

Search for NSI in neutrino propagation with IceCube DeepCore

Thomas Ehrhardt for the IceCube Collaboration

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- ▶ if SM is effective low-energy theory, with new physics at high energy scale Λ :

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \frac{O_5}{\Lambda^1} + \frac{O_6}{\Lambda^2} + \dots + \text{h.c.}$$

- ▶ if SM is effective low-energy theory, with new physics at high energy scale Λ :

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \frac{O_5}{\Lambda^1} + \frac{O_6}{\Lambda^2} + \dots + \text{h.c.}$$

after EWSB: Majorana
neutrino mass

Weinberg (1979)

- ▶ if SM is effective low-energy theory, with new physics at high energy scale Λ :

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \frac{O_5}{\Lambda^1} + \boxed{\frac{O_6}{\Lambda^2} + \dots} + \text{h.c.}$$

- ▶ $d \geq 6$ operators can give rise to **general neutrino interactions**
- ▶ strongest low-energy effects typically expected for $d = 6$ Bischer, Rodejohann (2019)
 - ▶ 10 possible operators for both **NC** and **CC** interactions; 2 usually considered as NSI



- ▶ neutrino propagation only sensitive to vector-type interactions

Bergmann, Grossmann, Nardi (1999)

Why consider NC (propagation) NSI?

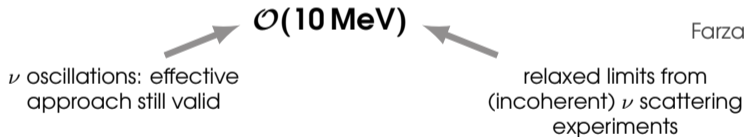
▶ **in principle:**

gauge invariance
↓
strong constraints on NC NSI from non-observation of charged lepton flavour violation

Gavela, Hernandez, Ota, Winter (2009),
Antusch, Baumann, Fernández-Martínez (2009)

▶ **but:**

models with neutral light new mediators able to avoid these bounds



Farzan, Tórtola (2018)

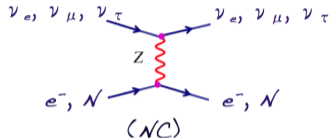
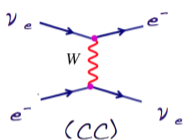
▶ **also:**

NC NSI from **CC NSI** via Fierz transformation

↑
difficult to obtain large NSI (ν production, detection)

Bischer, Rodejohann, Xu (2018)

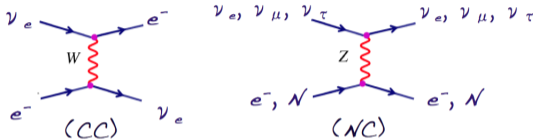
- flavour-asymmetric matter \Rightarrow flavour-dependent index of refraction in **coherent neutrino scattering**



interaction Hamiltonian:

$$H_I = H_Z^n + H_Z^p + H_Z^e + H_W^e$$

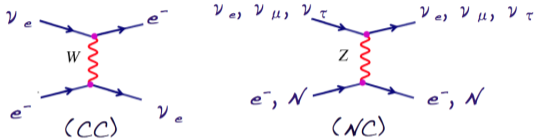
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interaction Hamiltonian:

$$H_I = H_Z^n + \underbrace{H_Z^p + H_Z^e}_{=0} + H_W^e$$

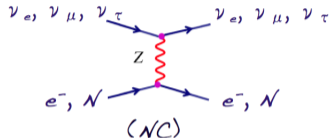
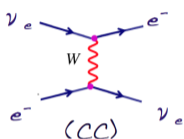
- flavour-asymmetric matter \Rightarrow flavour-dependent index of refraction in **coherent neutrino scattering**



interaction Hamiltonian:

$$H_I = \underbrace{H_Z^n}_{\text{unobservable}} + H_W^e$$

- flavour-asymmetric matter \Rightarrow flavour-dependent index of refraction in **coherent neutrino scattering**

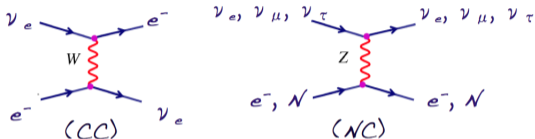


interaction Hamiltonian:

$$H_I = H_W^e = \text{diag}(\sqrt{2}G_F N_e, 0, 0)$$

(responsible for MSW & parametric enhancement)

- flavour-asymmetric matter \Rightarrow flavour-dependent index of refraction in **coherent neutrino scattering**



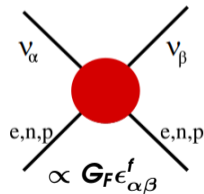
interaction Hamiltonian:

$$H_I = H_W^e = \text{diag}(\sqrt{2}G_F N_e, 0, 0)$$

(responsible for MSW & parametric enhancement)

