

Connections between low-energy CP violation, lepton number violating collider signals and genesis mechanisms

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CP³ Origins

Cosmology & Particle Physics

SDU 

Overview

- flavor symmetry with CP for low-energy CP phases δ, α, β
- correlation between δ, α, β and baryon asymmetry Y_B
- possible connection to lepton number violating signals at LHC

Low-energy CP phases in the lepton sector

Parametrization

$$U_{PMNS} = \tilde{U} \text{diag}(1, e^{i\alpha/2}, e^{i(\beta/2+\delta)})$$

with

$$\tilde{U} = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix}$$

and $s_{ij} = \sin \theta_{ij}$, $c_{ij} = \cos \theta_{ij}$

Experimental status

- first indications for Dirac phase $\delta \sim \frac{3}{2} \pi$
(Capozzi et al. ('16), Esteban et al. ('16), de Salas et al. ('17))
- no constraints on Majorana phases α and β

***Prediction of CP phases with a
flavor and a CP symmetry***

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- framework: 3 copies of charged leptons & Majorana neutrinos
- impose flavor symmetry G_f on space of 3 lepton generations
- G_f is **non-abelian**, **finite** and **discrete**
- impose also CP symmetry

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- G_f and CP must be broken at low energies and we assume **residual symmetries** G_e (G_ν) in the charged lepton (neutrino) sector
 - ↪ form of charged lepton mass matrix m_e & Majorana mass matrix m_ν is restricted by G_e and G_ν

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- G_f and CP must be broken at low energies and we assume **residual symmetries** G_e (G_ν) in the charged lepton (neutrino) sector
 - \hookrightarrow form of charged lepton mass matrix m_e & Majorana mass matrix m_ν is restricted by G_e and G_ν
 - in particular, G_ν contains CP, described by CP transformation X

$$X m_\nu X = m_\nu^*$$

\hookrightarrow contributions U_e & U_ν to lepton mixing U_{PMNS} are fixed

Prediction of CP phases with a flavor and a CP symmetry

(Harrison/Scott ('02), Grimus/Lavoura ('03), Feruglio/H/Ziegler ('12),
Holthausen/Lindner/Schmidt ('12), Chen et al. ('14), Grimus/Rebelo ('95))

G_f and CP

charged leptons

$$G_e$$

$$U_e$$

neutrinos

$$G_\nu = Z_2 \times \text{CP}$$

$$U_\nu = \Omega_\nu R(\theta) K_\nu$$

$$U_{PMNS} = U_e^\dagger \Omega_\nu R(\theta) K_\nu$$

[Masses do not play a role in this approach.]

Prediction of CP phases with a flavor and a CP symmetry

$$U_{PMNS} = U_e^\dagger \Omega_\nu R(\theta) K_\nu$$

- 3 unphysical phases are removed by $U_e \rightarrow U_e K_e$
- U_{PMNS} contains one free parameter θ
- possible permutations of rows and columns of U_{PMNS}

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Predictions

Mixing angles and **all CP phases** are predicted
in terms of one free parameter θ only,
up to permutations of rows/columns

Example

(H/Meroni/Molinaro ('14); see also Ding et al. ('14))

$\Delta(3 n^2), \Delta(6 n^2)$ and **CP**

charged leptons

$$G_e = Z_3$$

$$U_e$$

neutrinos

$$G_\nu = Z_2 \times \mathbf{CP}$$

$$U_\nu = \Omega_\nu R(\theta) K_\nu$$

$$U_{PMNS} = U_e^\dagger \Omega_\nu R(\theta) K_\nu$$

four different types of mixing patterns with different characteristics

Example

Case 3 b.1)

(H/Meroni/Molinaro ('14))

