

# Heavy neutrino searches with Icecube

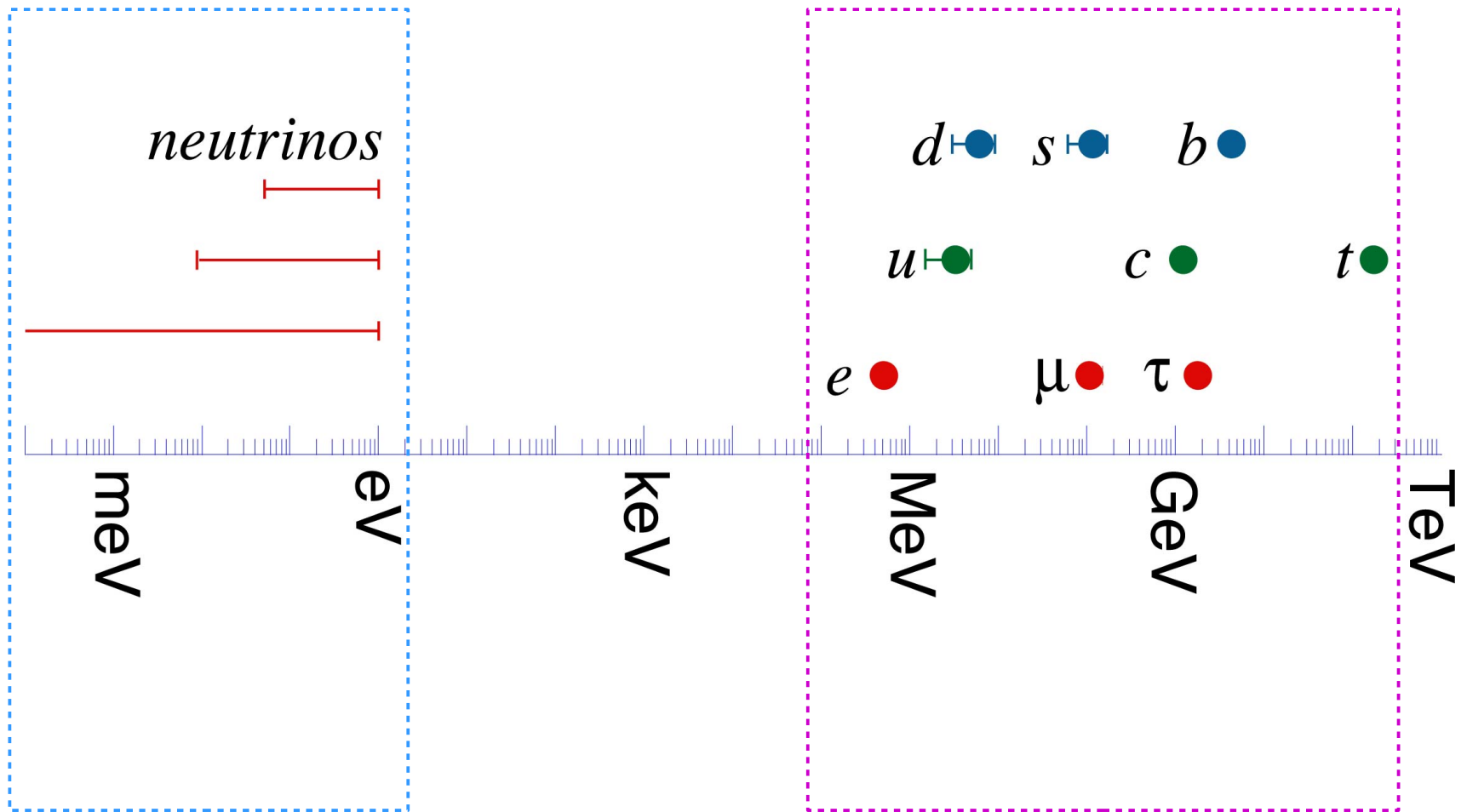
Pilar Coloma  
IFIC, UV/CSIC



PPNT'19, Uppsala  
Oct 8<sup>th</sup> 2019



# Neutrino masses



# Majorana or Dirac?

New fields are required to give neutrinos a mass. Two main ways:

1) Dirac mass: as for the rest of fermions in the SM

$$Y_\nu \bar{L}_L \tilde{\phi} \nu_R \rightarrow m_\nu \bar{\nu}_L \nu_R$$

$$Y_\nu \lesssim 10^{-12}$$

# Majorana or Dirac?

New fields are required to give neutrinos a mass. Two main ways:

1) Dirac mass: as for the rest of fermions in the SM

$$Y_\nu \bar{L}_L \tilde{\phi} \nu_R \rightarrow m_\nu \bar{\nu}_L \nu_R$$

$$Y_\nu \lesssim 10^{-12}$$

2) A Majorana mass. For example:

$$Y_\nu \bar{L}_L \tilde{\phi} \nu_R + \frac{1}{2} M \bar{\nu}_R^c \nu_R$$

$$(\cancel{I})$$

# Majorana or Dirac?

New fields are required to give neutrinos a mass. Two main ways:

1) Dirac mass: as for the rest of fermions in the SM

$$Y_\nu \bar{L}_L \tilde{\phi} \nu_R \rightarrow m_\nu \bar{\nu}_L \nu_R$$

$$Y_\nu \lesssim 10^{-12}$$

2) A Majorana mass. For example:

$$Y_\nu \bar{L}_L \tilde{\phi} \nu_R + \frac{1}{2} M \bar{\nu}_R^c \nu_R$$

( $\cancel{L}$ )