- NSI introduce additional potential $H_{\text{NSI}}(\mathbf{x}) = \sqrt{2}G_F \sum_{f=e,n,p} N_f(\mathbf{x}) \epsilon_{\alpha\beta}^f$

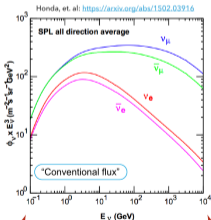
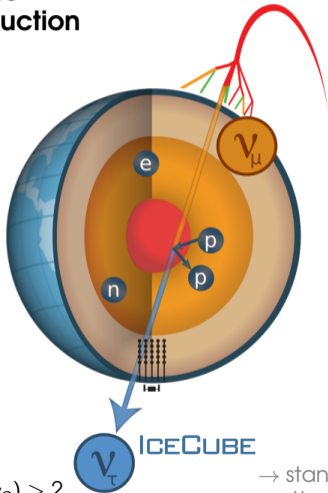
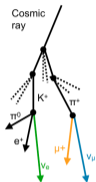
$$H_{\text{mat}}(\mathbf{x}) = \sqrt{2}G_F N_e(\mathbf{x}) \begin{pmatrix} 1 + (\epsilon_{ee}^{\oplus} - \epsilon_{\mu\mu}^{\oplus})(\mathbf{x}) & \epsilon_{e\mu}^{\oplus}(\mathbf{x}) & \epsilon_{e\tau}^{\oplus}(\mathbf{x}) \\ \epsilon_{e\mu}^{\oplus*}(\mathbf{x}) & 0 & \epsilon_{\mu\tau}^{\oplus}(\mathbf{x}) \\ \epsilon_{e\tau}^{\oplus*}(\mathbf{x}) & \epsilon_{\mu\tau}^{\oplus*}(\mathbf{x}) & (\epsilon_{\tau\tau}^{\oplus} - \epsilon_{\mu\mu}^{\oplus})(\mathbf{x}) \end{pmatrix}$$



- in Earth, $\epsilon_{\alpha\beta}^{\oplus}$ can be taken as constant:

$$\epsilon_{\alpha\beta}^{\oplus}(\mathbf{x}) \approx \epsilon_{\alpha\beta}^{\oplus} = \epsilon_{\alpha\beta}^e + \epsilon_{\alpha\beta}^p + 1.051 \epsilon_{\alpha\beta}^n$$

1. atmospheric neutrino production



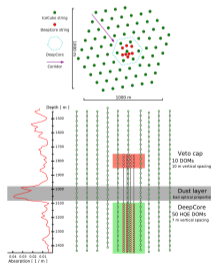
$$\Phi(\nu_\mu)/\Phi(\nu_e) = 2 \quad \Phi(\nu_\mu)/\Phi(\nu_e) > 2$$

2. neutrino oscillation

$$i \frac{d}{dx} |\nu\rangle = \left(U \frac{M_D^2}{2E} U^\dagger + H_{\text{mat}}(x) \right) |\nu\rangle$$

baselines
between 20 km
and 12 700 km

3. neutrino detection



- ▶ 8 dedicated densely instrumented DeepCore strings

- ▶ sensitivity to $\mathcal{O}(10 \text{ GeV})$ $\nu + \bar{\nu}$

- ▶ rich physics potential

→ $\nu_\mu + \bar{\nu}_\mu$ disappearance

→ $\nu_\tau + \bar{\nu}_\tau$ appearance

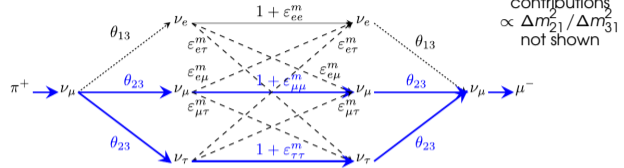
→ decoherence

→ standard & nonstandard matter effects

→ sterile neutrinos

NSI in atmospheric neutrino oscillations

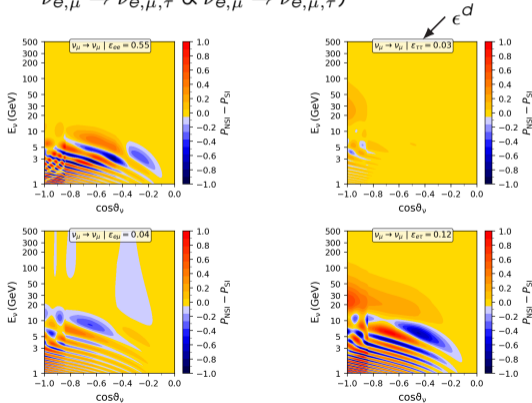
- ▶ here: effects of NSI couplings on ν_μ survival probability
- ▶ (observed: superposition of $\nu_{e,\mu} \rightarrow \nu_{e,\mu,\tau}$ & $\bar{\nu}_{e,\mu} \rightarrow \bar{\nu}_{e,\mu,\tau}$)



