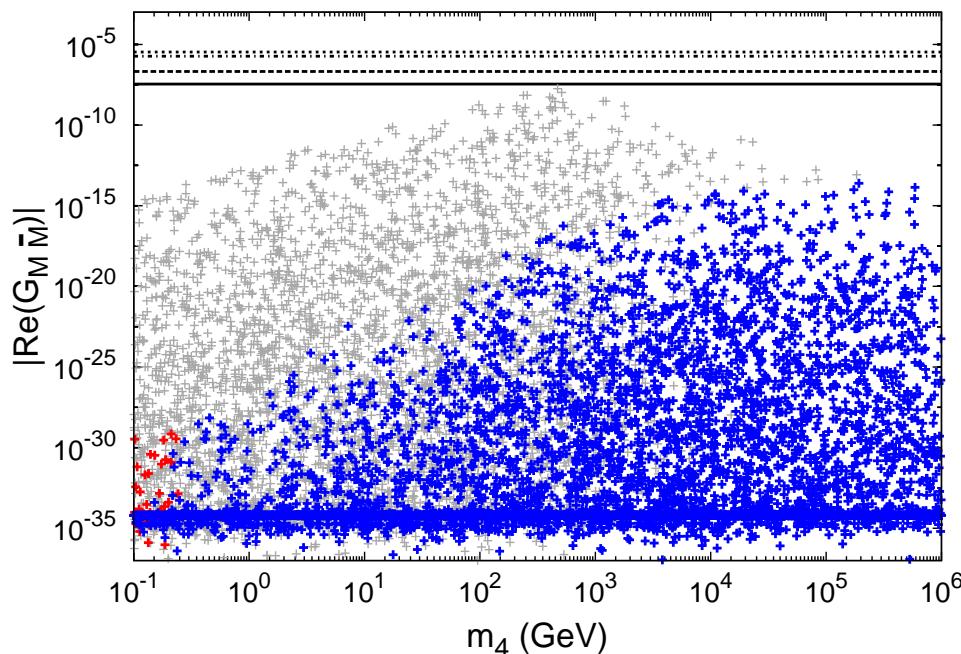
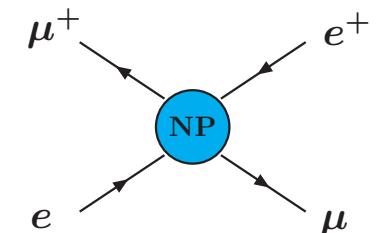


- Very brief overview of a **subset of observables**
 - [other observables: muonium, LFUV, in-flight conversion..., $\mu N \rightarrow e N \xrightarrow{H} \tau X$, ...]
[↪ M. Yamanaka]
- and **sub-sub set of New Physics models** aiming at **accounting for ν phenomena!**
- Lepton physics might offer valuable hints in **constructing and probing NP models**
 - New Physics can be manifest via **cLFV, EDMs, LNV, ... before direct discovery!**
 - High-intensity** data can provide information on the underlying NP model

 **Backup**

Sterile neutrinos: Muonium cLFV

- **Muonium:** hydrogen-like **Coulomb bound state** ($e^- \mu^+$); free of hadronic interactions!
- **Mu – $\bar{\text{Mu}}$ conversion**
Spontaneous conversion of a ($e^- \mu^+$) into ($e^+ \mu^-$)
- Also consider **cLFV Mu decay:** $\text{Mu} \rightarrow e^+ e^-$



- Large values of $G_{\text{M}\bar{\text{M}}}$ precluded due to conflict with $\text{CR}(\mu - e, \text{Au})$ and $\text{BR}(\mu \rightarrow 3e)$
Within **FNAL experimental reach??**
- Maximal values $\text{Mu} \rightarrow e^+ e^- \sim \mathcal{O}(10^{-25})$
Within experimental reach ?

