

The prompt atmospheric neutrino flux

Rikard Enberg

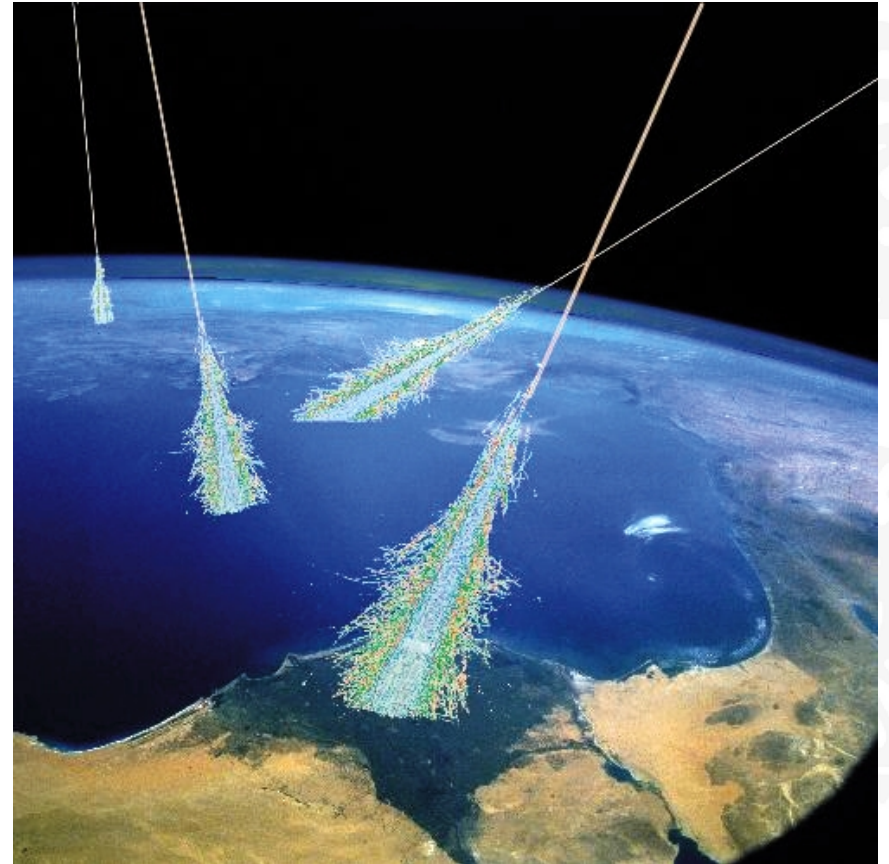
PPNT19, Uppsala, Oct 9, 2019



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Atmospheric neutrinos

- Cosmic rays bombard upper atmosphere and collide with air nuclei
- Very large CM energy \rightarrow Hadron production: pions, kaons, D-mesons ...
- Interaction & decay \Rightarrow cascade of particles
- Semileptonic decays \Rightarrow neutrino flux



Astropic of the day, 060814

Conventional neutrino flux

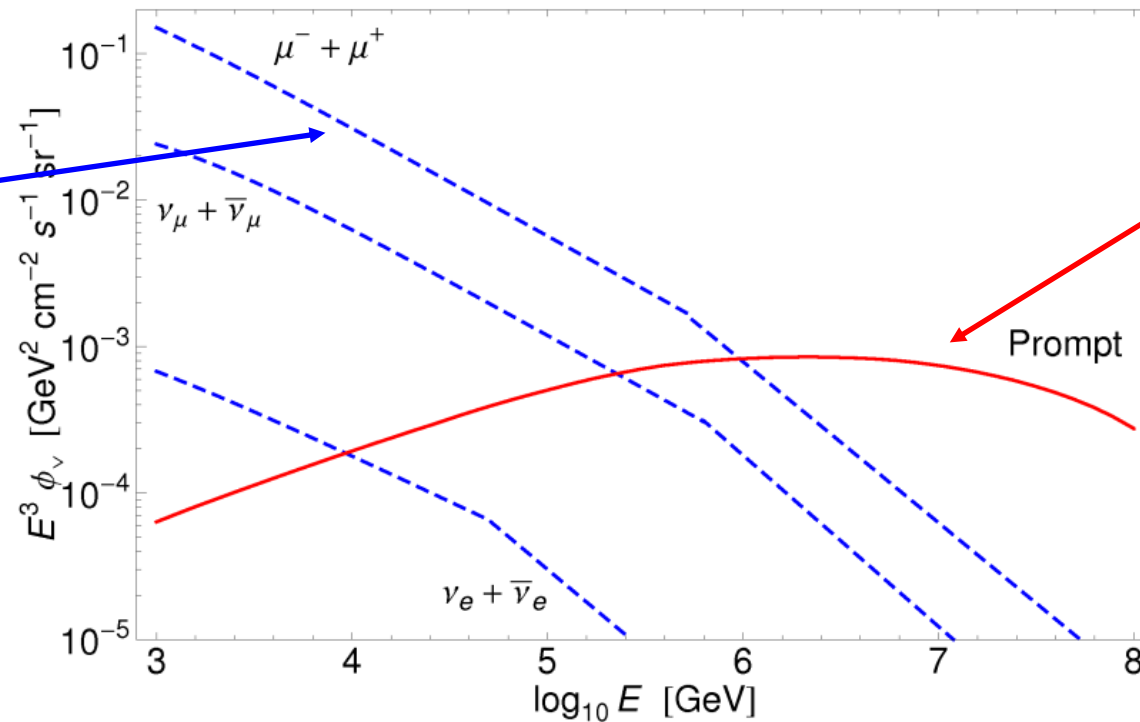
- Pions (and kaons) are produced in more or less every inelastic collision
- π^+ always decay to neutrinos: $BR(\pi^+ \rightarrow \mu^+ \nu_\mu) = 99.98 \%$
- *But π^\pm, K^\pm are long-lived* ($c\tau \sim 8$ meters for π^+)
 - \Rightarrow lose energy through collisions before decay
 - \Rightarrow neutrino energies are degraded
- This is called the **conventional neutrino flux**

Prompt neutrino flux

- Hadrons containing **heavy quarks** (*charm or bottom*) are **extremely short-lived**:
 - ⇒ decay before losing energy
 - ⇒ harder neutrino energy spectrum
- However, production cross-section is much smaller
- There is a cross-over energy above which prompt neutrinos dominate over the conventional flux
- This is called the **prompt neutrino flux**

Prompt vs conventional fluxes of atmospheric neutrinos

Pions & kaons:
long-lived
⇒ lose energy before decay



Charmed mesons:
short-lived
⇒ don't lose energy
⇒ harder spectrum

Prompt flux: Enberg, Reno, Sarcevic, arXiv:0806.0418 (**ERS**)

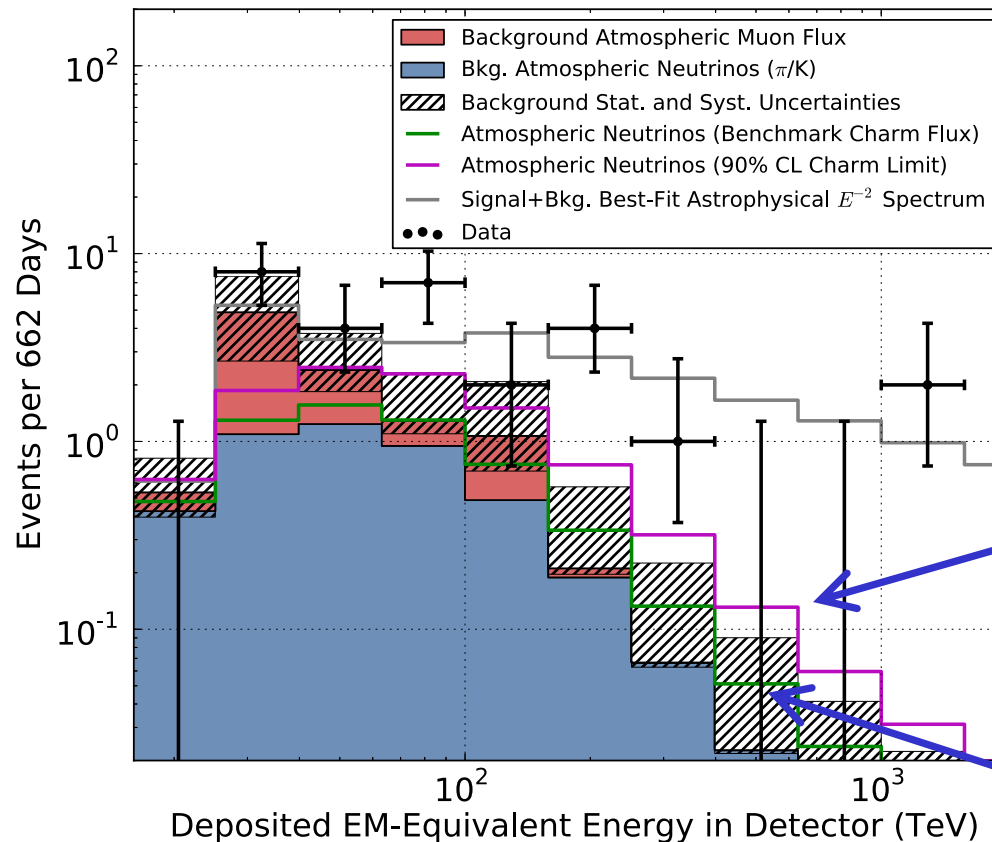
Conventional: Gaisser & Honda, Ann. Rev. Nucl. Part. Sci. **52**, 153 (2002)

Why are we interested?

- Atmospheric neutrinos are a background to extragalactic neutrinos
- Test beam for neutrino experiments
- Learn about cascades and the underlying production mechanism
- Higher energy pp collisions than in LHC: can maybe even learn something about QCD?