Kopp, Lindner, Ota, Sato (2008)

flavour-diagonal couplings

(note the different NSI coupling magnitudes)



flavour-changing couplings

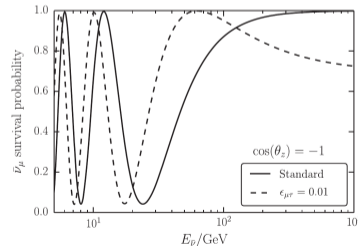
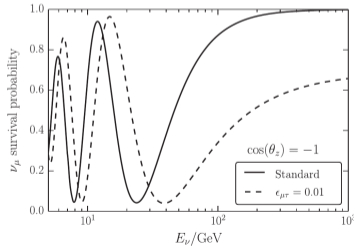
First low-energy DeepCore NSI search: physics

Esmaili, Smirnov (2013)

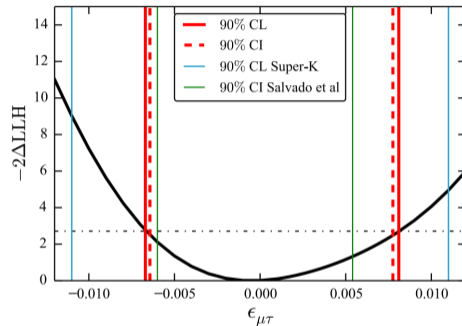
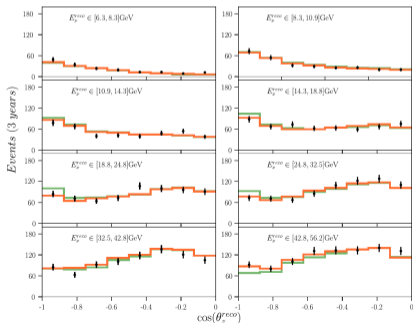
► assumption:

$$H_{\text{mat}}(x) = \sqrt{2}G_F N_{\theta}(x) \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 3\epsilon_{\mu\tau}^d \\ 0 & 3\epsilon_{\mu\tau}^d & 0 \end{pmatrix}$$

- 2ν system of $\nu_{\mu} \leftrightarrow \nu_{\tau}$ approximately realised for $\epsilon_{\theta\beta} \ll 1$, $E_{\nu} \gtrsim 20$ GeV
- solely $\epsilon_{\mu\tau}$ determines high-energy behaviour of 2ν system
- assuming NSI universality ($\epsilon_{\tau\tau} = \epsilon_{\mu\mu}$)
 \Rightarrow strongest constraints on $\epsilon_{\mu\tau}$



- ▶ 3-year sample of $\sim 5\text{k}$ upgoing events
- ▶ high purity of $\nu_\mu + \bar{\nu}_\mu$ CC: $\sim 70\%$
- ▶ atm. μ background $< 1\%$ level



@ 90% C.L. (Wilks):

$$-6.7 \times 10^{-3} < \epsilon_{\mu\tau}^d < 8.1 \times 10^{-3}$$

$$\Leftrightarrow -2.1 \times 10^{-2} < \epsilon_{\mu\tau}^\oplus < 2.4 \times 10^{-2}$$

event selection: medium-statistics osc.
sample extended to $E_{\text{reco,max}} = 100 \text{ GeV}$

data taking period: three years,
2012–2015 \Rightarrow 45k atm. ν 's + $\bar{\nu}$'s

cf. PRD99 032007 (2019)

► 5 **one-at-a-time** (model-dependent) searches:

- 2 searches for **non-universal** interactions: $\epsilon_{\theta\theta}^{\oplus} - \epsilon_{\mu\mu}^{\oplus}$ and $\epsilon_{\tau\tau}^{\oplus} - \epsilon_{\mu\mu}^{\oplus}$
- 3 searches for **flavour-changing** interactions:

$$\epsilon_{\theta\mu}^{\oplus} = |\epsilon_{\theta\mu}^{\oplus}| e^{i\delta_{\theta\mu}}, \epsilon_{\theta\tau}^{\oplus} = |\epsilon_{\theta\tau}^{\oplus}| e^{i\delta_{\theta\tau}}, \epsilon_{\mu\tau}^{\oplus} = |\epsilon_{\mu\tau}^{\oplus}| e^{i\delta_{\mu\tau}}$$

► the 6th fit allows for **more freedom in the NSI flavour structure:**

$$H_{\text{mat}}(x) = R_{12}(\varphi_{12}) R_{13}(\varphi_{13}) \sqrt{2} G_F N_e(x) \text{diag}(\epsilon_{\oplus}^{\oplus}, 0, 0) R_{13}^{\dagger}(\varphi_{13}) R_{12}^{\dagger}(\varphi_{12})$$