- first column is fixed via choice of residual Z_2 symmetry $Z_2(m)$ in neutrino sector
- in particular, solar mixing angle constrains m to be $m \approx \frac{n}{2}$
- free parameter θ is fixed by reactor mixing angle
- for $m = \frac{n}{2}$ we find lower limit on CP violation via **Dirac phase**

$$|\sin \delta| \gtrsim 0.71$$

and both **Majorana phases** α, β depend on $X = X(s)$ only

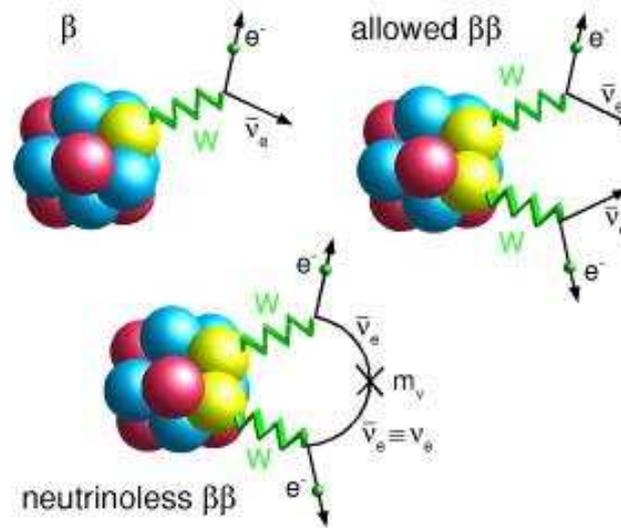
$$|\sin \alpha| = |\sin \beta| = |\sin 6 \phi_s| \quad \text{with} \quad \phi_s = \frac{\pi s}{n} \quad \text{and} \quad s = 0, \dots, n-1$$

Example

Case 3 b.1) and $n = 8, m = 4$: *(H/Meroni/Molinaro ('14))*
some viable choices of s

s	$\sin^2 \theta_{13}$	$\sin^2 \theta_{12}$	$\sin^2 \theta_{23}$	$\sin \delta$	$\sin \alpha = \sin \beta$
$s = 1$	0.0220	0.318	0.579	0.936	$-1/\sqrt{2}$
	0.0220	0.318	0.421	-0.936	$-1/\sqrt{2}$
$s = 2$	0.0216	0.319	0.645	-0.739	1
$s = 4$	0.0220	0.318	0.5	∓ 1	0

Low-energy CP phases in neutrinoless double beta decay



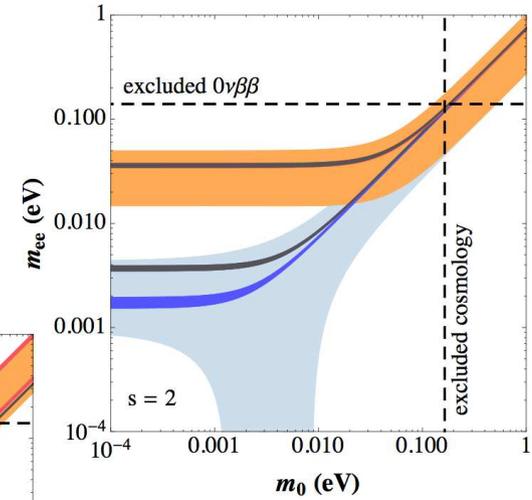
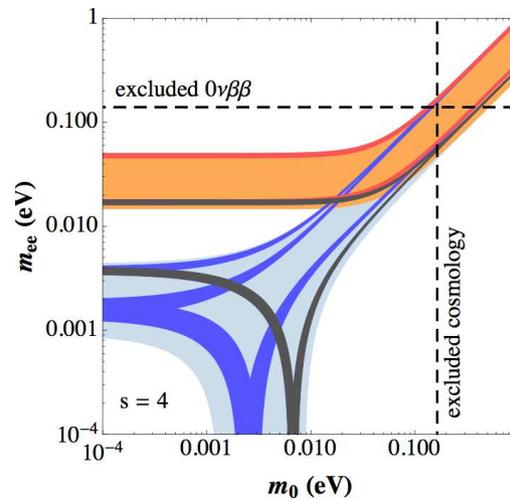
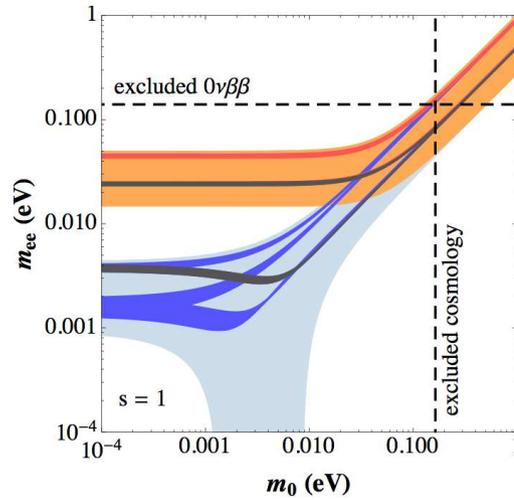
Dependence of neutrinoless double beta decay on **Majorana phases** α , β