The IceCube events from 2013

The significance is sensitive to the prompt flux prediction



IceCube, arXiv:1311.5238

Prompt flux (ERS calc)

Calculations of the prompt flux

More recent:

Bhattacharya, RE, Reno, Sarcevic, Stasto, [arXiv:1502.01076](#) (**BERSS**)

Garzelli, Moch, Sigl, [arXiv:1507.01570](#) (**GMS**)

Gauld, Rojo, Rottoli, Sarkar, Talbert, [arXiv:1511.06346](#) (**GRRST**)

Bhattacharya, RE, Jeong, Kim, Reno, Sarcevic, Stasto, [arXiv:1607.00193](#)
(**BEJKRSS**)

PROSA Collaboration (Garzelli et al), [arXiv:1611.03815](#)

Benzke, Garzelli, et al., [arXiv:1705.10386](#)

Older but widely used:

Thunman, Ingelman, Gondolo, [hep-ph/9505417](#)

Pasquali, Reno, Sarcevic, [hep-ph/9806428](#)

Martin, Ryskin, Stasto, [hep-ph/0302140](#) (**MRS**)

RE, Reno, Sarcevic, [arXiv:0806.0418](#) [hep-ph] (**ERS**)

Calculations of the prompt flux

More recent:

Compare with

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Important message

QCD is crucial for some astrophysical processes:

- **Atmospheric neutrinos**
- Neutrino-nucleon cross-section @ high energy
- (Interactions in astrophysical sources?
See arXiv:0808.2807 and arXiv:1407.2985)

For example:

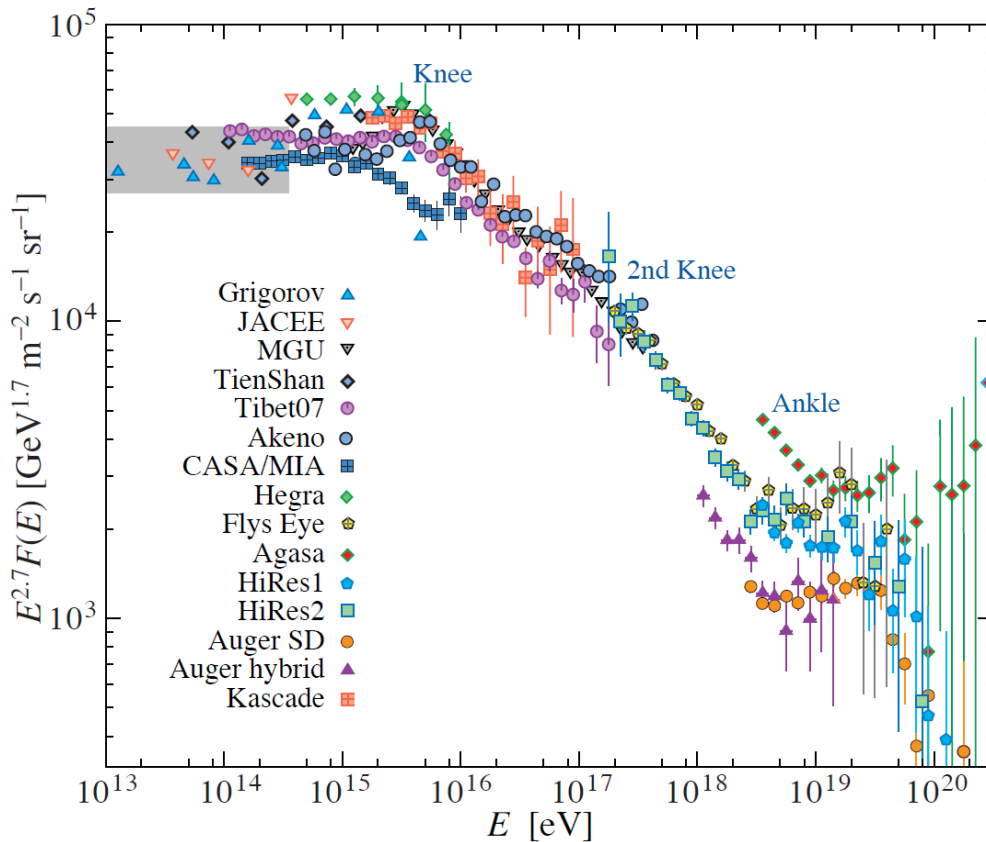
- What happens at **small Bjorken-x**? (Need very small x)
- **Forward** region (Hard to measure at colliders)
- **Fragmentation** of quarks \rightarrow hadrons (Non-perturbative, hard meas.)
- **Nuclear effects** in pA hard interactions

The calculation has many ingredients

- *Incident cosmic ray flux*
- Atmospheric density
- ***Cross section for heavy quarks in pp/pA collisions at extremely high energy (perturbative QCD)***
- Rescattering of nucleons, hadrons (hadronic xsecs) (scattering lengths)
- Decay spectra of charmed mesons & baryons (decay lengths)
- Cascade equations and their solution (Semi-analytic: spectrum-weighted Z-moments)

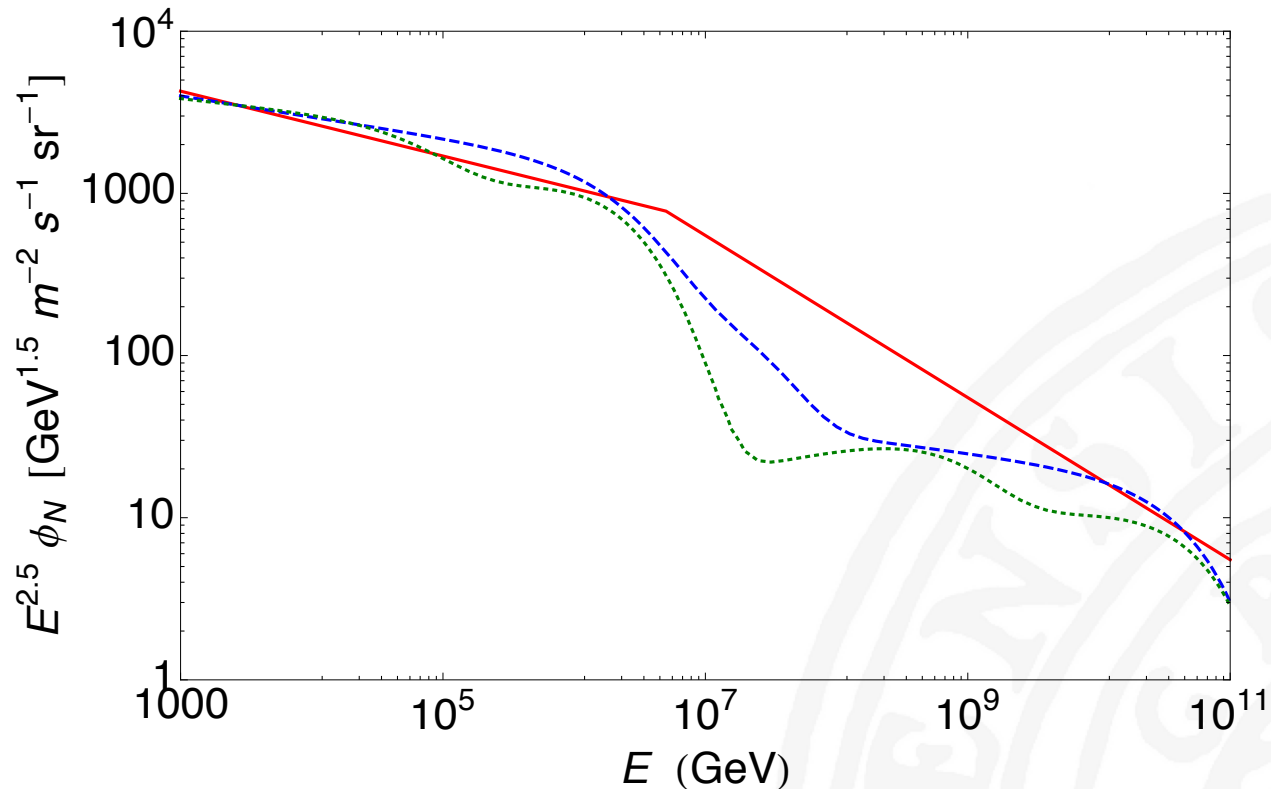
Cosmic rays (CR)

Plot from Particle Data Group



- Knees and ankles → seems natural to associate different sources with different energy ranges of the CR flux
- Highest energies: Extragalactic origin? → GRBs, AGNs, or more exotic
- Lower energies: Galactic origin? → SNRs etc

Incident cosmic ray flux: nucleons



Solid red = Broken power law (old standard)

Dashed blue = Gaisser all proton (H3p)

Dotted green = Gaisser, Stanev, Tilav (GST4)

Calculating the neutrino flux: Particle production

Particle physics inputs: energy distributions

$$\frac{dn(k \rightarrow j; E_k, E_j)}{dE_j} = \frac{1}{\sigma_{kA}(E_k)} \frac{d\sigma(kA \rightarrow jY, E_k, E_j)}{dE_j}$$
$$\frac{dn(k \rightarrow j; E_k, E_j)}{dE_j} = \frac{1}{\Gamma_k} \frac{d\Gamma(k \rightarrow jY; E_j)}{dE_j}$$

along with interaction lengths, or cooling lengths

$$\lambda_N(E) = \frac{\rho(h)}{\sigma_{NA}(E) n_A(h)}$$

→ Need the charm production cross section $d\sigma/dx_F$

Problem with QCD in this process

Charm cross section in LO QCD:

$$\frac{d\sigma_{\text{LO}}}{dx_F} = \int \frac{dM_{c\bar{c}}^2}{(x_1 + x_2)s} \sigma_{gg \rightarrow c\bar{c}}(\hat{s}) G(x_1, \mu^2) G(x_2, \mu^2)$$

where

$$x_{1,2} = \frac{1}{2} \left(\sqrt{x_F^2 + \frac{4M_{c\bar{c}}^2}{s}} \pm x_F \right)$$

CM energy is large: $s = 2E_p m_p$ so $x_1 \sim x_F$ and $x_2 \ll 1$

$x_F=1:$	$E=10^5 \rightarrow x \sim 4 \cdot 10^{-5}$	$x_F=0:$	$E=10^5 \rightarrow x \sim 6 \cdot 10^{-3}$
	$E=10^6 \rightarrow x \sim 4 \cdot 10^{-6}$		$E=10^6 \rightarrow x \sim 2 \cdot 10^{-3}$
	$E=10^7 \rightarrow x \sim 4 \cdot 10^{-7}$		$E=10^7 \rightarrow x \sim 6 \cdot 10^{-4}$

Very small x needed for forward processes (large x_F)!