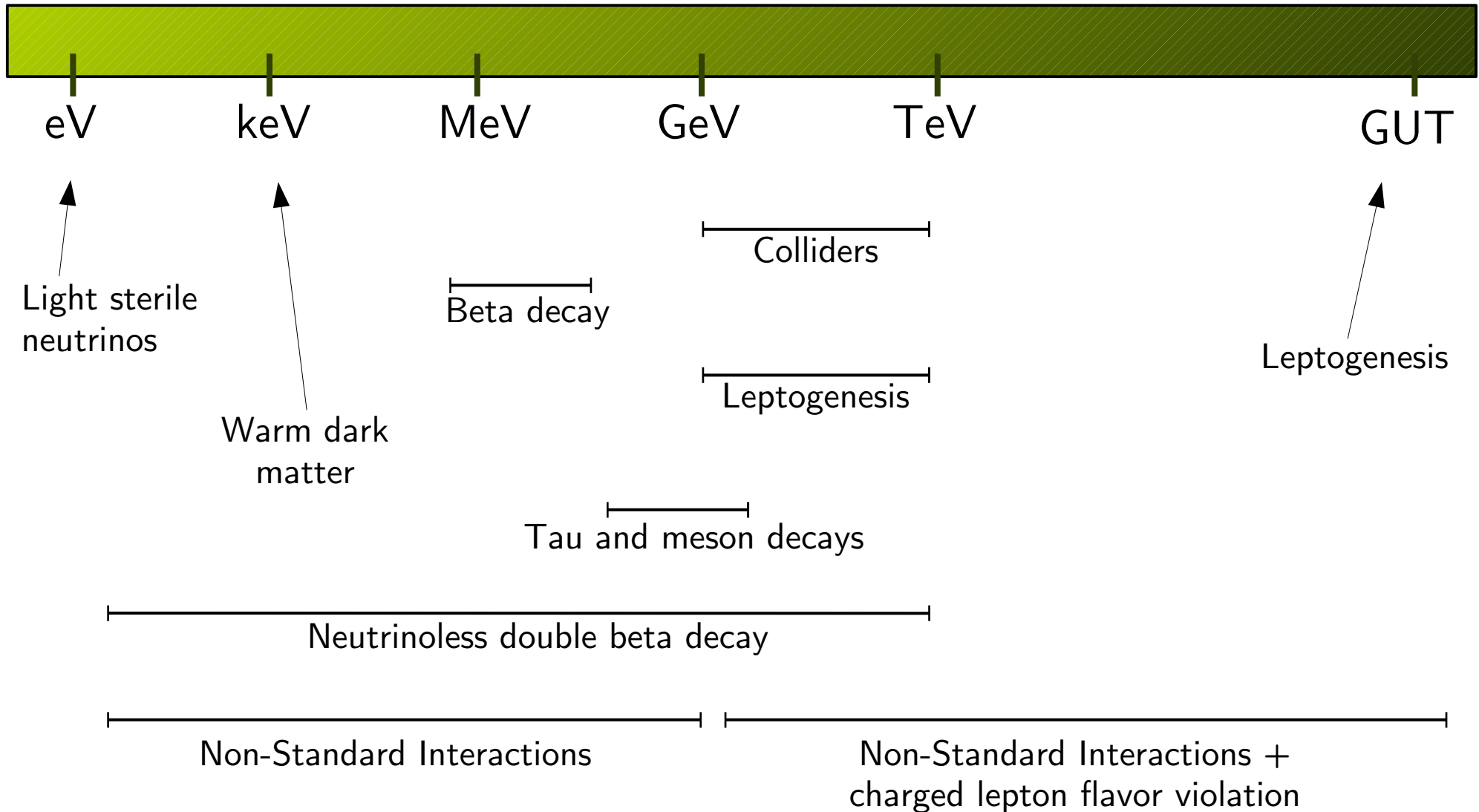
$$M \gg Y_\nu v$$

$$\rightarrow m_\nu = Y_\nu^\dagger M^{-1} Y_\nu v^2$$

Type I Seesaw:

Minkowski '77, Gell-Mann, Ramond, Slansky '79, Yanagida '79, Mohapatra, Senjanovic '80