$$m_{ee} = \left| \cos^2 \theta_{12} \cos^2 \theta_{13} m_1 + \sin^2 \theta_{12} \cos^2 \theta_{13} e^{i\alpha} m_2 + \sin^2 \theta_{13} e^{i\beta} m_3 \right|$$

Low-energy CP phases in neutrinoless double beta decay

Case 3 b.1) and $n = 8, m = 4$

(H/Molinaro ('16))



Further approaches with CP symmetries

- many more flavor symmetries have been combined with CP (*Di Iura/H/Meloni ('15), Li/Ding ('15), Ballett et al. ('15), Rong ('16), Yao/Ding ('16), ...*)
- study of more general CP symmetries (*Everett et al. ('15)*)
- two CP symmetries in neutrino sector (*Chen et al. ('14)*)
- textures of CP transformations (*Chen et al. ('16)*)
- considerable efforts in building models in recent past (*Altarelli, Antusch, Branco, Chen, de Medeiros Varzielas, Ding, Everett, Feruglio, Gehrlein, Girardi, Gonzalez Felipe, Grimus, H, He, Joaquim, King, Lavoura, Luhn, Mahanthappa, Meloni, Meroni, Mohapatra, Nishi, Päs, Pascoli, Petcov, Rodejohann, Smirnov, Spinrath, Tanimoto, Valle, ...*)

***Correlation between low-energy CP
phases and Y_B***

Correlation between low-energy CP phases and Y_B

- type-I seesaw mechanism with 3 right-handed (RH) neutrinos N
- Lagrangian

$$\mathcal{L}_l = -Y_D \bar{l} H^c N - \frac{1}{2} \overline{N^c} M_R N + \text{h.c.}$$

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- leptogenesis (flavored, unflavored, resonant) can explain the measured baryon asymmetry of the Universe
 $Y_B = (8.65 \pm 0.09) \times 10^{-11}$ (*Planck ('15)*)
through decay of RH neutrinos
- one of the Sakharov conditions is C and CP violation
- rule of thumb

$$Y_B \sim 10^{-3} \epsilon \eta \quad \text{with } \epsilon \text{ CP asymmetry, } \eta \text{ efficiency factor}$$

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- three possibilities to realize G_ν
 - a) M_R invariant under G_ν , Y_D invariant under G_f and CP
 - b) M_R invariant under G_f and CP, Y_D invariant under G_ν
 - c) M_R and Y_D invariant under G_ν

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(H/Molinaro ('16))

- scenario with unflavored leptogenesis
- CP phases are constrained by G_ν
- light and heavy neutrino masses are strongly correlated

$$m_i \propto \frac{1}{M_i}$$

- mixing among RH neutrinos gives PMNS mixing matrix

$$U_{PMNS} = U_\nu = U_R = \Omega_\nu R(\theta) K_\nu$$

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- crucial ingredient: **corrections $\mathcal{O}(\kappa)$ to Y_D**

(Jenkins/Manohar ('08), Bertuzzo et al. ('09), H/Molinaro/Petcov ('09), Aristizabal Sierra et al. ('09))