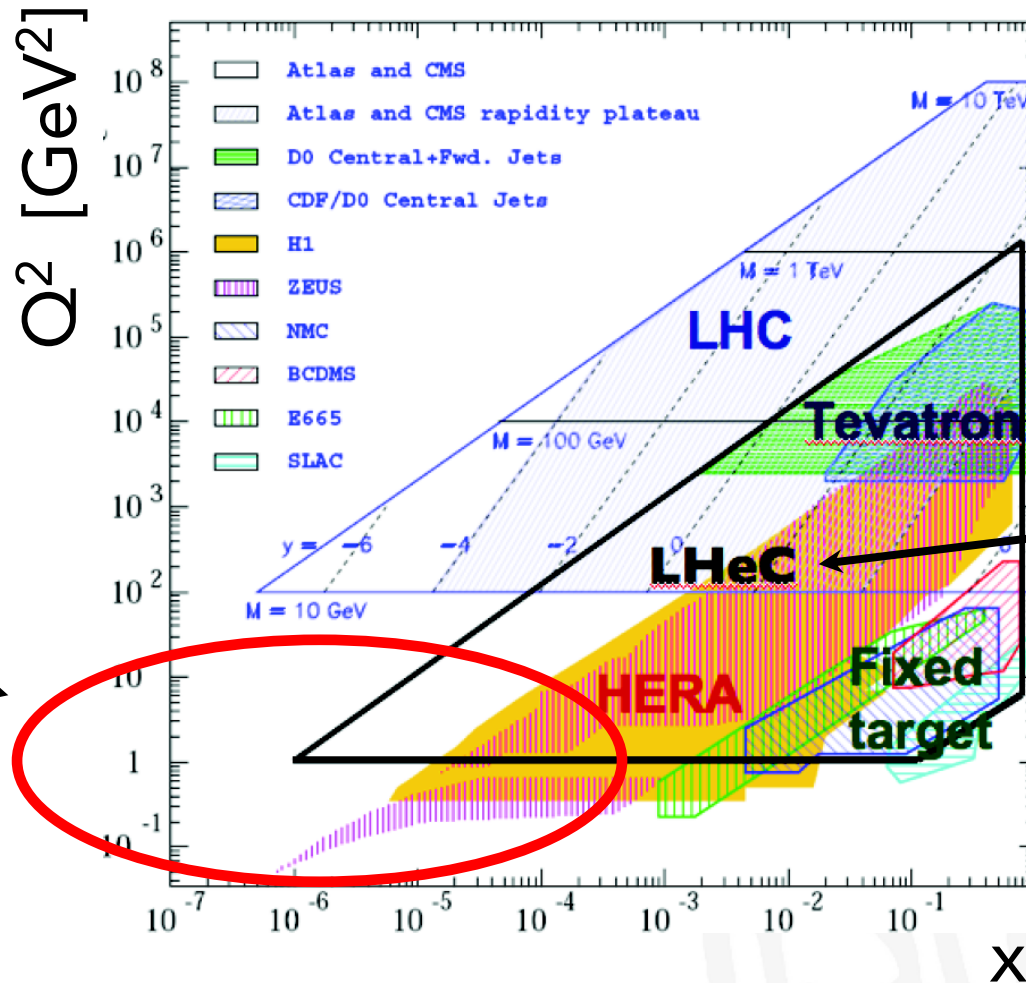
Problem with QCD at small x

- Parton distribution functions poorly known at small x
- At small x , must resum large logs: $\alpha_s \ln(1/x)$
- If logs are resummed (**BFKL**):
power growth $\sim x^{-\lambda}$ of gluon distribution as $x \rightarrow 0$
- Unitarity would be violated (T-matrix > 1)

How small x do we know?

- We haven't measured anything at such small x
- E.g. the MSTW pdf has $x_{\min} = 10^{-6}$
- **But that is an extrapolation!**
- HERA pdf fits: $Q^2 > 3.5 \text{ GeV}^2$ and $x > 10^{-4}$
- See Gao, Harland-Lang, Rojo, [arXiv:1709.04922](https://arxiv.org/abs/1709.04922) for more on pdfs

Kinematic plane

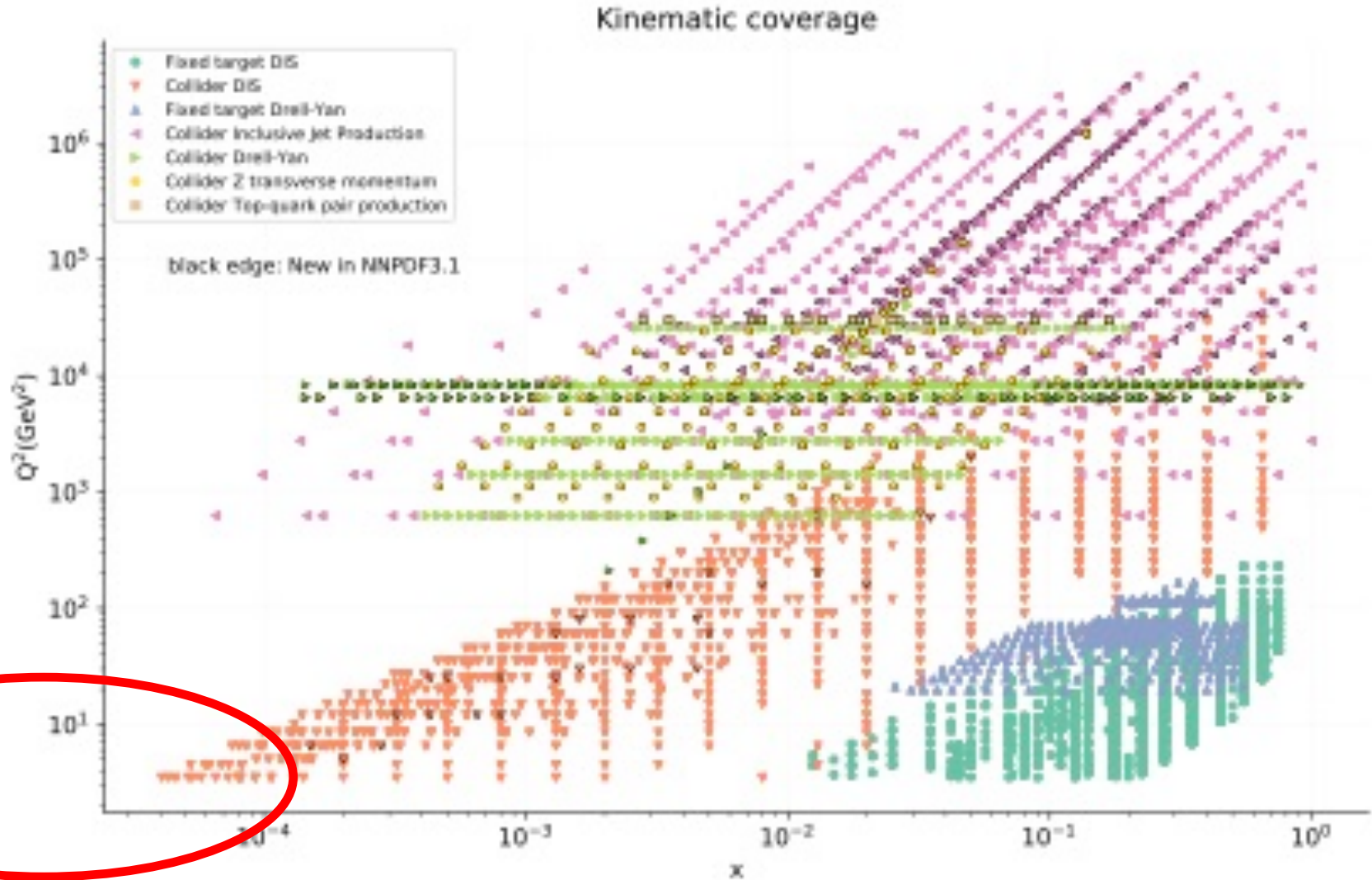


Where we need pdfs

Note LHeC!