# Scale of new physics



# Why the GeV scale?

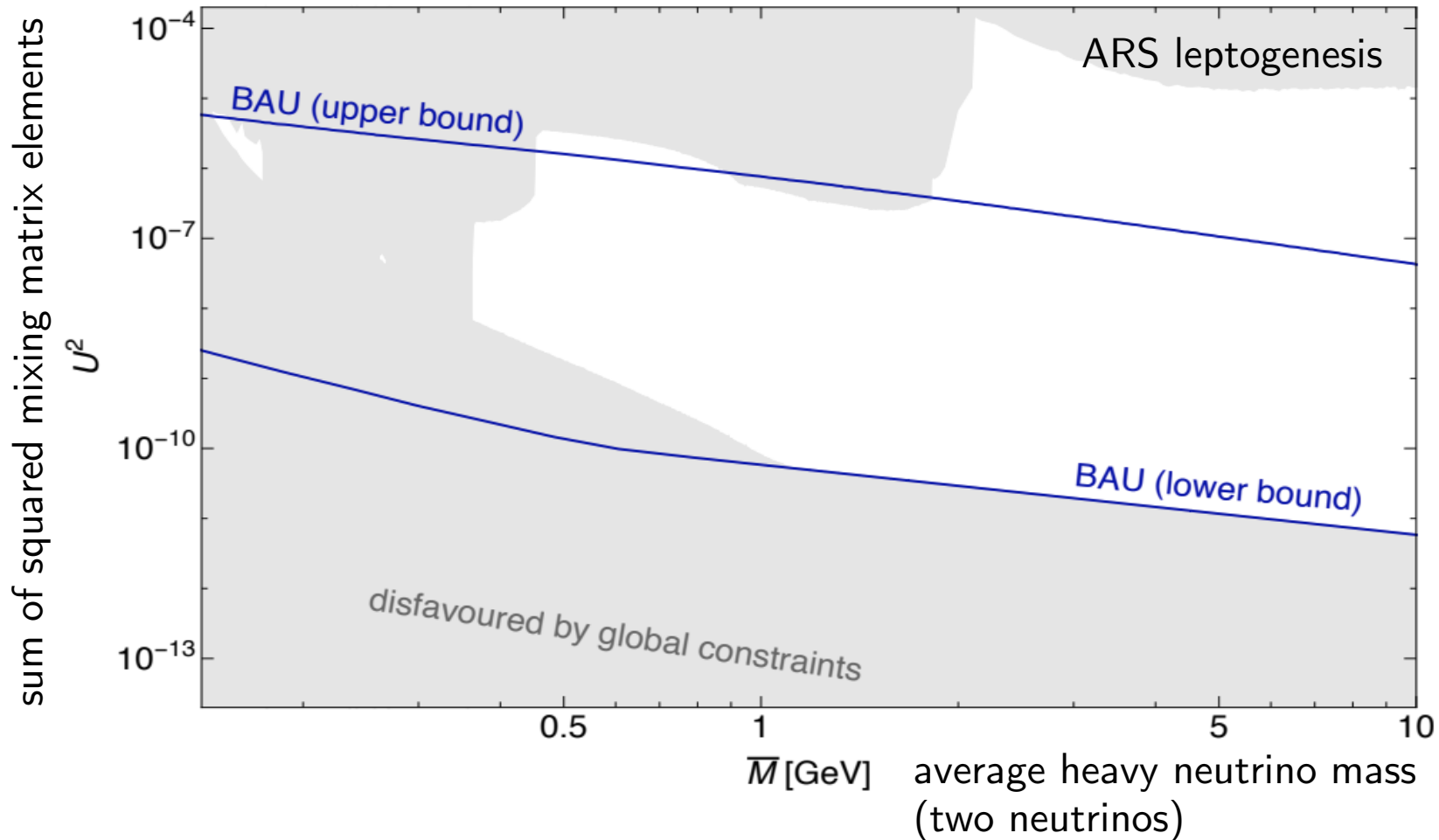
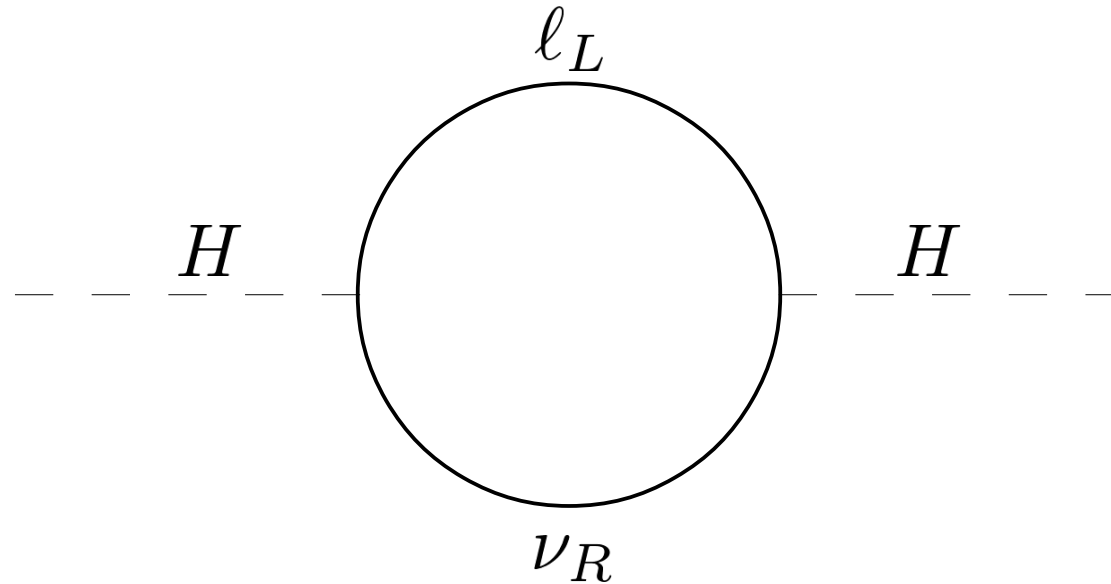


Figure from the review: Drewes, Garbretch, Hernandez, Kekic, Lopez-Pavon, Racker, Rius, Salvado, Teresi, 1711.02862