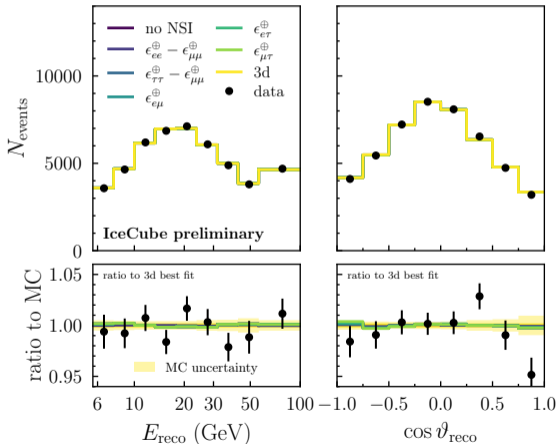
**Euler rotation
angles**

**overall matter
potential strength**

Esteban et al.
(2018)

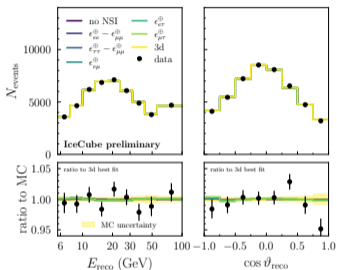
- ▶ rather weak deviations between the seven fits (including no NSI)

- goodness of fit (trials) approx. 20 % in all cases



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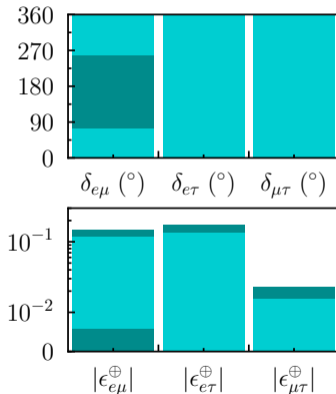
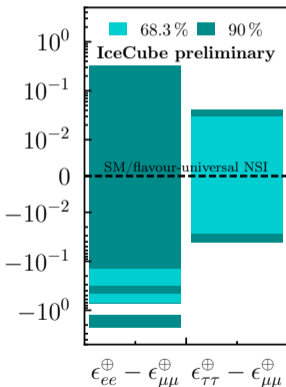
- goodness of fit (trials) approx. 20% in all cases



▶ **@ 90% C.L. (Wilks):**

all fits compatible with either flavour-universal NSI or no NSI
normal ordering (NO) assumed

68.3% (90%) allowed region



- ▶ NSI $e - \mu$ non-universality rescales Wolfenstein's SM matter potential:

$$V_{CC} \rightarrow V_{CC} \times \left(1 + \epsilon_{ee}^{\oplus} - \epsilon_{\mu\mu}^{\oplus} \right)$$

- ▶ degenerate with mass ordering (lacking sensitivity to θ_{12} & δ_{CP}):

$$\begin{aligned} \epsilon_{ee}^{\oplus} - \epsilon_{\mu\mu}^{\oplus} &\leftrightarrow - \left(\epsilon_{ee}^{\oplus} - \epsilon_{\mu\mu}^{\oplus} \right) - 2 \\ &\iff \Delta m_{31}^2 \leftrightarrow -\Delta m_{32}^2 \end{aligned}$$

Coloma, Schwetz (2016)

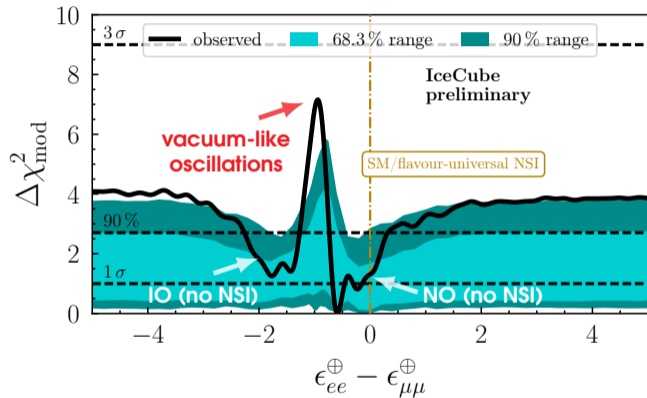
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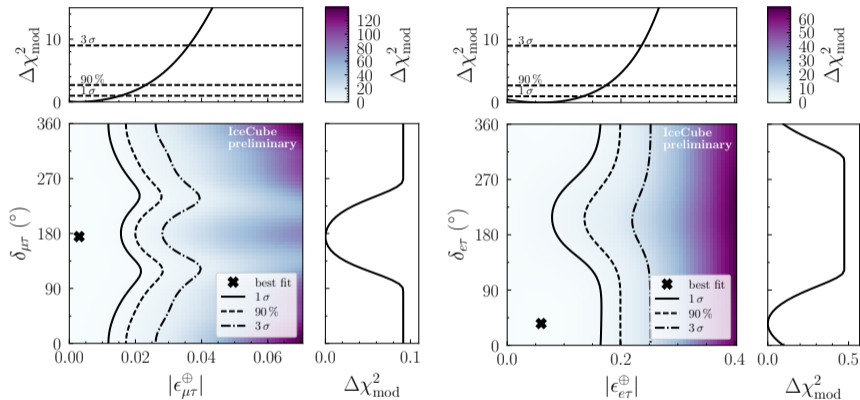
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Coloma, Schwetz (2016)