HERA: $x_{\min} \sim 10^{-4}$ used for PDF fits ($Q^2 \sim 3.5 \text{ GeV}^2$)

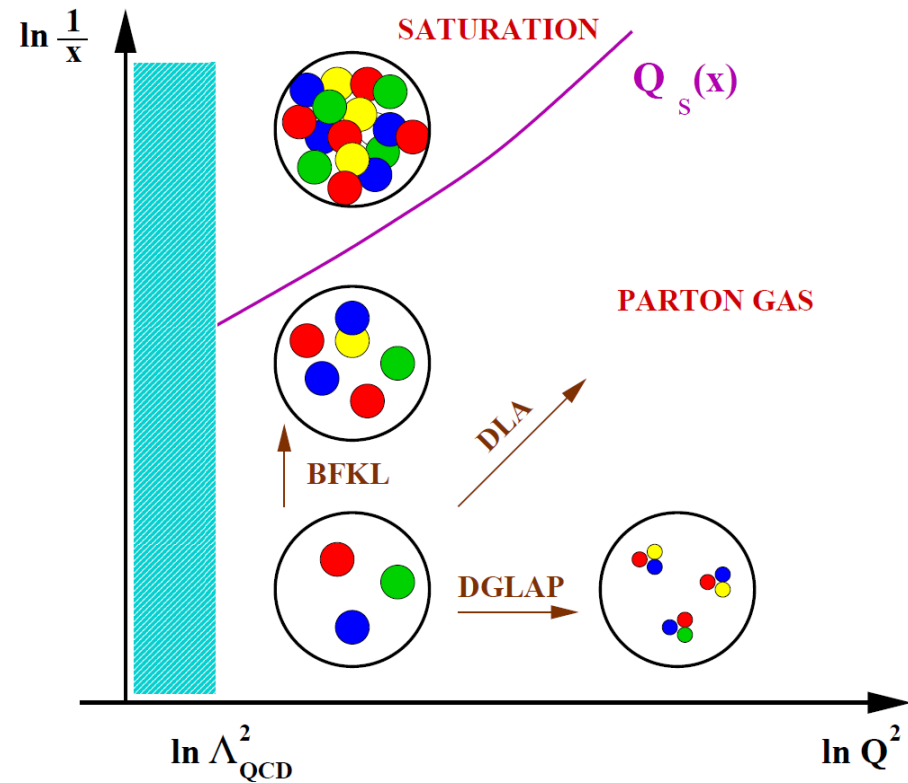
Kinematic plane of NNPDF3.1



R.D. Ball et al, arXiv:1706.00428

Parton saturation

- **Saturation** at small x :
 - Number of gluons in the nucleon becomes so large that gluons recombine
 - Reduction in the growth



- This is sometimes called the **color glass condensate**
- Non-linear QCD evolution: **Balitsky-Kovchegov equation**

Bhattacharya et al (BEJKRSS, 2016): Redo QCD calculations in many ways

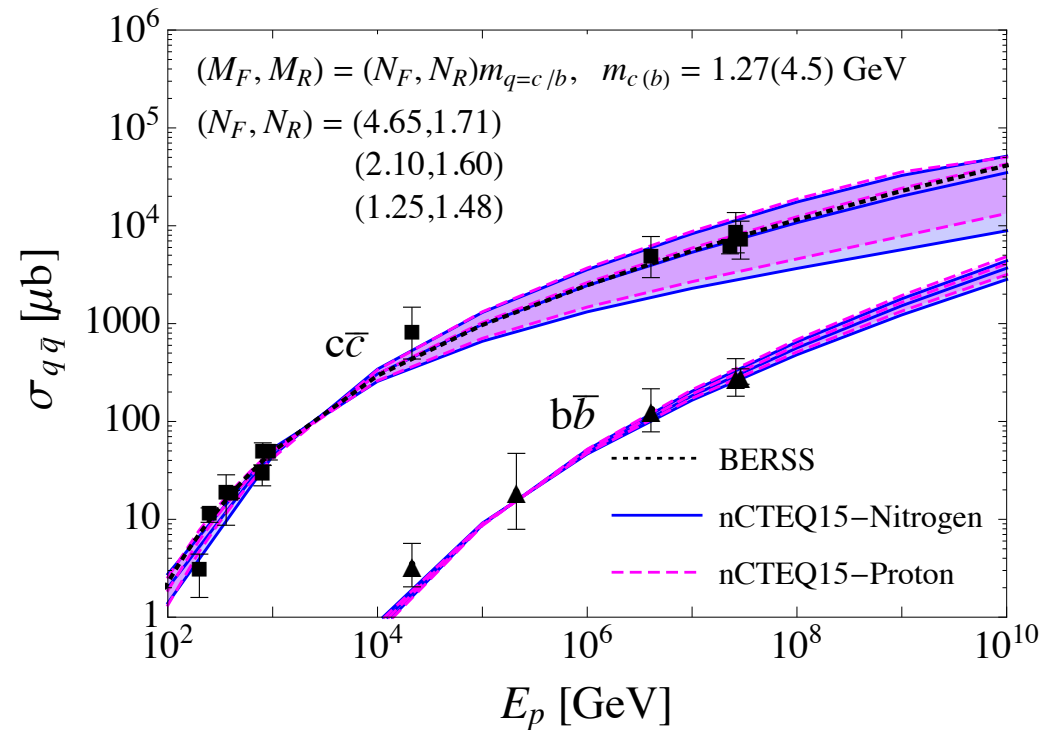
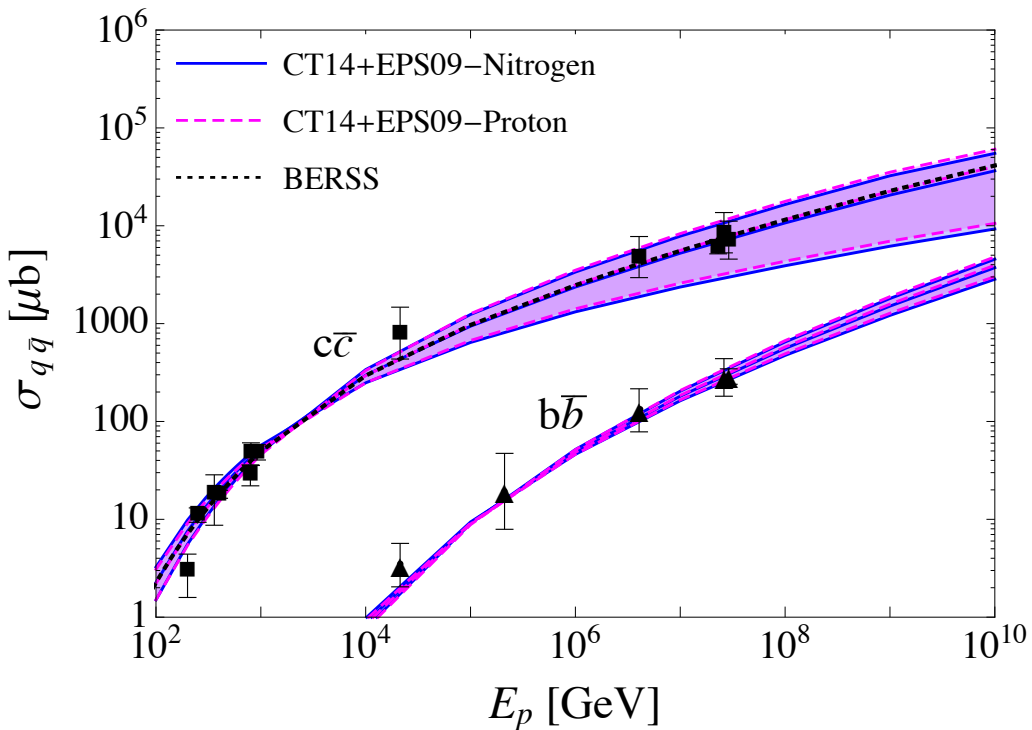
- ***Standard NLO QCD with newest PDFs***
 - BERSS updated with RHIC/LHCb input, uses Nason, Dawson, Ellis and Mangano, Nason, Ridolfi
- ***Dipole picture with saturation***
 - Approximate solution of Balitsky-Kovchegov equation
 - Update of ERS calc with new HERA fits + other dipoles
- ***kT factorization with and without saturation***
 - Resums large logs, $\alpha_s \log(1/x)$ with BFKL
 - Off-shell gluons, unintegrated PDFs (+ subleading...)
 - Kutak, Kwiecinski, Martin, Sapeta, Stasto (permutations)

Include scale variations, PDF errors, charm mass, etc
→ Plausible upper and lower limits on x_{sec}

Also include nuclear shadowing

- Partons are not in a free nucleon, but in a nucleus!
- Estimate shadowing with nuclear PDFs (nCTEQ15 and EPS09)
- Reduces flux by 10–30% at the highest energies
- Larger effect on the flux than on the total $\sigma(\text{cc})$ due to asymmetric $x_{1,2}$

$\sigma(cc)$ and $\sigma(bb)$



Data from RHIC, LHC and lower energies
 Total cross sections well described by all calculations
 (at high energies), nuclear shadowing small

(Error bands=scale variations and PDF uncertainties)

Differential cross sections (LHCb)

LHCb measured D-meson production at 7 and 13 TeV

Kinematical range: $p_T < 8 \text{ GeV}$, $0 < y < 4.5$

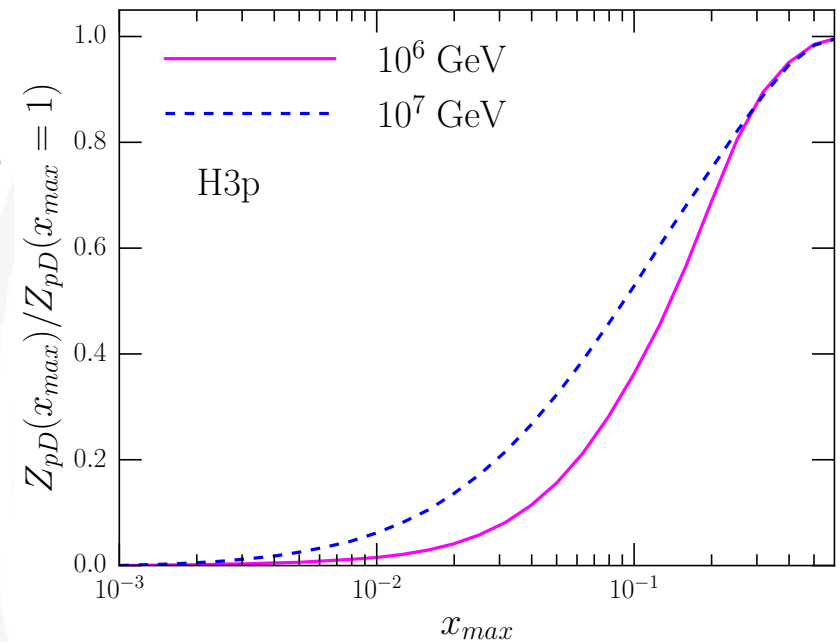
The flux is mostly sensitive to *large y and small p_T* .