# Why a low-scale seesaw?



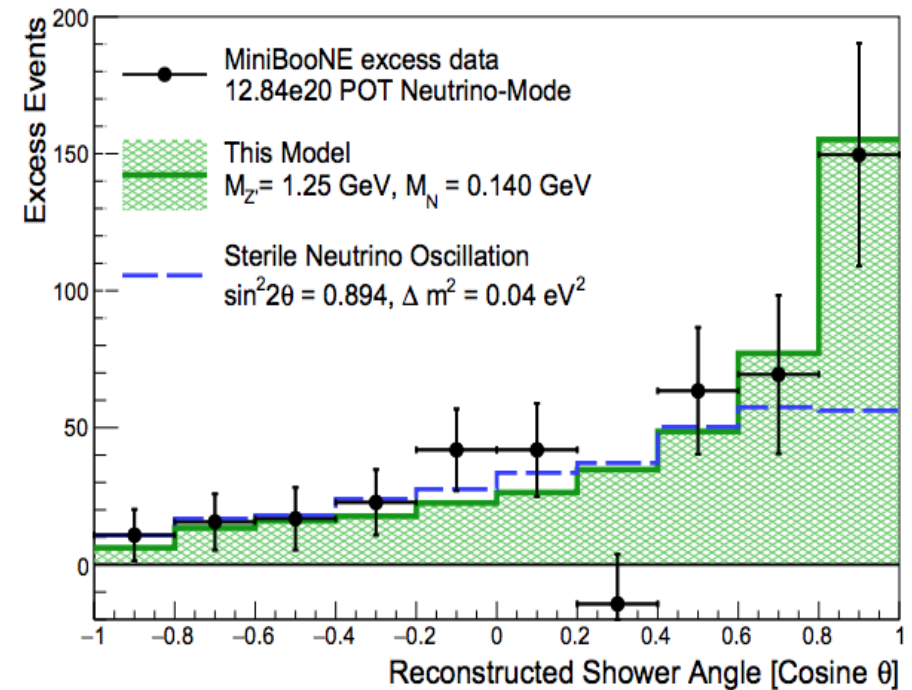
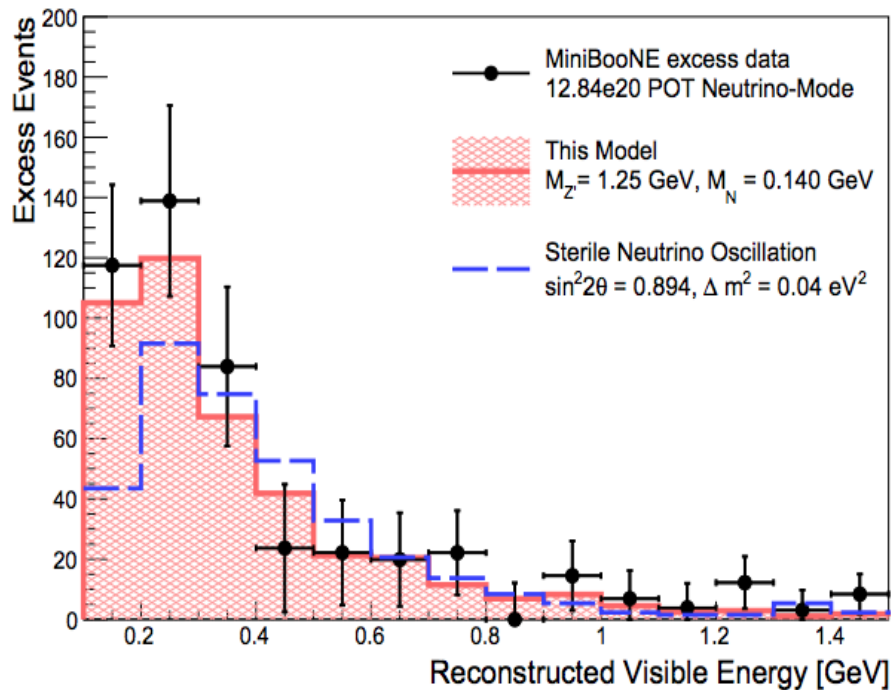
$$\delta\mu^2 \approx \frac{Y_\nu^2}{4\pi^2} M_R^2 \log\left(\frac{q}{M_R}\right)$$

Vissani, hep-ph/9709409



# Further motivation...?

Figure from Ballett, Pascoli, Ross-Lonergan, 1808.02915



See also:

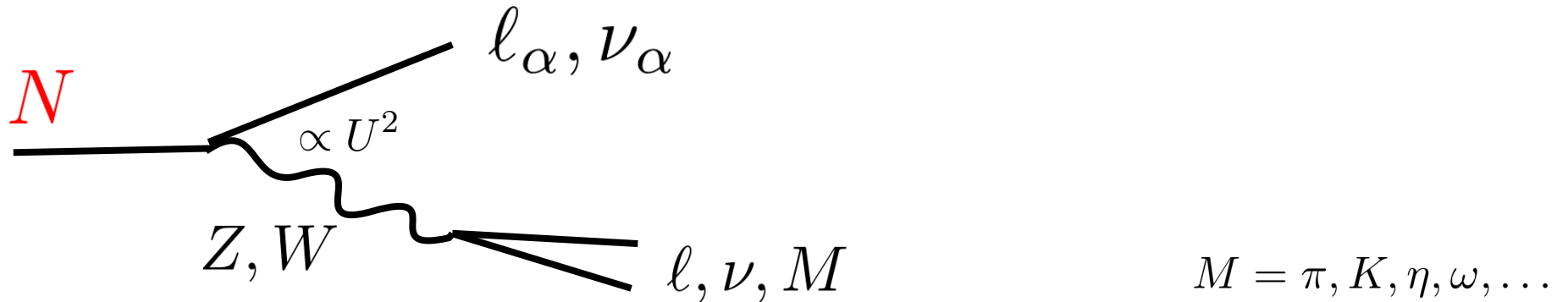
Bertuzzo, Jana, Machado, Zukanovich Funchal, 1807.09877

Fischer, Hernandez-Cabezudo, Schwetz, 1909.09561

Gninenko, 1009.5536 and 0902.3802

Palomares-Ruiz, Pascoli, Schwetz, hep-ph0505216

# Direct searches for GeV neutrinos



- Direct searches can be divided into two main categories:
  - Peak searches
  - Displaced decays: fixed target experiments, colliders, ...

$$c\tau \sim \text{few} \left( \frac{\text{GeV}}{m_N} \right)^5 \left( \frac{10^{-4}}{U^2} \right) \text{ m}$$

For detailed calculations of heavy neutrino decay channels see, e.g.: Ballett, Boschi, Pascoli, 1905.00284; Bondarenko et al, 1805.08567