90% C.L.: $[-2.26, -1.27] \oplus [-0.74, 0.32]$

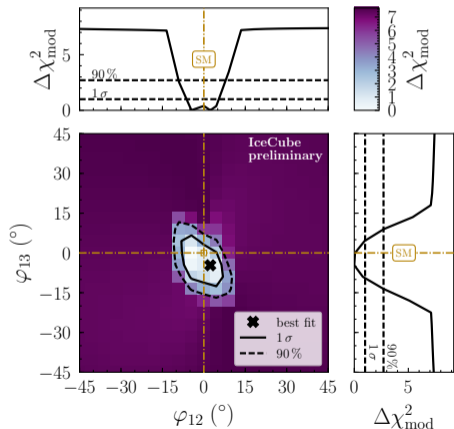


90 % limits:

- ▶ $|\epsilon_{\mu\tau}^\oplus| \leq 0.023$
- ▶ $|\epsilon_{e\tau}^\oplus| \leq 0.17$

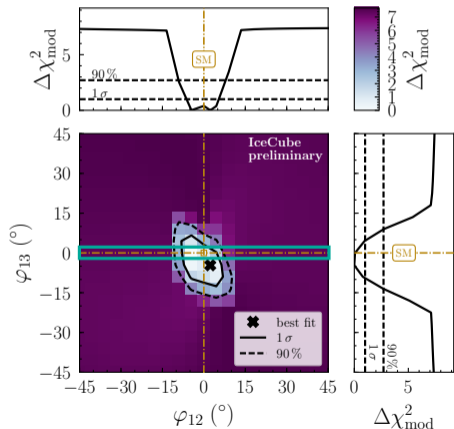
- ▶ correlated effects of magnitudes and complex phases
 - ▶ complex phases unconstrained

projection of 3d profile onto $(\varphi_{12}, \varphi_{13})$, i.e., minimised over ϵ_{\oplus}



- ▶ $\mathcal{O}(10^5)$ points tested; refined projections onto subdimensions
- ▶ at 90% C.L., data pins down both matter potential's strength (ϵ_{\oplus}) and flavour projection angles $(\varphi_{12}, \varphi_{13})$
(for two degenerate matter-potential eigenvalues, CP-conserving NSI, cf. Esteban et al. (2018))
- ▶ not shown:
@ 90%: $\epsilon_{\oplus} \in [-1.2, -0.3] \oplus [0.2, 1.4]$

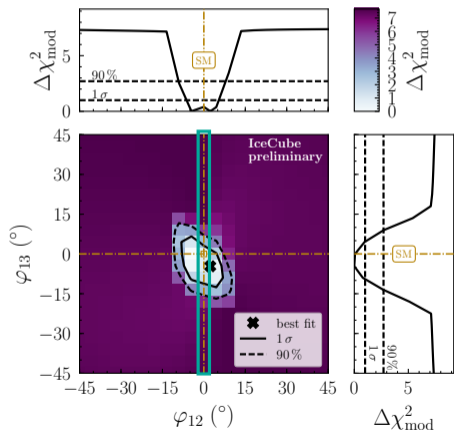
projection of 3d profile onto $(\varphi_{12}, \varphi_{13})$, i.e., minimised over ϵ_{\oplus}



$$\varphi_{13} = 0:$$

\Rightarrow no flavour-changing τ NSI

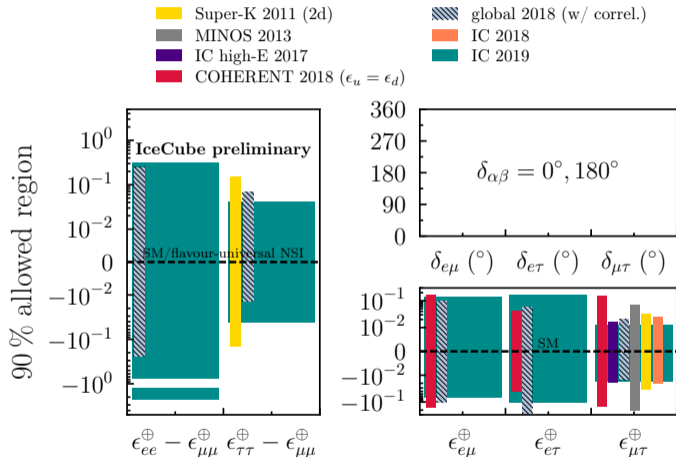
projection of 3d profile onto $(\varphi_{12}, \varphi_{13})$, i.e., minimised over ϵ_{\oplus}