Cumulative fraction of Z-moment
as function of x_F :

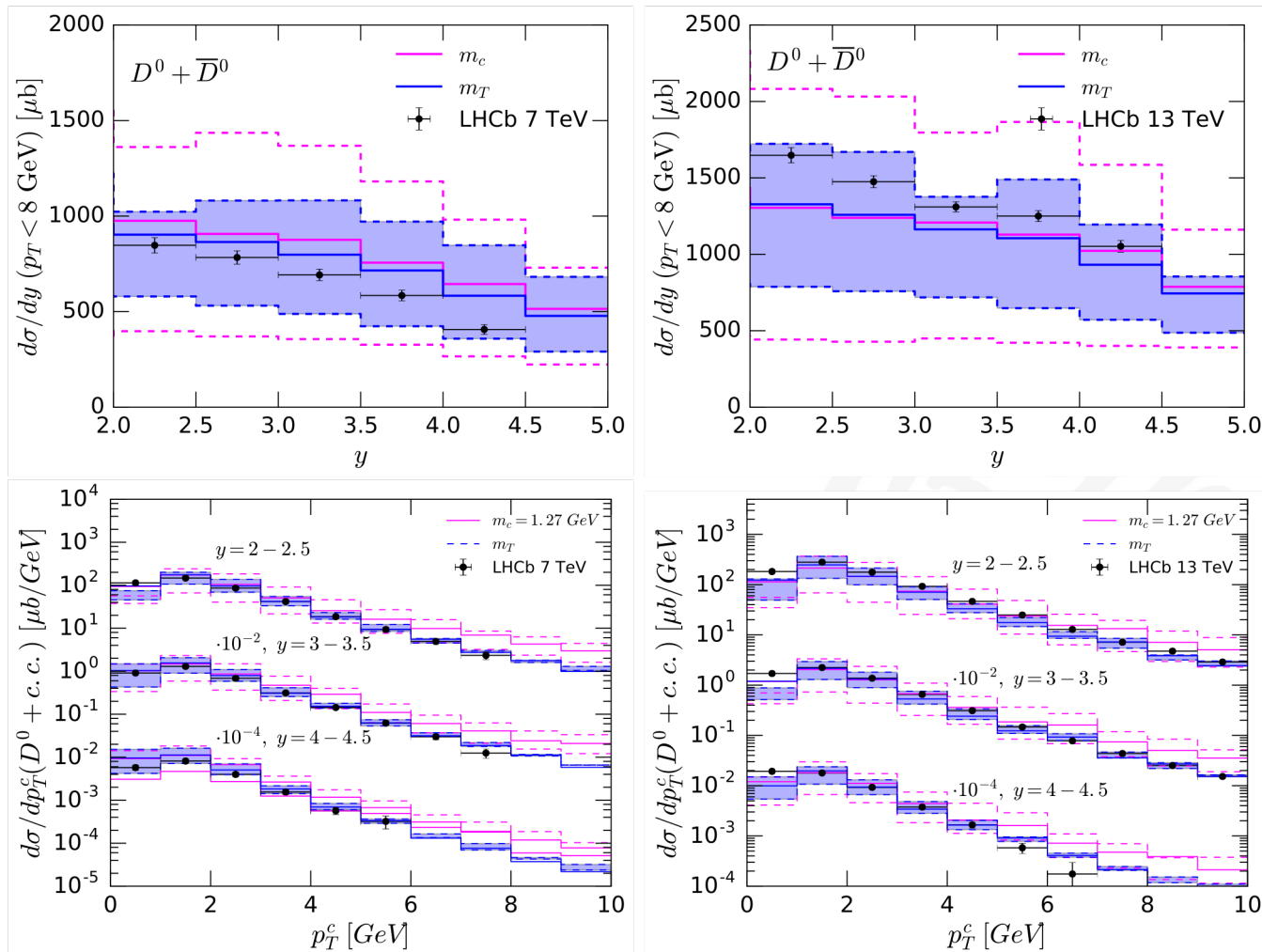
Estimate: 90% of Z_{pD} given by

$y > 4.9$ for $E_p = 10^6 \text{ TeV}$

$y > 5.7$ for $E_p = 10^7 \text{ TeV}$



Comparison of NLO QCD



Data from LHCb: arXiv:1302.2864 and arXiv:1510.01707

Prompt ν_μ ($=\nu_e=\mu$) fluxes

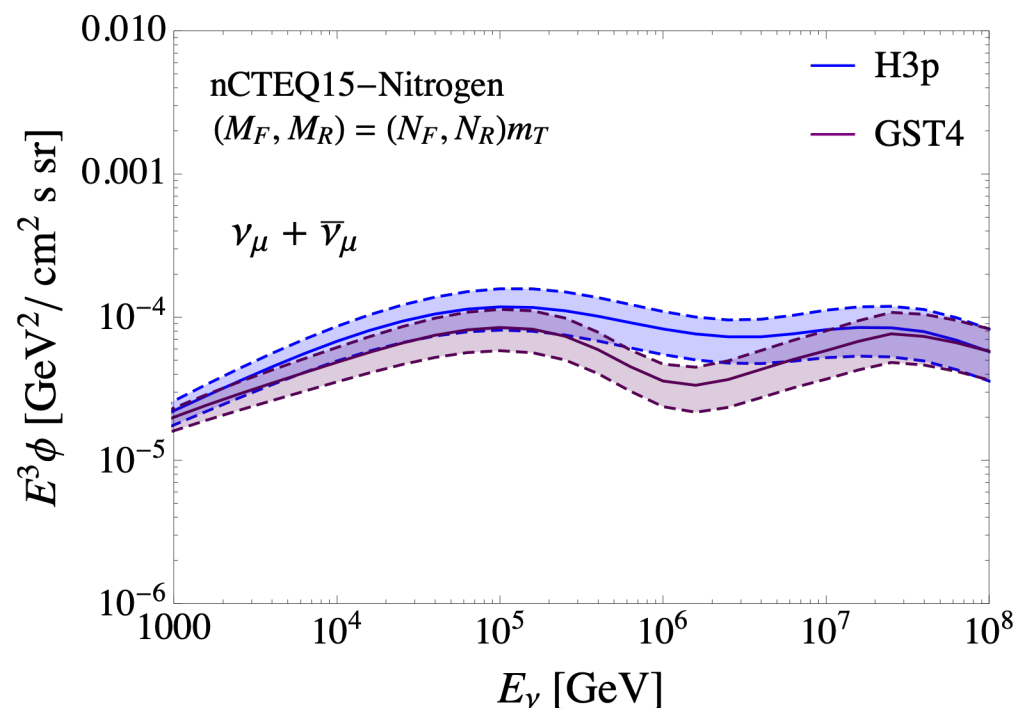
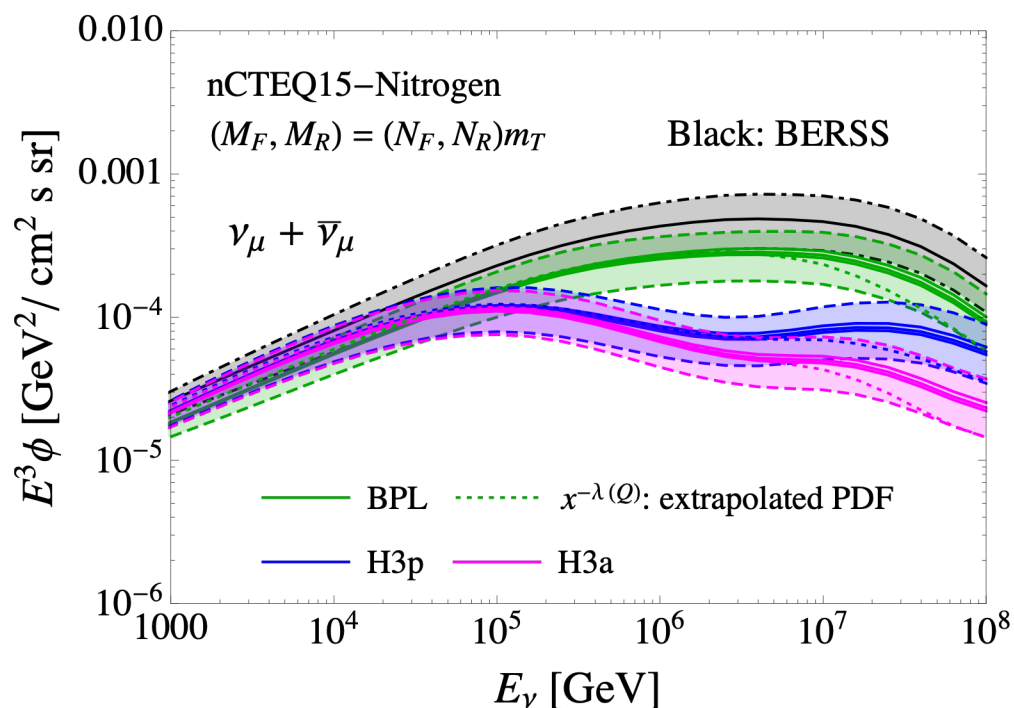
We have calculated prompt neutrino fluxes using all these variations in QCD, nuclear effects, cosmic ray fluxes.