# Direct searches for GeV neutrinos

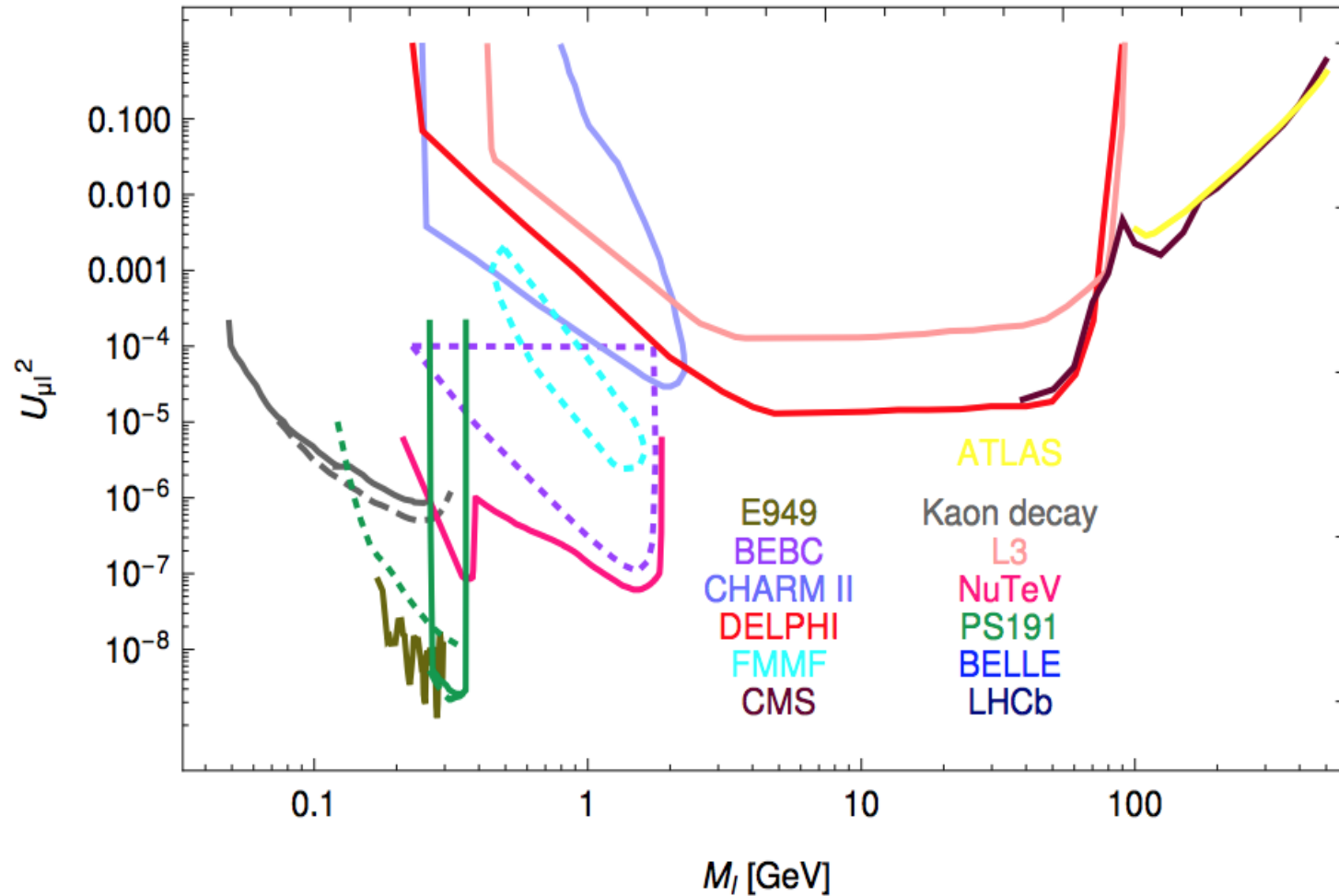


Figure from Drewes and Garbrecht, 1502.00477

(See also Ruchayskiy and Ivashko, 1112.3319, Atre et al, 0901.3589)