"3+1" toy model [Abada, De Romeri and AMT, '15]

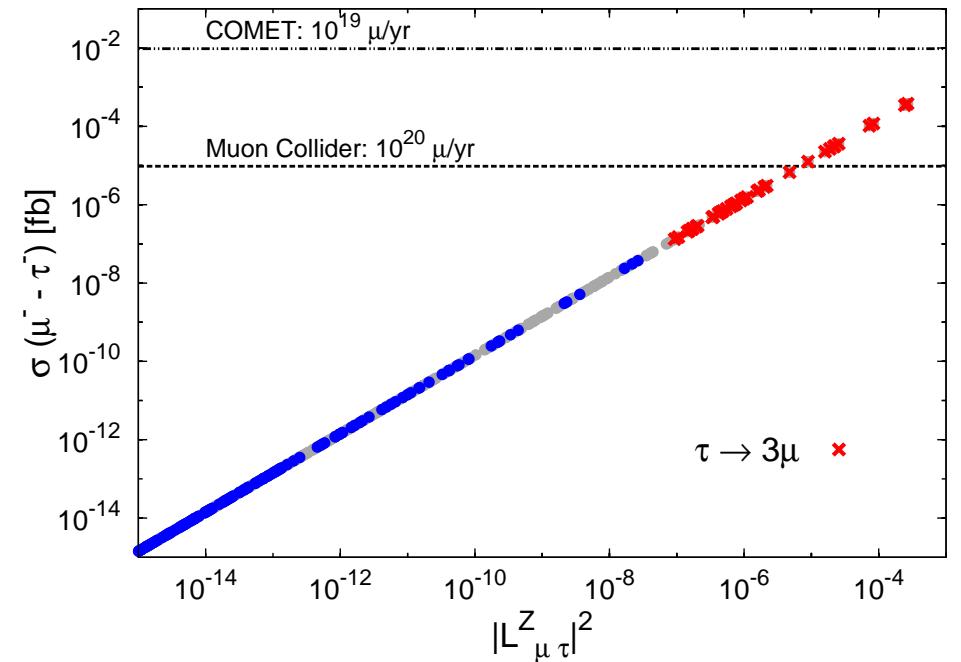
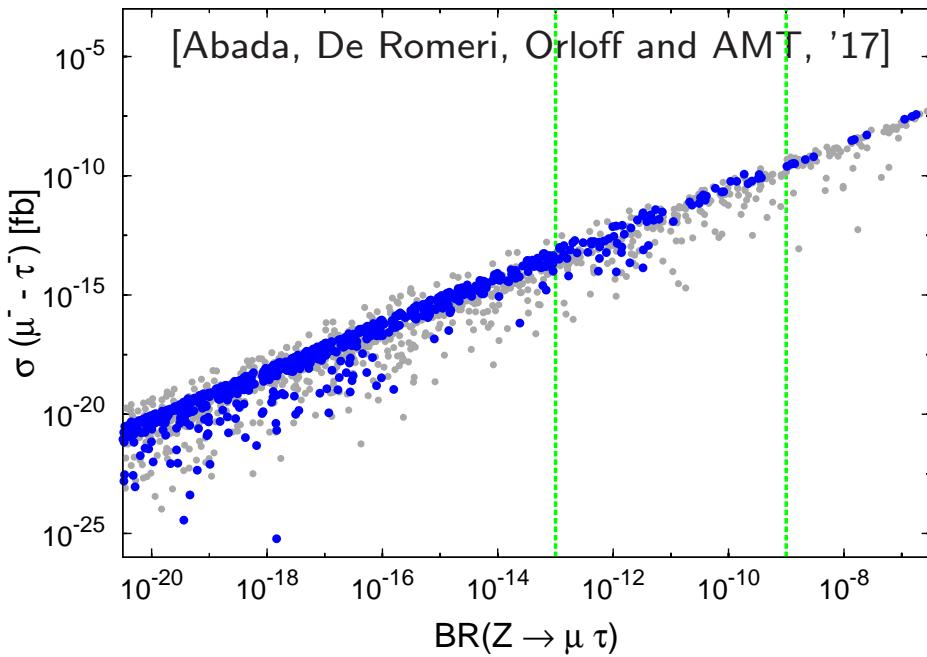
cLFV in-flight conversion

- Energetic beam of leptons (e, μ) directed on fixed (dense) target

$$e + N \rightarrow \mu + N, e + N \rightarrow \tau + N \text{ and } \mu + N \rightarrow \tau + N$$

$$N_{\text{signal}}(\ell_i \rightarrow \ell_j) = N_{\ell_i} \times \sigma(\ell_i \rightarrow \ell_j) \times T_m \times N_{p+n} \times \text{BR}(\tau \rightarrow \mu\nu\nu)$$