Also compare to other calculations:

- RE, Reno, Sarcevic (ERS) [0806.0418](#)
- Bhattacharya et al (BERSS), [1502.01076](#)
- Garzelli, Moch, Sigl, [1506.08025](#)
- Gauld, Rojo, Rottoli, Sarkar, Talbert, [1511.06346](#)

→ estimate of theoretical uncertainties

NLO QCD

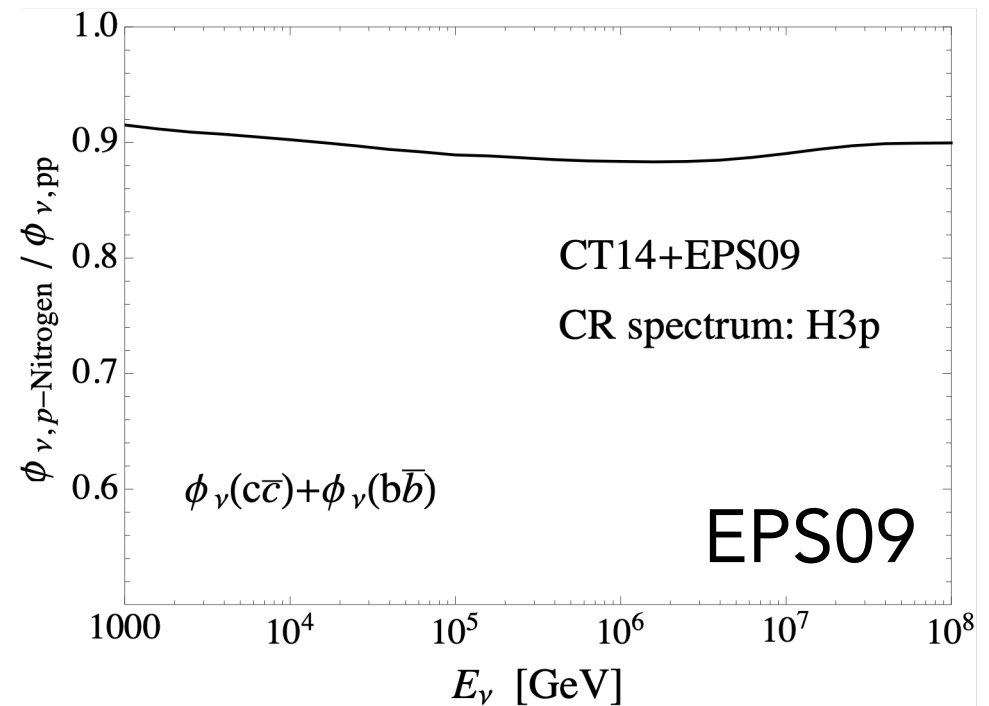
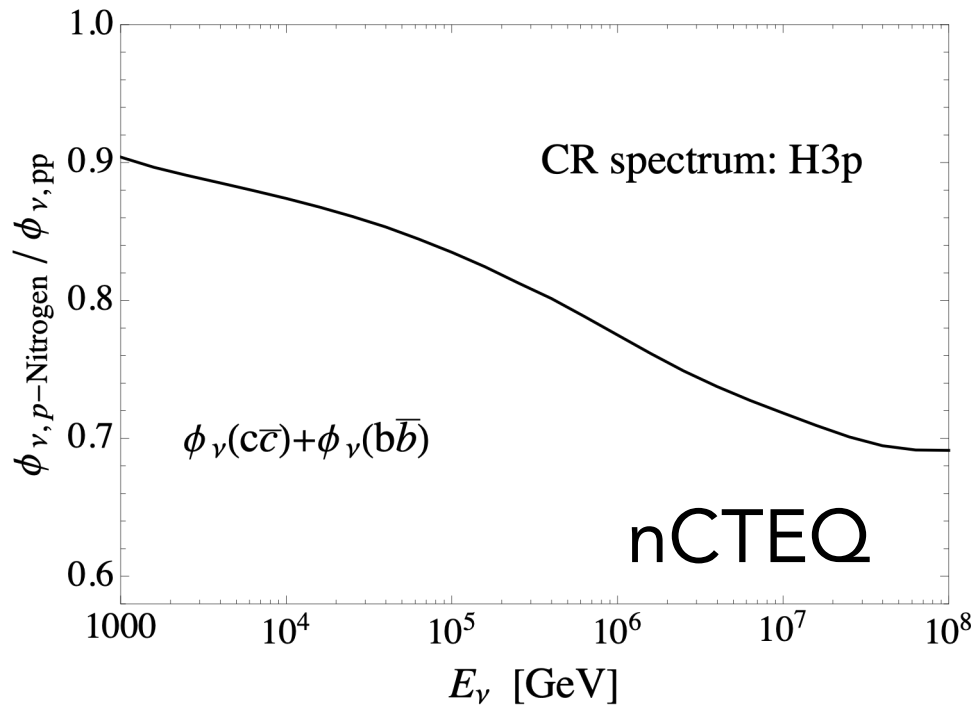


Compare with our BERSS NLO QCD and different cosmic ray fluxes

Difference to BERSS: bb now included, modified fragmentation fractions, nuclear effects (here: nCTEQ15)

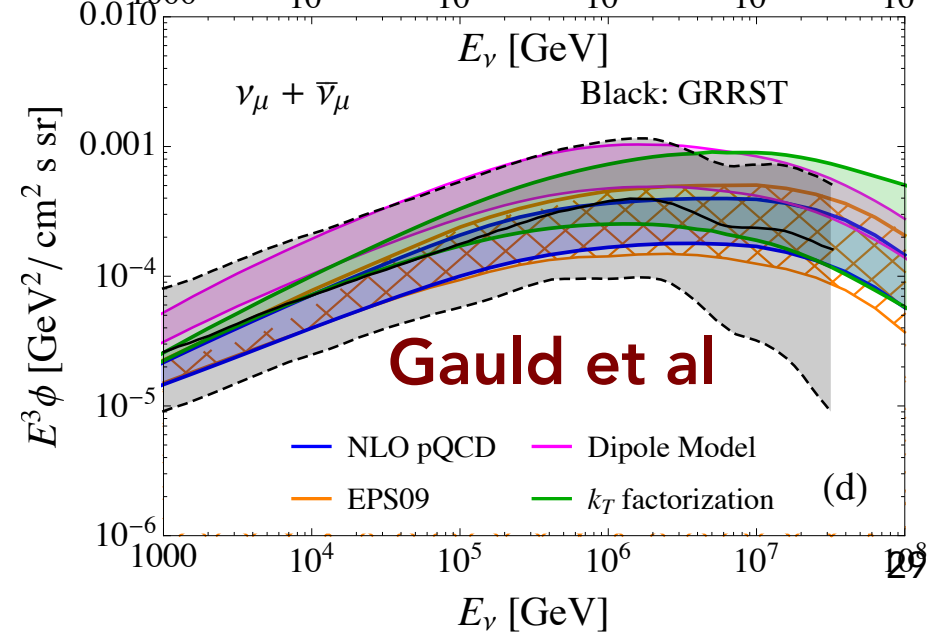
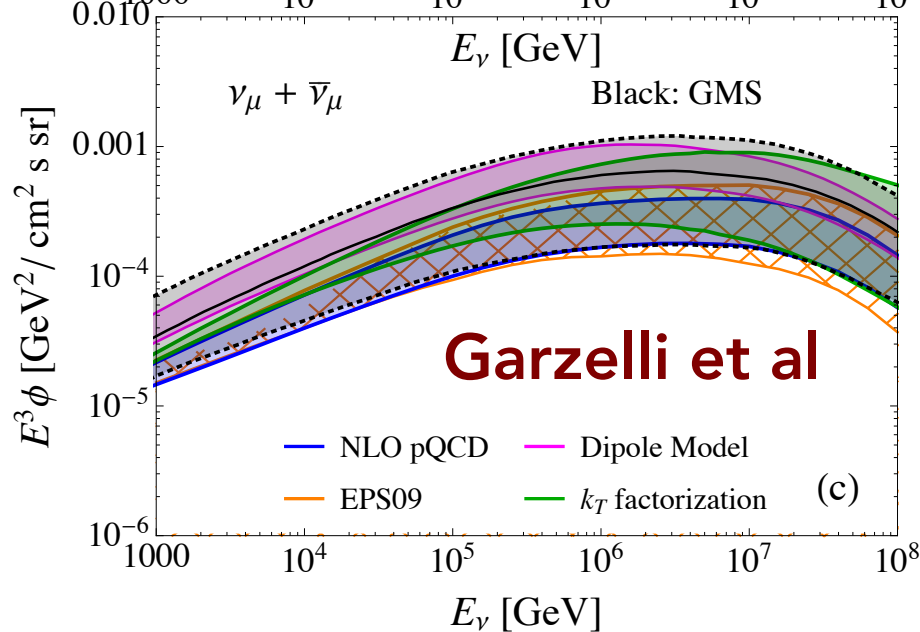
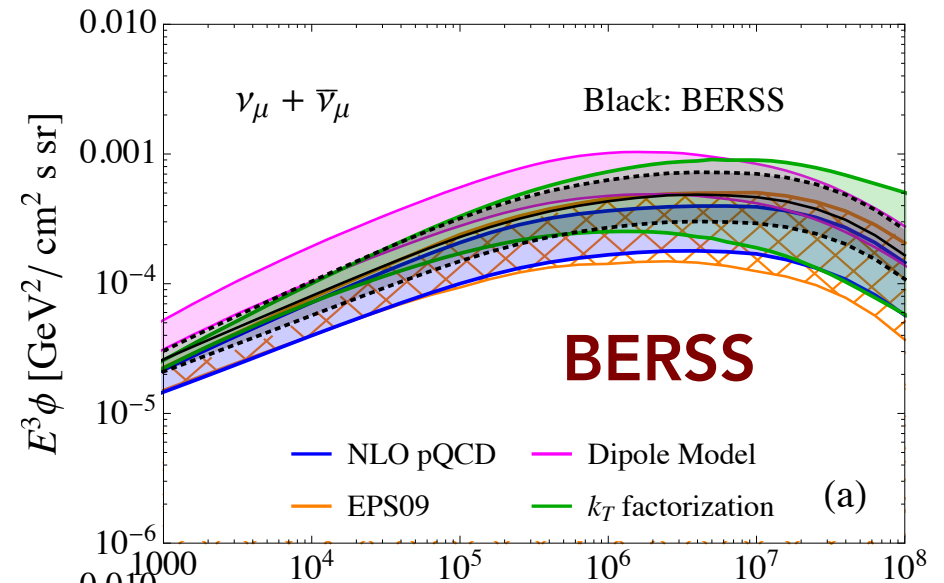
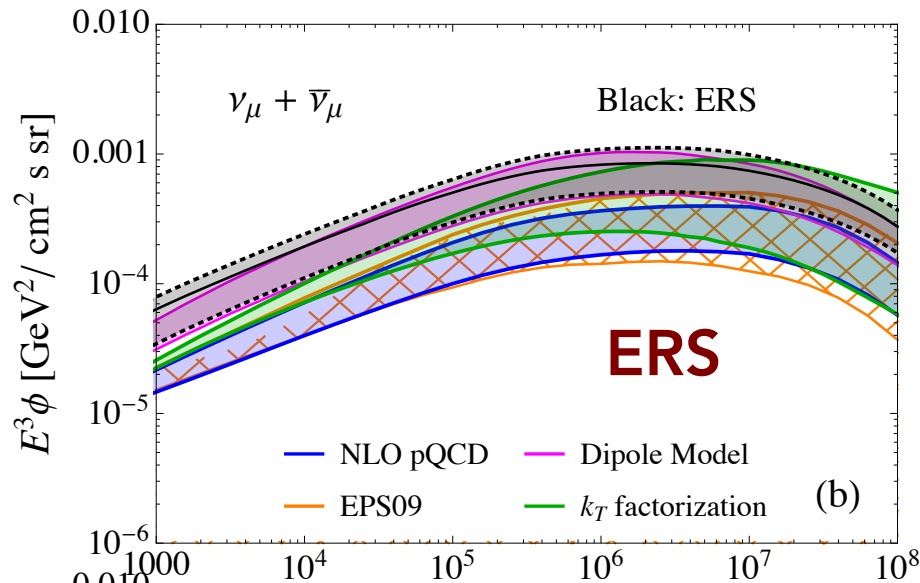
Overall: (30%, 40%, 45%) lower than BERSS at (10^3 , 10^6 , 10^8) GeV²⁷