# Direct searches for GeV neutrinos

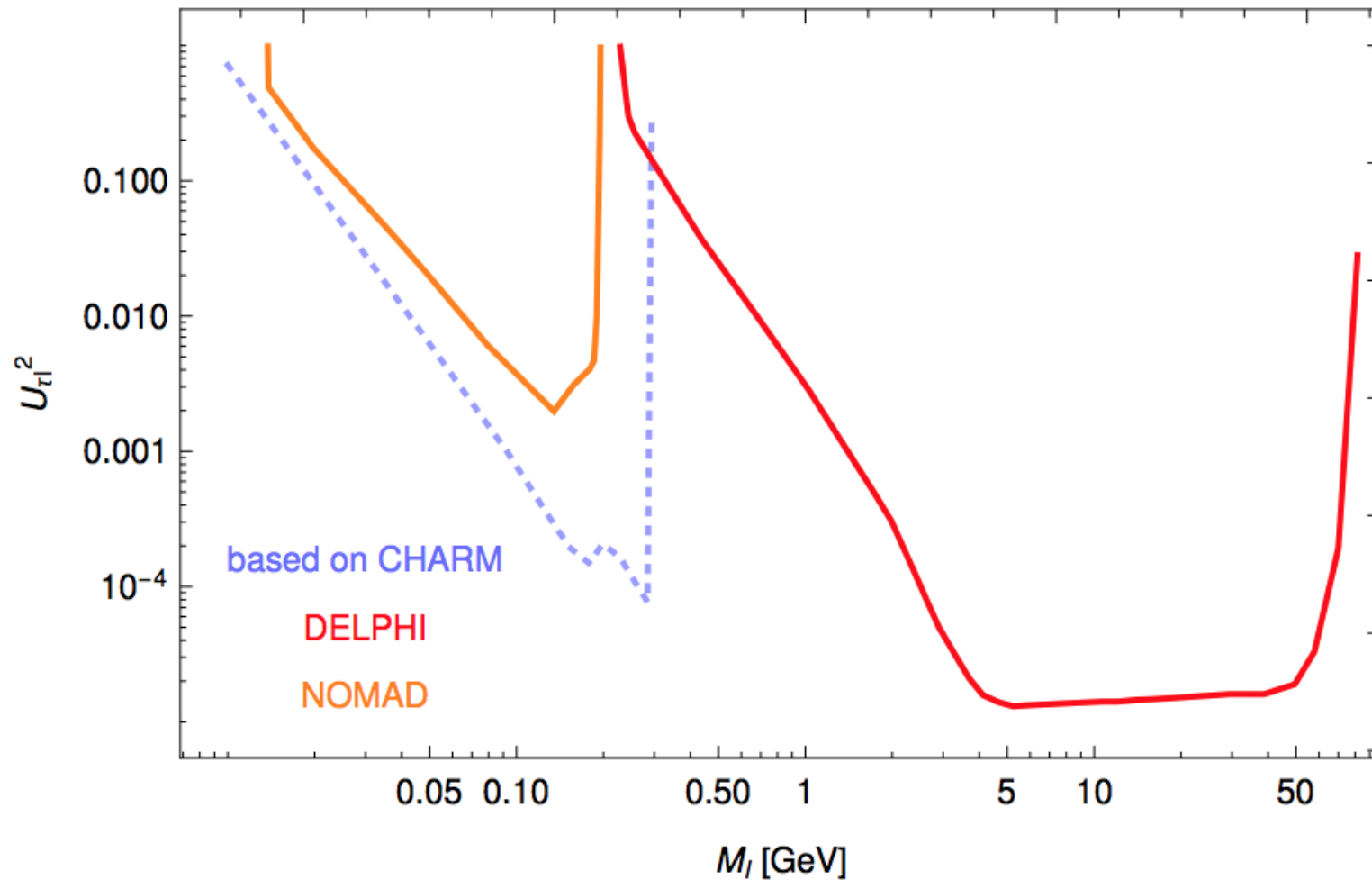
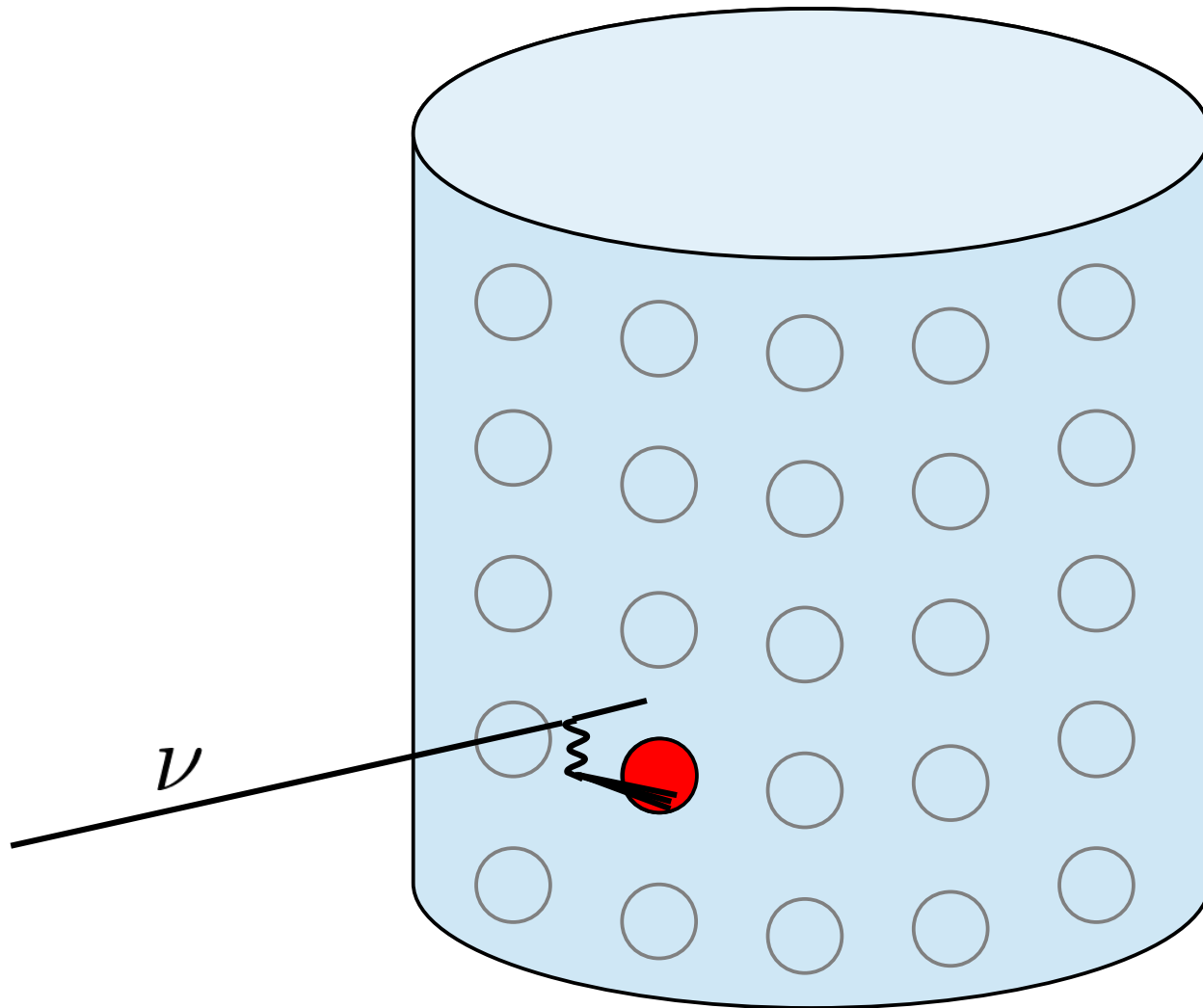