$$\varphi_{12} = 0:$$

\Rightarrow no flavour-changing μ NSI

DeepCore NSI search: limit comparison

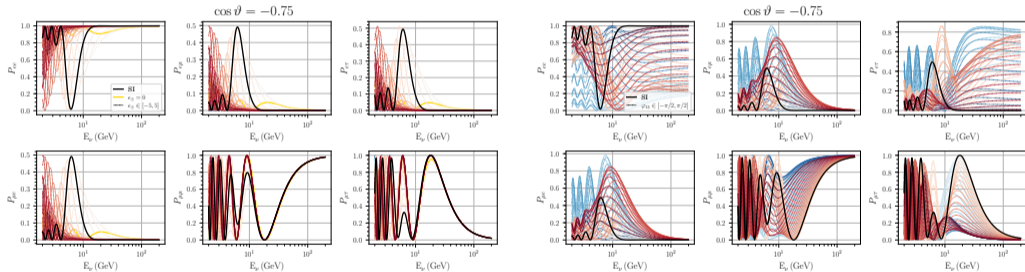


References:

- ["Super-K 2011"](#)
- ["MINOS 2013"](#)
- ["IC high-E 2017"](#)
- ["COHERENT 2018"](#)
- ["global 2018"](#)
- ["IC 2018"](#)
- ["IC 2019"](#): this work

- ▶ most common assumption:
real NSI
(= no new sources of CPV)

- ▶ broad NSI impact on atmospheric neutrinos
 - ▶ signals in all oscillation channels & across wide range of energies: from GeV to TeV



⇒ **promising search strategy:**

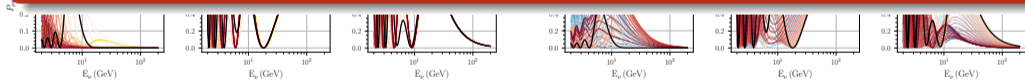
- ▶ use standard matter effects at few GeV as "anchor"
- ▶ look for incompatibility with high-energy data
- ▶ also reduce summation over $\nu + \bar{\nu}$

- ▶ broad NSI impact on atmospheric neutrinos
 - ▶ signals in all oscillation channels & across wide range of energies: from GeV to TeV

⇒ need unified low- (DeepCore) + high-energy (IceCube) event sample

▶ technically challenging, but generally beneficial to BSM searches

▶ active topic within IceCube



⇒ promising search strategy:

- ▶ use standard matter effects at few GeV as "anchor"
- ▶ look for incompatibility with high-energy data
- ▶ also reduce summation over $\nu + \bar{\nu}$

- ▶ IceCube offers wide range of possibilities to search for new neutrino interactions
 - ▶ focus so far on **propagation (matter effects)**

well-motivated
theoretically

constraints
independent of
new-physics scale

- ▶ DeepCore 2nd generation NSI analysis performed
 - ▶ 10× statistics increase — extended energy range — all neutrino flavours

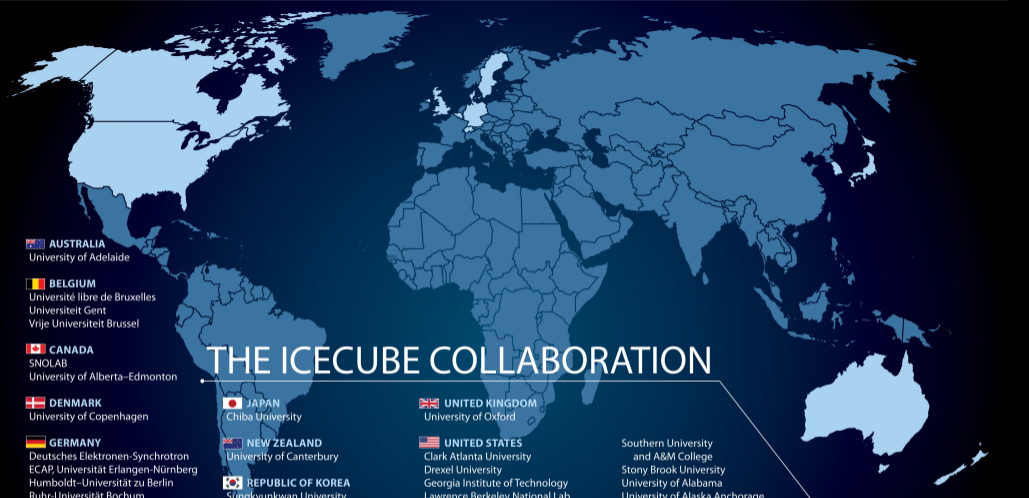
⇒ set limits on all effective NSI couplings for Earth matter
+ tested NSI flavour structure in three-dim. parameter space

- ▶ work toward next-generation IceCube & DeepCore NSI analyses ongoing
 - ▶ model assumptions
 - ▶ statistical approach
 - ▶ event selection
 - ▶ systematics

- ▶ IceCube offers wide range of possibilities to search for new neutrino interactions
 - ▶ focus so far on **propagation (matter effects)** well-motivated theoretically constraints independent of new-physics scale