Influence of nuclear shadowing

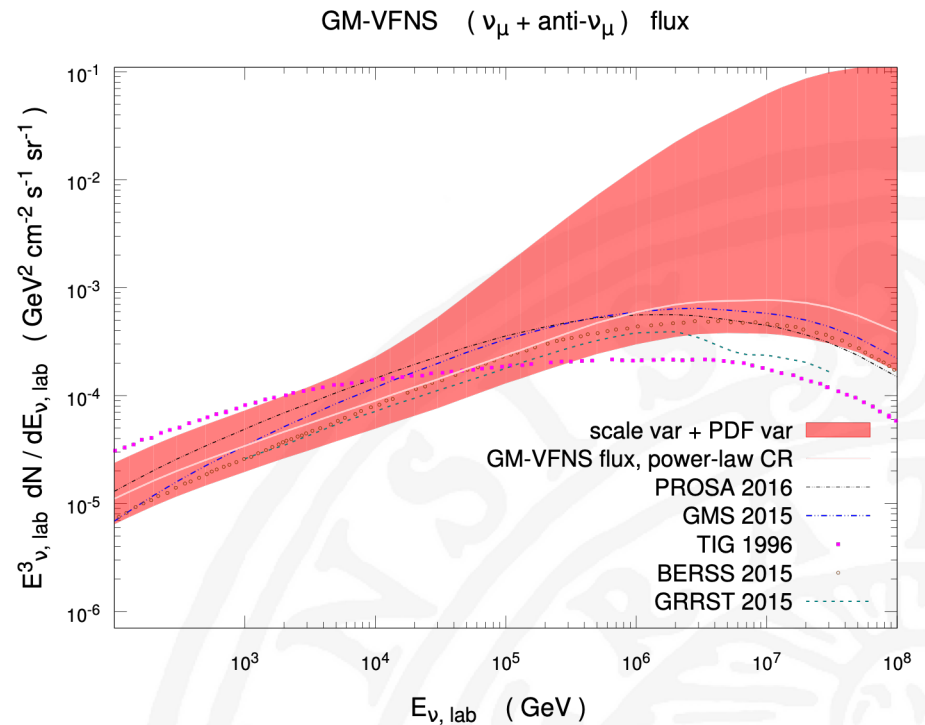
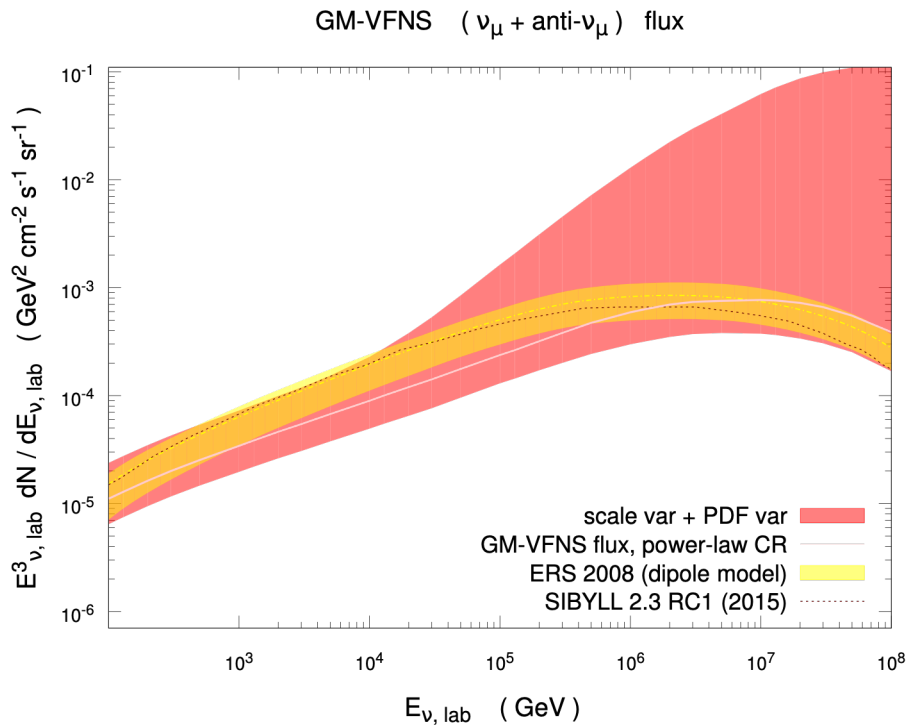


- Ratio of NLO QCD flux with and without nuclear effects
- 20–30% suppression from 10^5 to 10^8 GeV for nCTEQ
(only 4–13% for total cross section)
 - But much less for EPS (frozen at $x=10^{-6}$)

And now everything, using broken power law



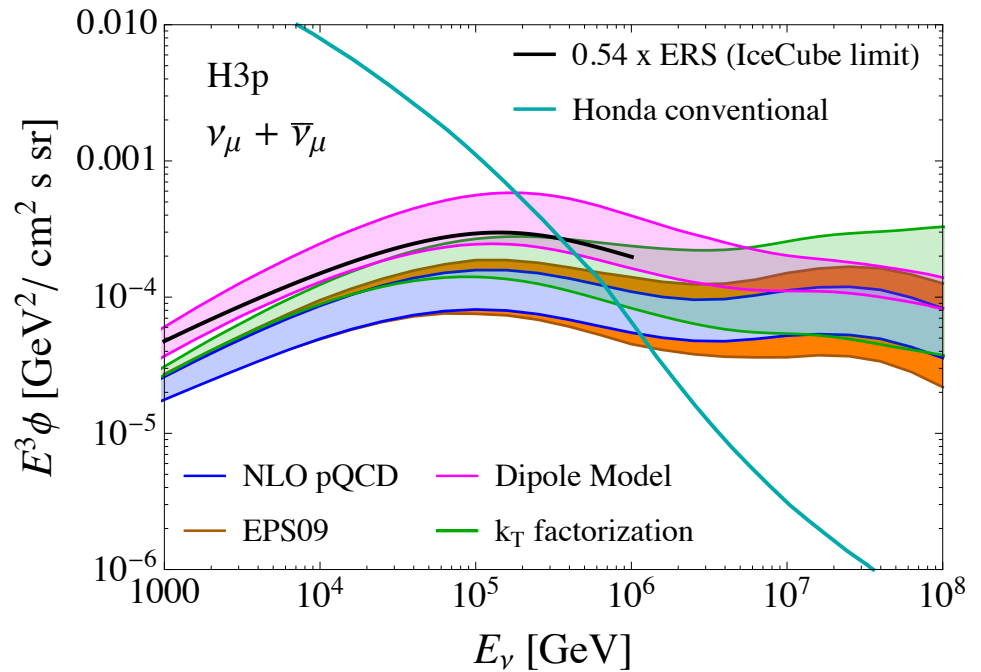
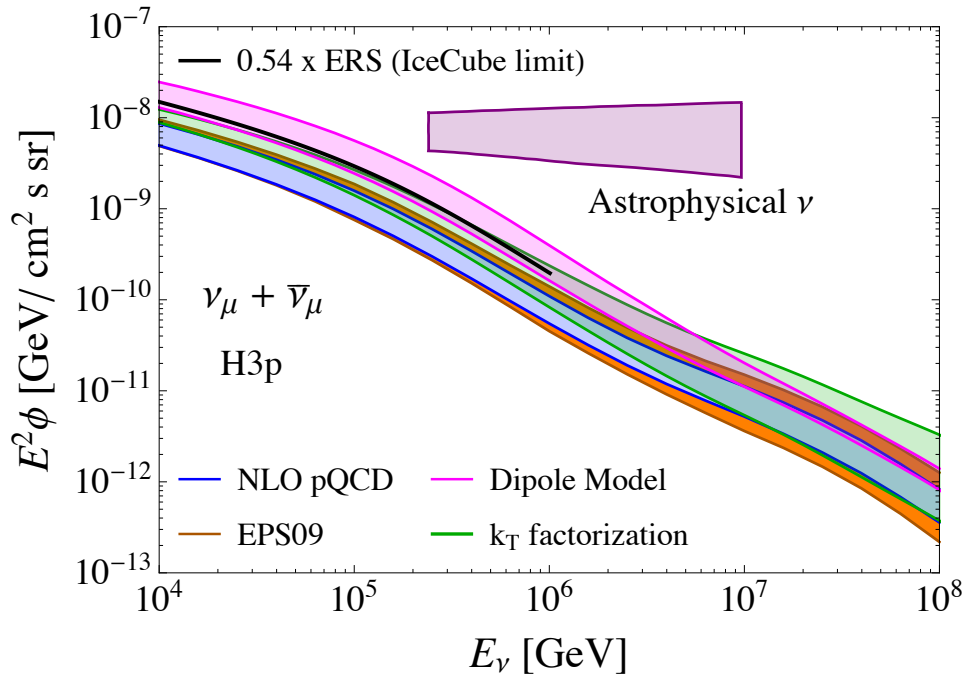
Benzke et al GM-VFNS calculation



pQCD calculation in “General Mass–Variable Flavor Number Scheme” (GM-VFNS)
M. Benzke, M. V. Garzelli, B. Kniehl, G. Kramer, S. Moch, G. Sigl, [arXiv:1705.10386](https://arxiv.org/abs/1705.10386)

The large pdf uncertainty at large energy arises from a particular set of CTEQ pdf fits (ct14nlo) – not constrained by data (but other sets don’t show this – situation unclear)

And what does IceCube say?