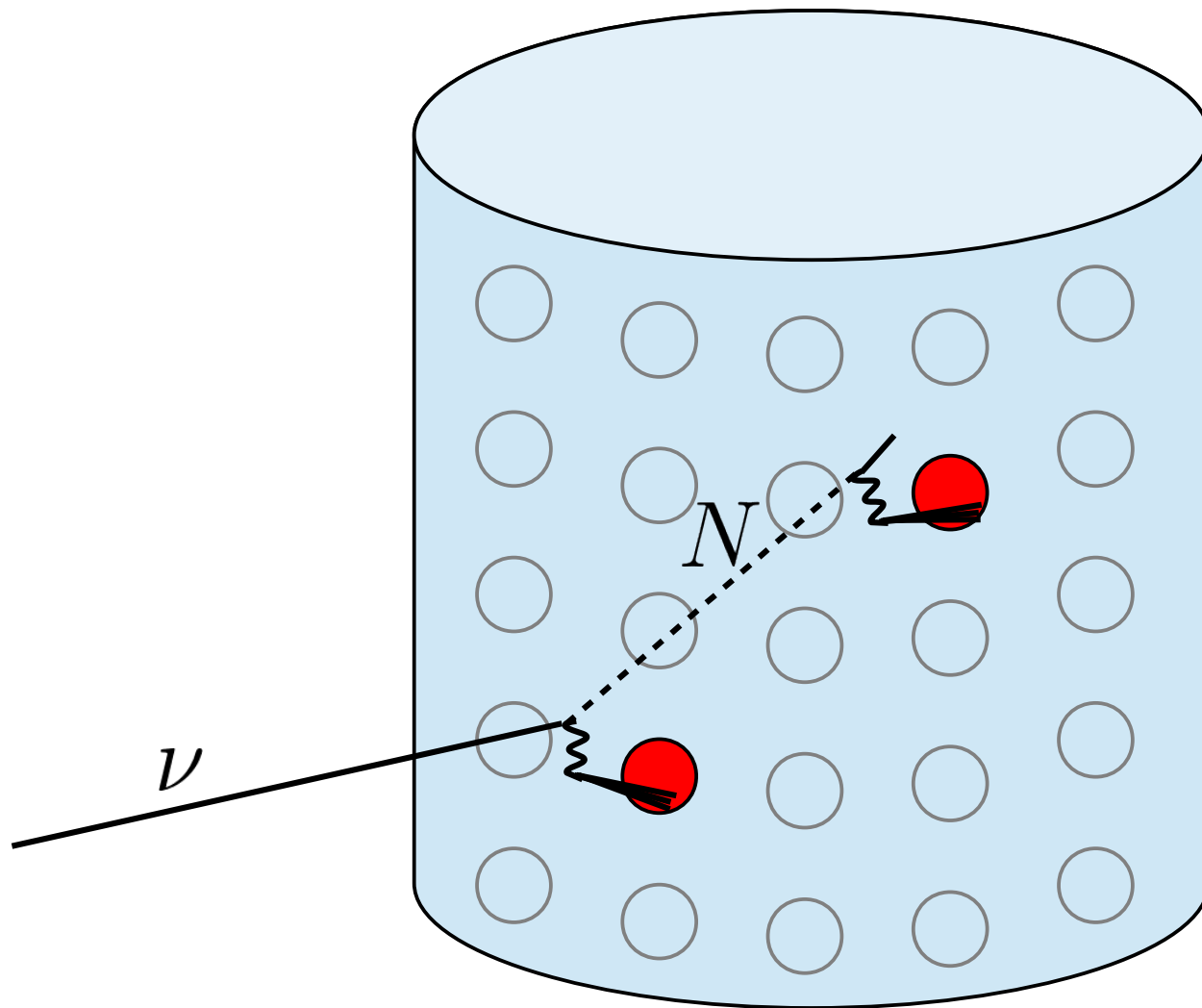
Figure from Drewes and Garbrecht, 1502.00477

(See also Ruchayskiy and Ivashko, 1112.3319, Atre et al, 0901.3589)