- $\sigma(\ell_i \rightarrow \ell_j)$: elastic interactions with nuclei, Z -penguin dominated cLFV



- Large values of $\ell_i \rightarrow \ell_j$ precluded due to conflict with $\text{CR}(\mu - e, \text{Au})$ and $\text{BR}(\ell_j \rightarrow 3\ell_i)$
- N_{signal} beyond experimental sensitivity - even for very intense lepton beams

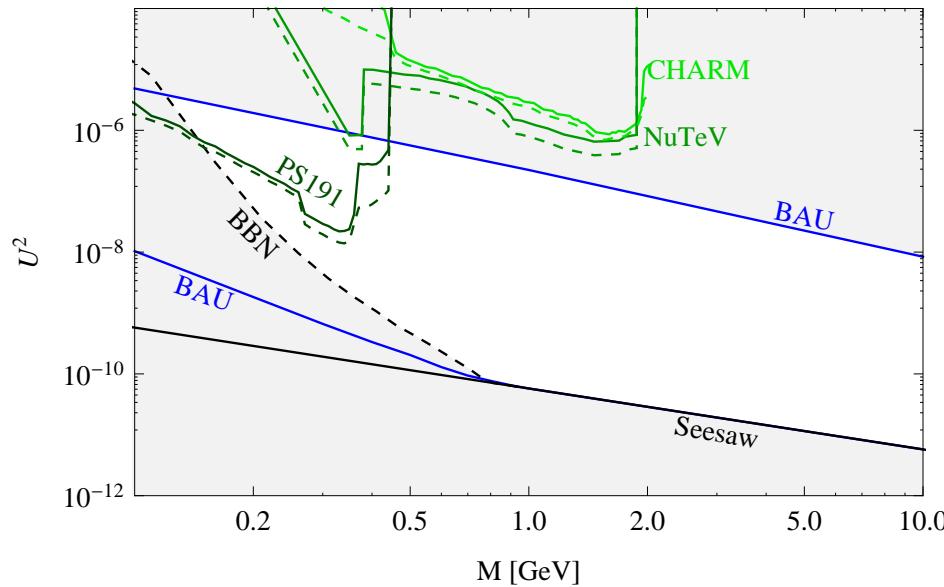
cLFV and ν_s : ν MSM

- Minimal “type I seesaw-like” extension: SM + 3 ν_R

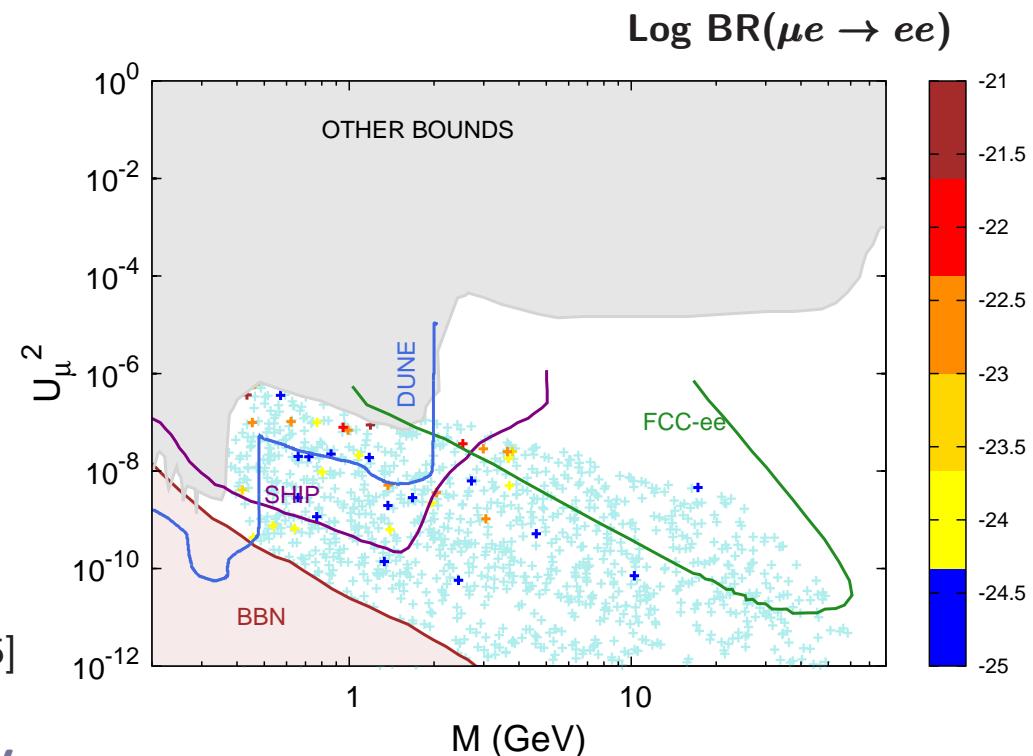
New states account for m_ν^{light} , offer DM candidate, allow for BAU via leptogenesis