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THE ICECUBE COLLABORATION

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University of Adelaide

 **BELGIUM**

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Universiteit Gent
Vrije Universiteit Brussel

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(FWO-Vlaanderen)

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Swedish Polar Research Secretariat

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US National Science Foundation (NSF)

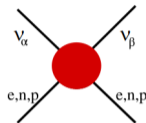


icecube.wisc.edu

backup

- ▶ a common approach: invoke scalar extensions of the SM

after integrating out new
(heavy) scalar boson ϕ



valid for $M_\phi^2 \gg Q^2$ †

analogous to Fermi
4-point interaction

- ▶ BSM theories yield effective coupling parameters:

$$\mathcal{L}_{\text{NSI}}^{\text{eff}} = -2\sqrt{2}G_F \epsilon_{\alpha\beta}^{ff',C} (\bar{\nu}_\alpha \gamma^\mu \nu_\beta) (\bar{f} \gamma_\mu P_C f')$$

$$\begin{aligned} \alpha, \beta &= e, \mu, \tau \\ f, f' &= e, u, d \\ C &= L, R \end{aligned}$$

- ▶ in the following:

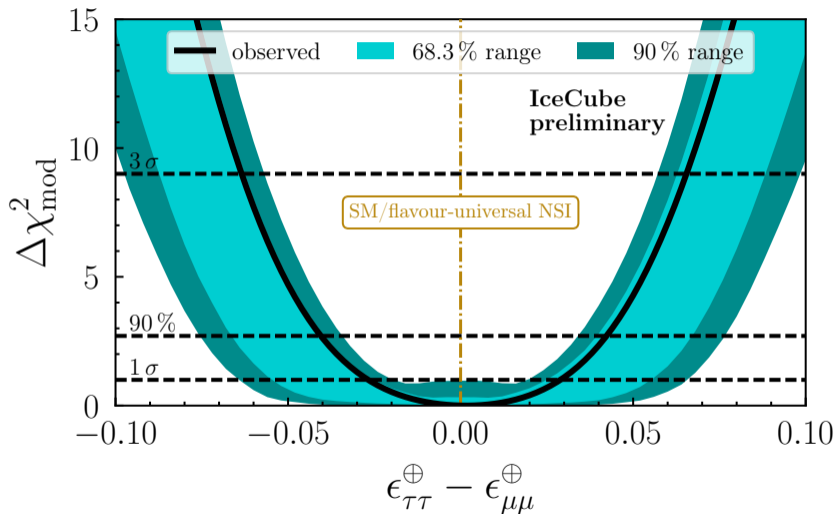
$$\epsilon_{\alpha\beta}^f \equiv \epsilon^{ff,V} = \epsilon^{ff,L} + \epsilon^{ff,R}$$

(vector interaction)

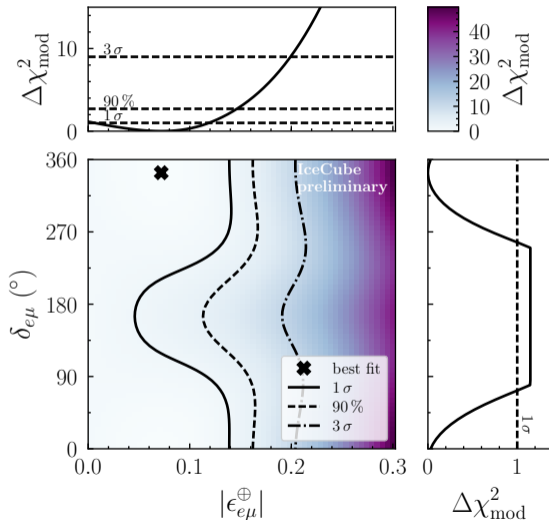
NC-like: $f = f'$

† $Q^2 = 0$ in coherent forward scattering \Rightarrow sensitivity independent of new physics scale

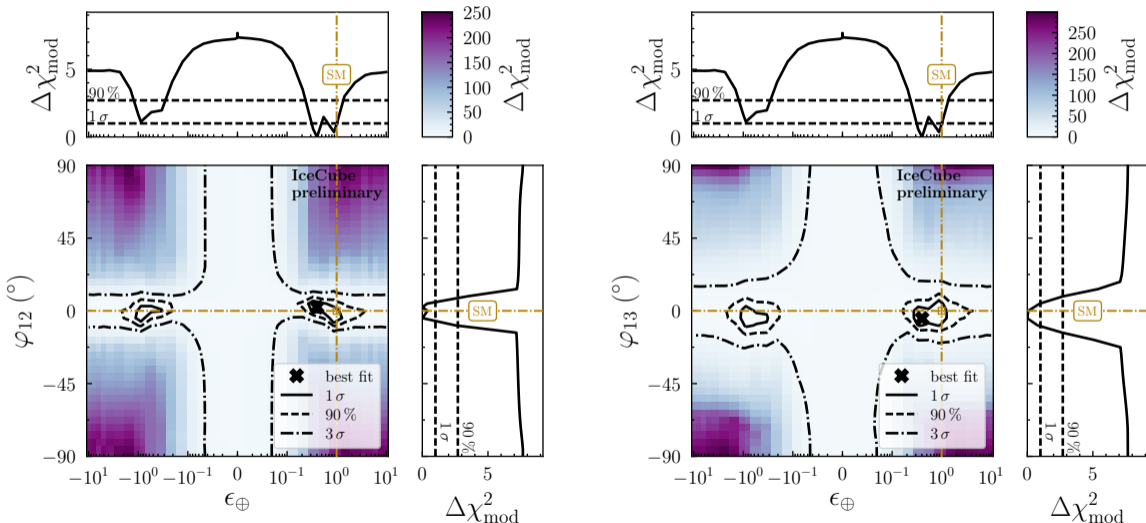
DeepCore NSI search: $\epsilon_{\tau\tau}^{\oplus} - \epsilon_{\mu\mu}^{\oplus}$