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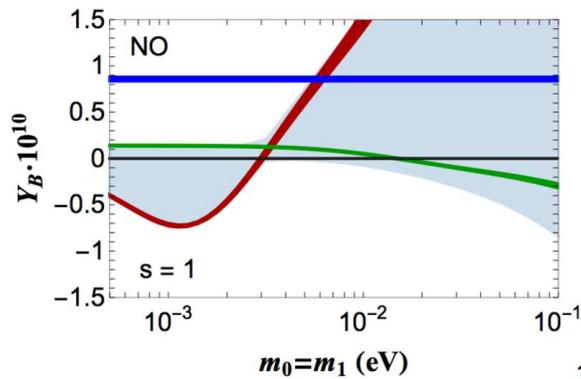
(H/Molinaro ('16))

- scenario with unflavored leptogenesis
- CP phases are constrained by G_ν
- crucial ingredient: **corrections $\mathcal{O}(\kappa)$ to Y_D**
- then: size of ϵ is small, since $\epsilon \propto \kappa^2$,
sign of ϵ can be fixed, since CP phases are constrained

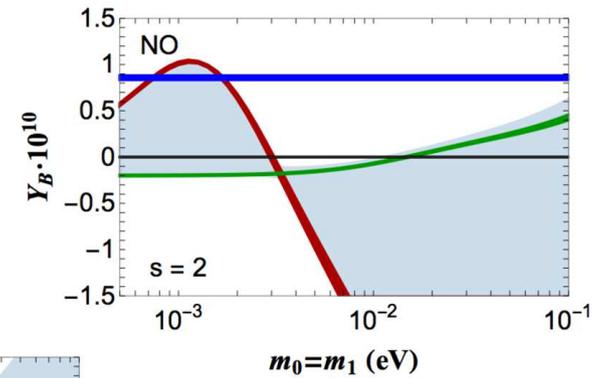
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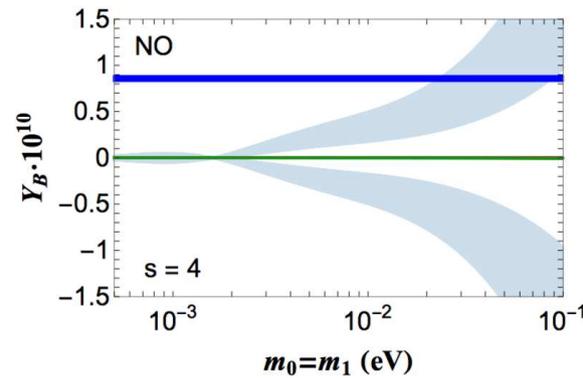
(H/Molinaro ('16))



only $\sin \delta$
non-zero



$\sin \alpha < 0$



$\sin \alpha > 0$

for light neutrinos with normal ordering

Correlation between low-energy CP phases and Y_B

b) M_R invariant under G_f and CP, Y_D invariant under G_ν

(Dev/H/Molinaro (in preparation))

- scenario suitable for resonant leptogenesis, since RH neutrino masses are degenerate
- CP phases are constrained by G_ν
- light neutrino masses depend on Yukawa couplings