⇒ tiny Yukawa couplings; heavily constrained parameter space (th, cosmo, exp..)

[Canetti et al, '13]



[Abada et al, '15]



- ν MSM: very difficult prospects for cLFV

Hints of an organising principle: ν in Left-Right models

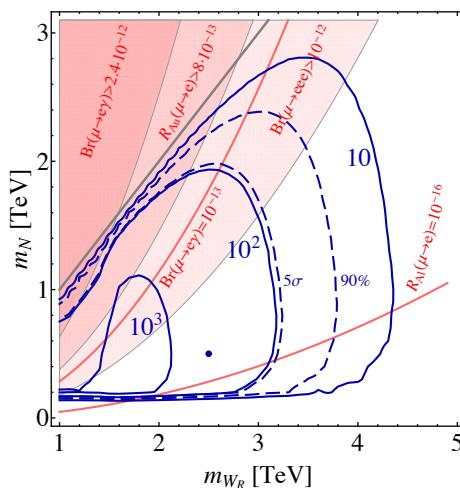
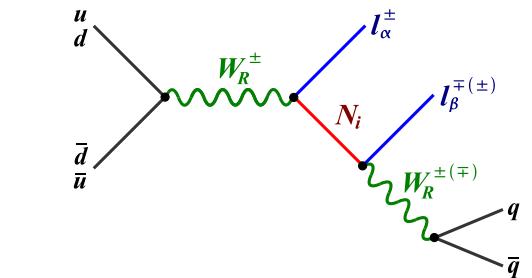
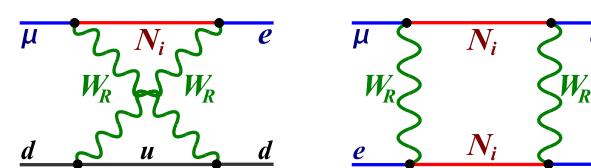
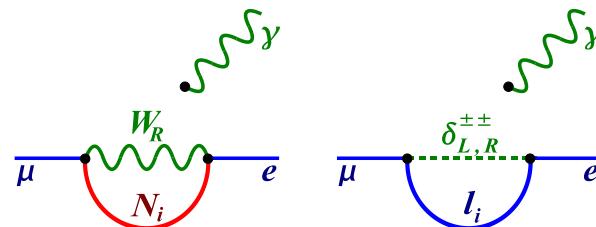
★ Minimal Left-Right extension of the SM (non-SUSY)

► extend SM gauge group: $SU(2)_L \otimes U(1) \Rightarrow SU(2)_L \otimes SU(2)_R \otimes U(1)_{B-L}$