DeepCore NSI search: $\epsilon_{e\mu}^\oplus$



DeepCore NSI search: ϵ_{\oplus} and flavour structure



DeepCore NSI search: nuisance parameters

parameter	prior
<i>neutrino flux and cross sections:</i>	
ν_e/ν_μ ratio	1.0 ± 0.05
$\nu/\bar{\nu}$ ratio (σ)	0.0 ± 1.0
$\Delta\gamma_\nu$	0.0 ± 0.1
effective livetime (years)	...
M_A^{CCQE} (GeV)	$0.99^{+0.248}_{-0.149}$
M_A^{res} (GeV)	1.12 ± 0.22
NC normalisation	1.0 ± 0.2
<i>oscillation:</i>	
θ_{23}	...
Δm_{32}^2	...
<i>detector:</i>	
Optical Eff., Overall (%)	100 ± 10
Optical Eff., Lateral (σ)	0.0 ± 1.0
Optical Eff., Head-on	...
Bulk ice, scattering (%)	100 ± 10
Bulk ice, absorption (%)	100 ± 10

cf. PRD99 032007 (2019)

Previous NSI searches with IceCube



2018

low-energy

- ▶ **DeepCore** 3 year study using upgoing "golden" ν_μ CC only

⇒ test only $\epsilon_{\mu\tau}$:

$$-6.7 \times 10^{-3} < \epsilon_{\mu\tau}^d < 8.1 \times 10^{-3} \text{ (90\% C.L.)} \quad \text{Phys.Rev. D97 (2018) no.7, 072009}$$

2017

high-energy

- ▶ **external** study using 1 year dataset of high-energy through-going muon tracks (0.3 TeV–20 TeV)

⇒ limit on $\epsilon_{\mu\tau}$, marginalised over $\epsilon_{\tau\tau} - \epsilon_{\mu\mu}$:

$$-6 \times 10^{-3} < \epsilon_{\mu\tau}^d < 5.4 \times 10^{-3} \text{ (90\% C.I.)} \quad \text{JHEP 1701 (2017) 141}$$

2013

high-energy+

low-energy

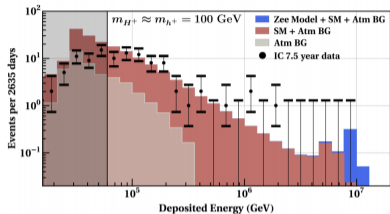
- ▶ **external** study using combination of 1 year high-energy (0.1 TeV–10 TeV) and 1 year low-energy (20 GeV–100 GeV) muon neutrinos

⇒ constrain both $\epsilon_{\mu\tau}$ and $\epsilon_{\tau\tau} - \epsilon_{\mu\mu}$:

$$-6.1 \times 10^{-3} < \epsilon_{\mu\tau}^d < 5.6 \times 10^{-3} \text{ (90\% C.I.)}$$

$$-3.6 \times 10^{-2} < \epsilon_{\tau\tau}^d - \epsilon_{\mu\mu}^d < 3.1 \times 10^{-2} \text{ (90\% C.I.)} \quad \text{JHEP 1306 (2013) 026}$$

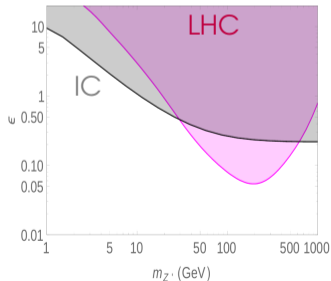
UHE resonance in $\bar{\nu}_\alpha e$ ("Zee burst")



arXiv:1908.02779

total νN cross section of high-energy astrophysical neutrinos from contained showers (6-year HESE)

$$\sigma_{\nu N}^{NSI} \lesssim \sigma_{\nu N}^{tot,cas} - \sigma_{\nu N}^{CC,IC} - \sigma_{\nu N}^{NC,SM}$$



arXiv:1907.07700