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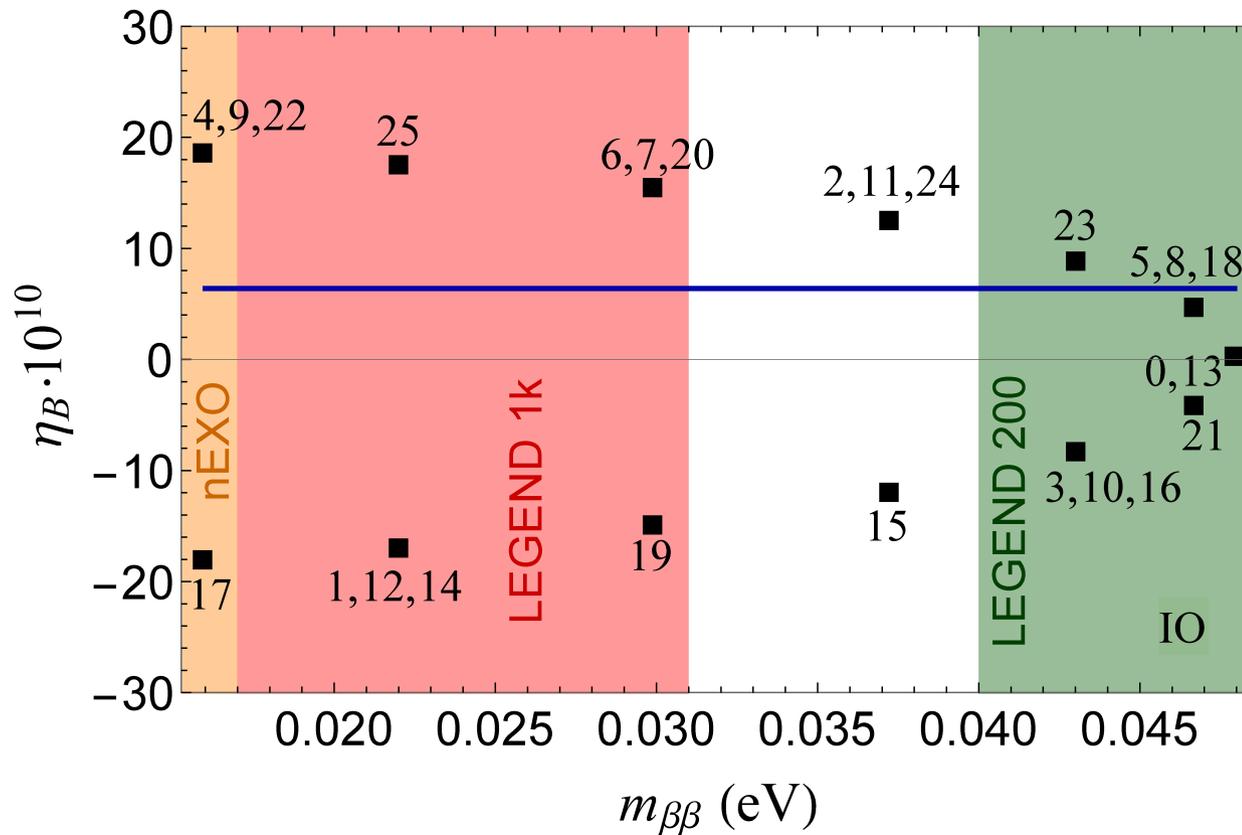
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- CP phases are constrained by G_ν
- light neutrino masses depend on Yukawa couplings
- crucial ingredient: **corrections $\mathcal{O}(\kappa)$ to M_R**
in order to split RH neutrino masses
for non-vanishing CP asymmetry

Example

Case 1) and $n = 26$

(Dev/H/Molinaro (in preparation))



for light neutrinos with inverted ordering and $m_0 = m_3 = 0$

Correlation between low-energy CP phases and Y_B

c) M_R and Y_D invariant under G_ν

- suitable scenario for flavored leptogenesis
- examples can be found in literature
(Mohapatra/Nishi ('15), Chen et al. ('16), Yao/Ding ('16))
- in such scenarios corrections of $\mathcal{O}(\kappa)$ are not needed for non-vanishing CP asymmetries
- however, correlation between low-energy CP phases and Y_B is not as strong as in other scenarios

***Possible connection to lepton
number violating signals at LHC***

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We consider

b) M_R invariant under G_f and CP, Y_D invariant under G_ν

(Dev/H/Molinaro (in preparation))

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- RH neutrino masses M_N can be in range of few 100 GeV

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production at LHC is too small, since it depends on small Yukawa couplings $y \lesssim 10^{-6}$

Possible connection to lepton number violating signals at LHC

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(Dev/H/Molinaro (in preparation))