► RH neutrinos automatically included $M_\nu \approx \begin{pmatrix} y_M v_L & y_D m_{EW} \\ y_D^T m_{EW} & y_M v_R \end{pmatrix}$

bi-doublet and triplet Higgs; new Z_R , W_R bosons

► New contributions to cLFV observables at low- and high-energies



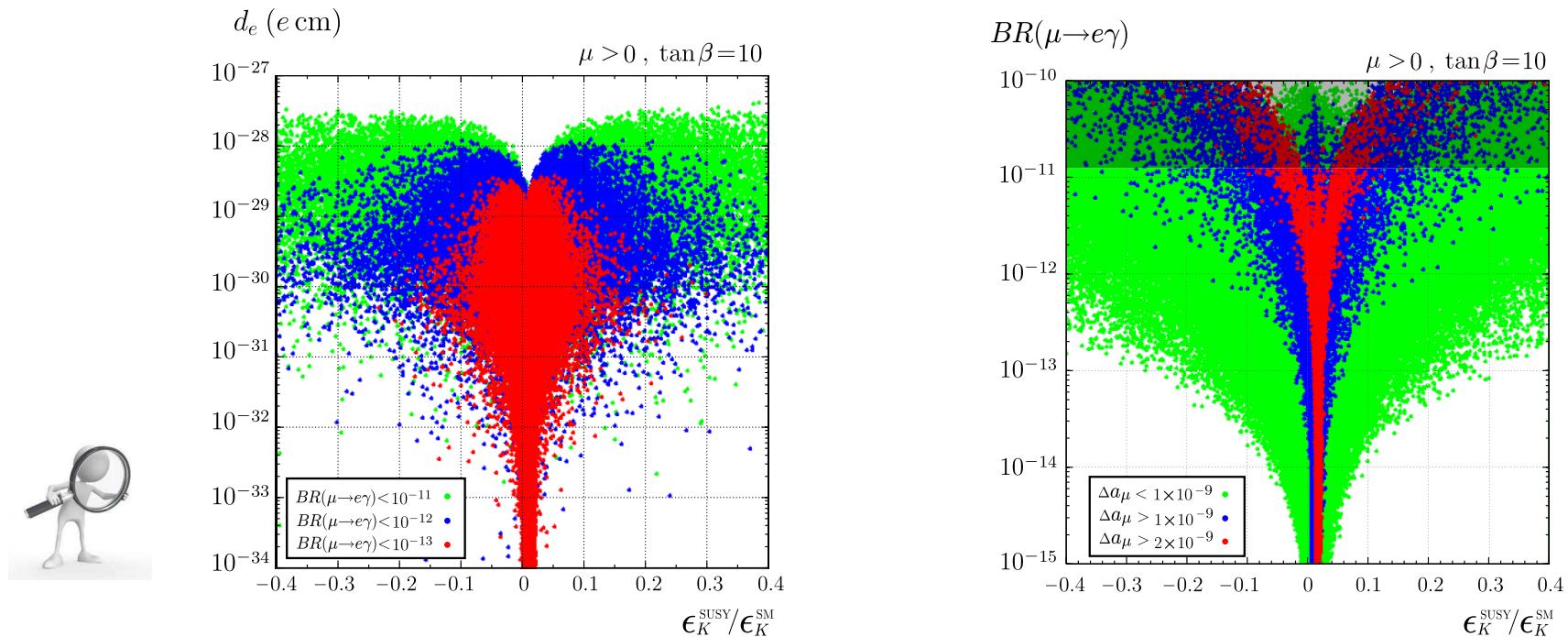
- If LHC \sqrt{s} above heavy neutrino threshold:
dilepton LFV signatures $pp \rightarrow W_R \rightarrow e^\pm \mu^\mp + 2 \text{ jets}$
- Complementarity studies of LHC signatures and
low-energy rare decays

[Das et al, 1206.0656]

Hints of an organising principle: SUSY seesaw and GUTs

★ Supersymmetric Grand Unified Theories - “Type I”

- ▶ Reduce arbitrariness of Yukawa couplings $Y^q, Y^\ell, Y^\nu \dots$



- ▶ SU(5) + RH neutrinos SUSY GUTs
- ▶ Correlated CP violation and flavour violating observables
in lepton and hadron sectors!