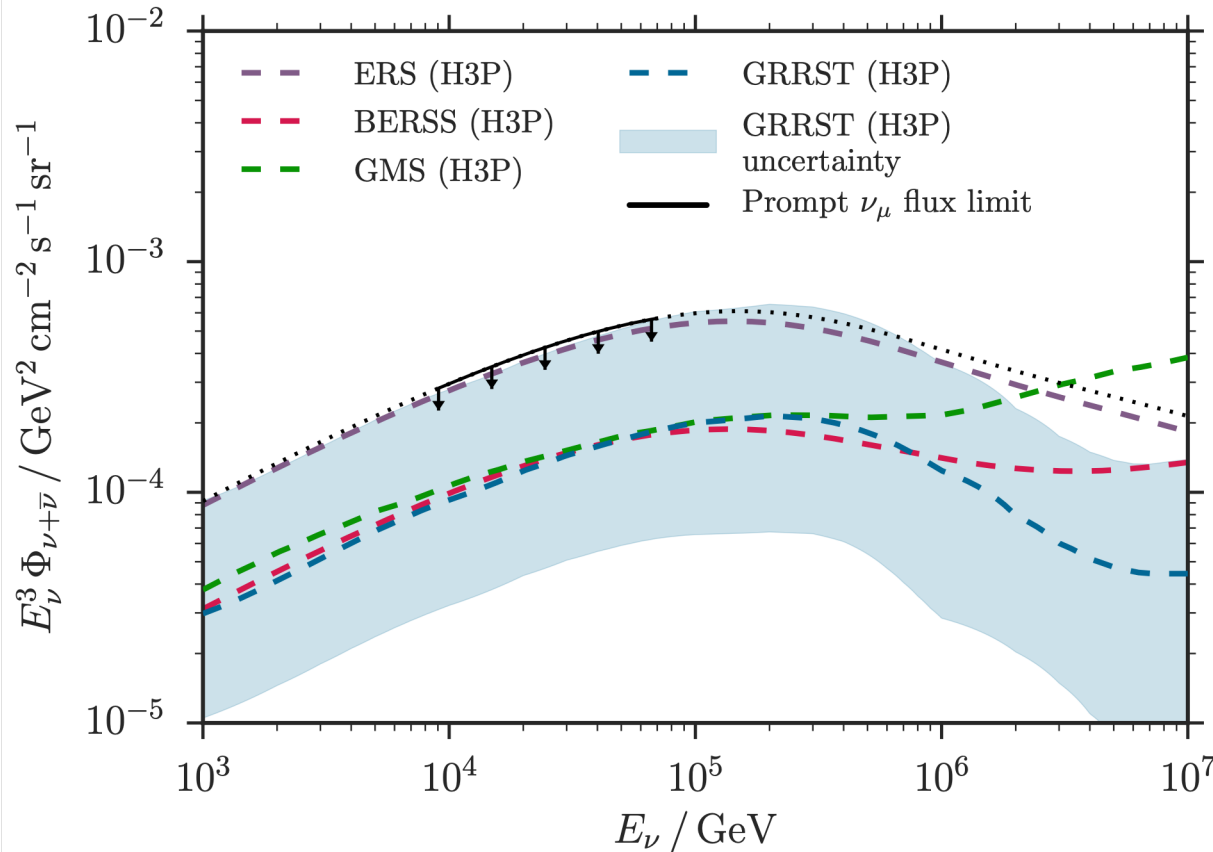


One IceCube limit (3 yr data) on the prompt flux sets a limit at 90% CL of

$0.54 \times$ (*ERS modified with H3p CR's*)

Best fit is $\phi_{\text{prompt}} = 0$

Another IceCube analysis (6 yrs)



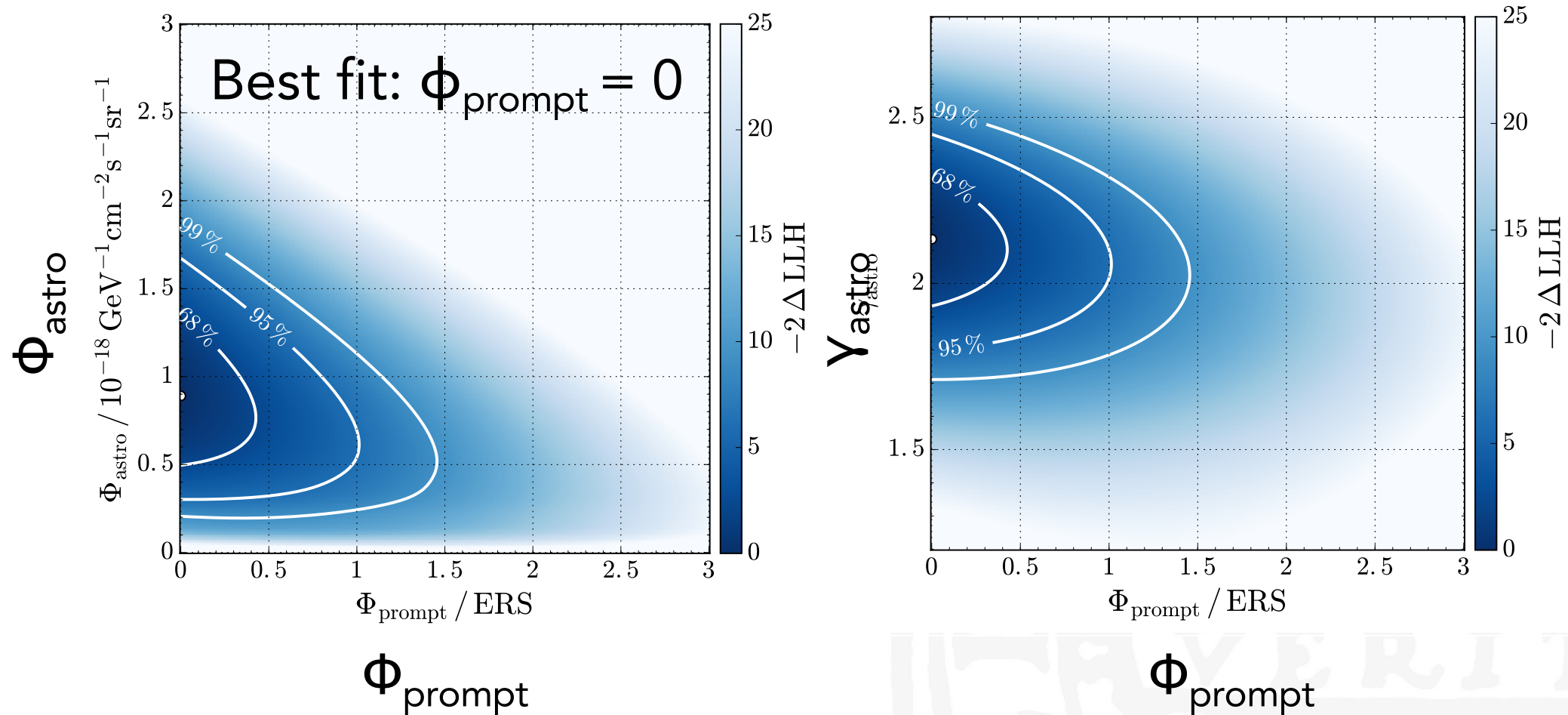
IceCube, arXiv:1607.08006

Limits on various calculations:

Model	Flux limit
ERS (H3p)	1.06
GMS (H3p)	≈ 2.9
BERSS (H3p)	≈ 3.0
GRRST (H3p)	≈ 3.1

Limit: $1.06 \times \text{ERS}$
 Best fit: $\phi_{\text{prompt}} = 0$

IceCube fits to Φ_{astro} and Φ_{prompt}



IceCube, arXiv:1607.08006

IceCube angular analysis:

Sample	Best Fit (ERS)	1σ Interval (90% CL)	$\sigma(\Phi_{\text{prompt}} > 0)$
Uncorrected	4.93	4.05-5.87 (3.55-6.56)	9.43
Marginalized Ang. Corr.	3.19	1.64-5.48 (0.98-7.26)	3.46

- Best fit $\Phi_{\text{prompt}} = 3.19 \times \text{ERS}$
- $0.98\text{--}7.26 \times \text{ERS}$ @ 90% CL
- $\Phi_{\text{prompt}} = 0$ excluded at 3.46σ

IceCube, arXiv:1506.07981

Conclusions

- The prompt neutrino flux poses one of the questions in neutrino astroparticle physics
 - How large is the flux?
 - Why hasn't it been discovered?
 - What is the proper way to calculate it?
- We think we know what we don't know about how to calculate it – more accelerator and cosmic ray data needed!