- RH neutrino masses M_N can be in range of few 100 GeV
- problem:
production at LHC is too small, since it depends on small Yukawa couplings $y \lesssim 10^{-6}$
- solution:
extension by $U(1)$ gauge symmetry, e.g. $B - L$ symmetry
then production via Z' gauge boson becomes possible

$$pp \rightarrow Z' \rightarrow N_i N_j$$

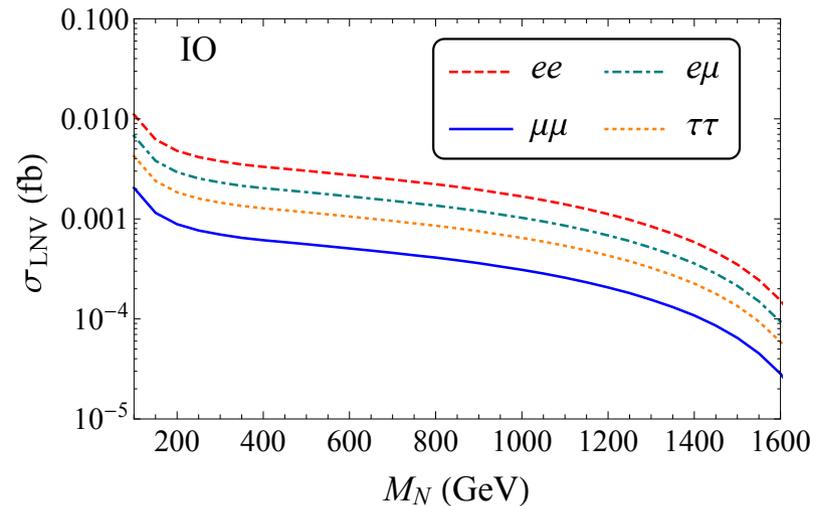
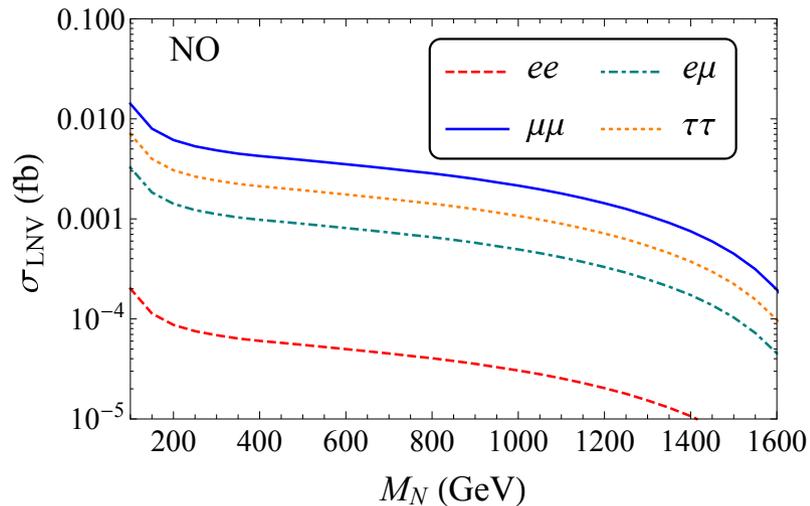
*(e.g. Deppisch/Desai/Valle ('13), Helo/Hirsch/Kovalenko ('13), Blanchet et al. ('09),
Blanchet/Dev/Mohapatra ('10))*

Possible connection to lepton number violating signals at LHC

(Dev/H/Molinaro (in preparation))

Study of lepton number violating process

$$pp \rightarrow Z' \rightarrow N_i N_i \rightarrow \ell_\alpha^\pm \ell_\beta^\pm + 4j$$



ratio of $\sigma_{\text{LNV}}^{\alpha\beta}$ can distinguish between normal and inverted ordering

Conclusions

- flavor and CP symmetry can constrain low-energy CP phases δ, α, β
- in scenarios with RH neutrinos we can correlate δ, α, β with the sign of the baryon asymmetry Y_B of the Universe
- connection to lepton number violating signals at LHC exists

Thank you for your attention.