[Buras et al, 1011.4853]

► Model-independent approach:
New Physics and low-energy observables

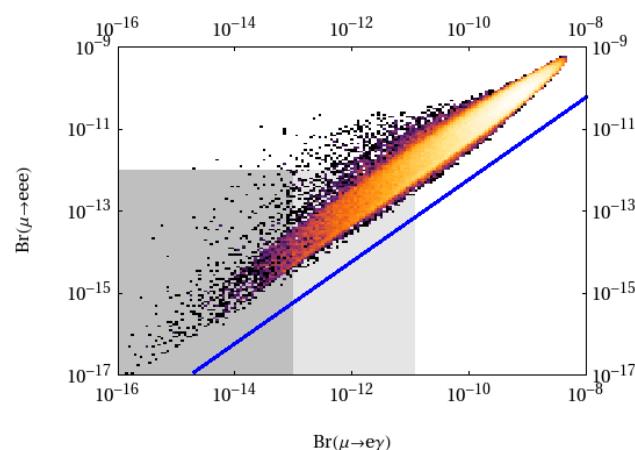
Models of New Physics: some more examples

► “Geometric” flavour violation - extra dimensional Randall-Sundrum models

$e - \mu$ bounds constrain NP scale beyond LHC reach: $T_{KK} \gtrsim 4$ TeV ($\sim KK^{(1st)} \gtrsim 10$ TeV)
future sensitivities: exclude (general) anarchic RS models up to 8 TeV ($\sim m_{KK-g} \gtrsim 20$ TeV)

[Beneke et al, 1508.01705]

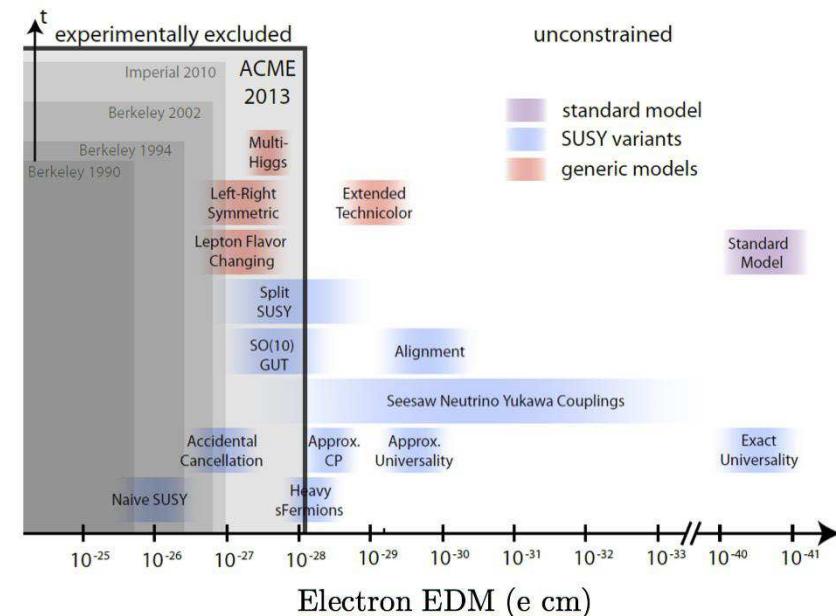
► cLFV and compositeness - Little(st) Higgs



distinctive patterns for ratios of observables (testability!)

$BR(\mu \rightarrow e\gamma)$ - disfavour important regions in parameter space!

[Blanke et al, '09]



► And more observables to test them!

[From A. West, PIC2015]

Effective approach

- \mathcal{L}^{eff} - “vestigial” (new) interactions of “heavy” fields with **SM** at low-energies

$$\mathcal{L}^{\text{eff}} = \mathcal{L}^{\text{SM}} + \sum_{n \geq 5} \frac{1}{\Lambda^{n-4}} \mathcal{C}^n(g, Y, \dots) \mathcal{O}^n(\ell, q, H, \gamma, \dots)$$

- **Dimension 5** - $\Delta\mathcal{L}^5$ (Weinberg): neutrino masses (**LNV**, $\Delta L = 2$)

a unique operator $\mathcal{O}_{ij}^5 \sim (L_i H)(H L_j)$

- **Dimension 6** - $\Delta\mathcal{L}^6$: kinetic corrections, **cLFV**, EWP tests, **EDMs**, t physics...

some examples (**dipole and 3-body**)

Dipole: $\mathcal{O}_{\ell_i \ell_j \gamma}^6 \sim L_i \sigma^{\mu\nu} e_j H F_{\mu\nu}$ radiative decays $\ell_i \rightarrow \ell_j \gamma$, χ -flipping ℓ^\pm dipole moments, ...

4 fermion: $\mathcal{O}_{\ell_i \ell_j \ell_k \ell_l}^6 \sim (\ell_i \gamma_\mu P_{L,R} \ell_j)(\ell_k \gamma^\mu P_{L,R} \ell_l)$ 3-body decays $\ell_i \rightarrow \ell_j \ell_k \ell_l$, ...

$\mathcal{O}_{\ell_i \ell_j q_k q_l}^6 \sim (\ell_i \gamma_\mu P_{L,R} \ell_j)(q_k \gamma^\mu P_{L,R} q_l)$ $\mu - e$ in Nuclei, meson decays, ...

Vector/scalar: $\mathcal{O}_{H H \ell_i \ell_j}^6 \sim (H^\dagger i \overleftrightarrow{D}_\mu H)(\ell_i \gamma_\mu \ell_j)$ 3-body decays $\ell_i \rightarrow \ell_j \ell_k \ell_l$, ...

[\leadsto presentation by G.M. Pruna, ...]

- **Higher order** - $\Delta\mathcal{L}^{7,8,\dots}$: $0\nu2\beta$, ν (transitional) magnetic moments, NSI, unitarity violation...

Constraining \mathcal{L}^{eff} : cLFV example



- Apply **experimental** bounds on (e.g.) **cLFV observables** to **constrain** $\frac{\mathcal{C}_{ij}^6}{\Lambda^2}$

Hypotheses on:

1. **size of “new couplings”**

⇒ **Natural couplings**

$$\mathcal{C}_{ij}^6 \sim \mathcal{O}(1)$$

2. **scale of “new physics”**

⇒ **Natural scale** - delicate..

direct discovery $\Lambda \sim \text{TeV}$

Effective coupling (example)	Bounds on Λ (TeV) (for $ \mathcal{C}_{ij}^6 = 1$)	Bounds on $ \mathcal{C}_{ij}^6 $ (for $\Lambda = 1$ TeV)	Observable
$\mathcal{C}_{e\gamma}^{\mu e}$	6.3×10^4	2.5×10^{-10}	$\mu \rightarrow e\gamma$
$\mathcal{C}_{e\gamma}^{\tau e}$	6.5×10^2	2.4×10^{-6}	$\tau \rightarrow e\gamma$
$\mathcal{C}_{e\gamma}^{\tau\mu}$	6.1×10^2	2.7×10^{-6}	$\tau \rightarrow \mu\gamma$
$\mathcal{C}_{\ell\ell,ee}^{\mu eee}$	207	2.3×10^{-5}	$\mu \rightarrow 3e$
$\mathcal{C}_{\ell\ell,ee}^{e\tau ee}$	10.4	9.2×10^{-5}	$\tau \rightarrow 3e$
$\mathcal{C}_{\ell\ell,ee}^{\mu\tau\mu\mu}$	11.3	7.8×10^{-5}	$\tau \rightarrow 3\mu$
$\mathcal{C}_{(1,3)H\ell}^{\mu e}, \mathcal{C}_{He}^{\mu e}$	160	4×10^{-5}	$\mu \rightarrow 3e$
$\mathcal{C}_{(1,3)H\ell}^{\tau e}, \mathcal{C}_{He}^{\tau e}$	≈ 8	1.5×10^{-2}	$\tau \rightarrow 3e$
$\mathcal{C}_{(1,3)H\ell}^{\tau\mu}, \mathcal{C}_{He}^{\tau\mu}$	≈ 9	$\approx 10^{-2}$	$\tau \rightarrow 3\mu$

[Feruglio et al, 2015]

- Despite its generality, caution in interpreting these effective limits!

- limits assume **dominance of one operator**; NP leads to several (interference...)

- contributions from **higher order operators** may be non-negligible if Λ is low...

- **multiple “new physics” scales**: $\mathcal{L}^{\text{eff}} = \mathcal{L}^{\text{SM}} + \frac{1}{\Lambda_{\text{LNV}}} \mathcal{C}^5(m_\nu) + \frac{1}{\Lambda_{\text{LFV}}^2} \mathcal{C}^6(\ell_i \leftrightarrow \ell_j) + \dots$

- Can these limits be used to extract information about $\mathcal{C}^5(m_\nu)$??