# Heavy neutrinos produced inside the detector







# Double-bangs are expected at UHE

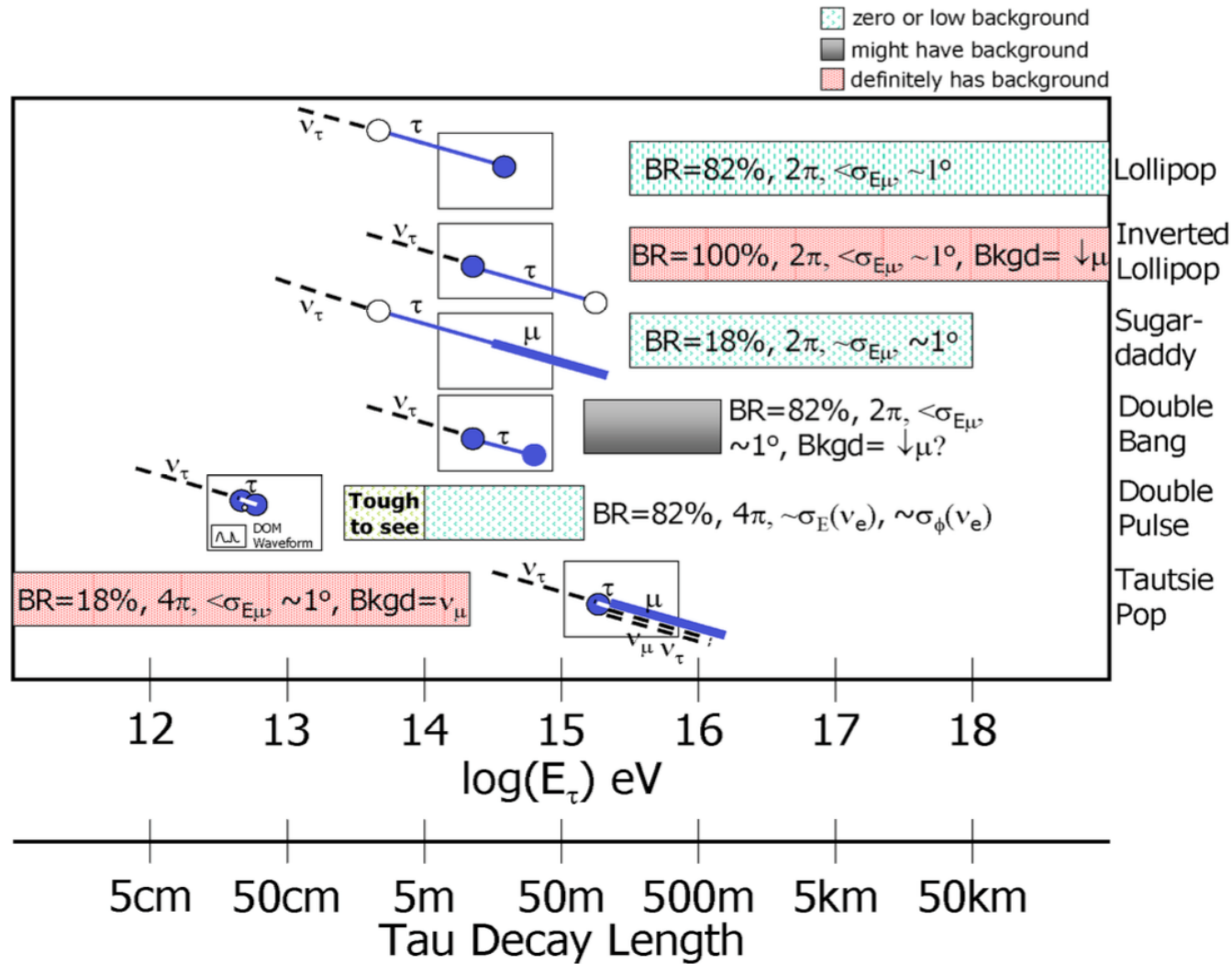
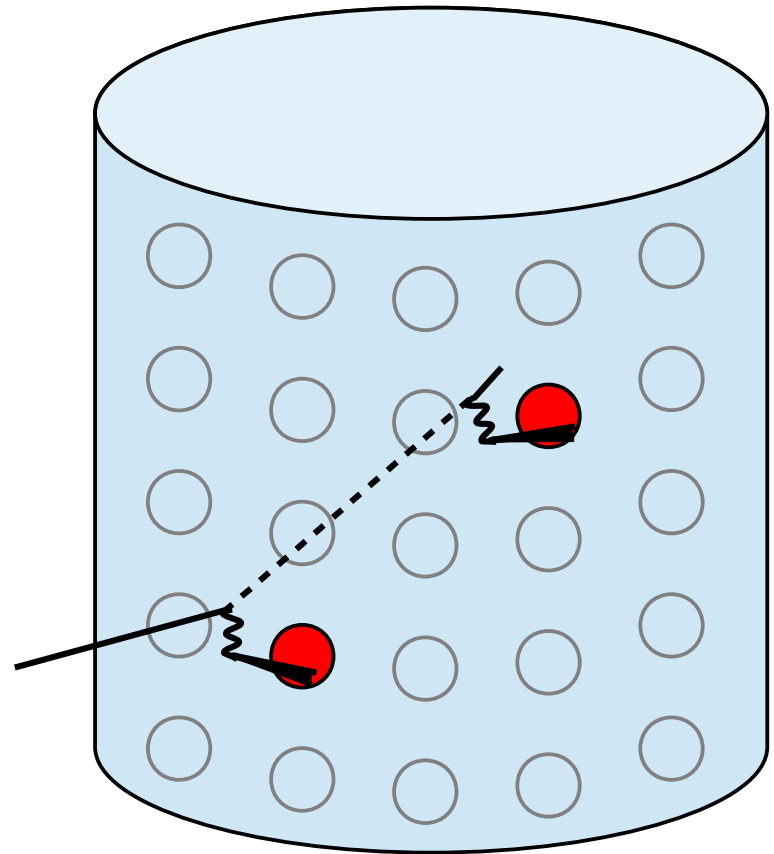


Figure from: Tau Neutrinos in IceCube, D. F. Cowen, TeVPA'06 proceedings

See also Icecube coll., 1509.06212

# Low-energy double bangs?

- Key requirements:
  - Trigger has to go off during first shower: 3-4 DOMs hit
  - Minimum energy/distance to reach a DOM (limited by ice absorption): 36 m
  - Minimum separation between the two showers (limited by time resolution, 20 m)



# Signal vs background

$$\frac{dN_{ev}}{dt} \simeq \rho_{ice} BR_{vis} \int dE_\nu dE_N d\cos\theta \frac{d\phi_{\nu_\mu}}{dE_\nu d\cos\theta} P_{\nu_\mu \rightarrow \nu_\tau} \frac{d\sigma_{\nu_\tau N}}{dE_N} V_{det}(L_{lab}, \cos\theta)$$

Effective volume (depends on the neutrino decay length and detector geometry)

Order-of-magnitude estimates **for the signal**, for 1 GeV heavy neutrino:

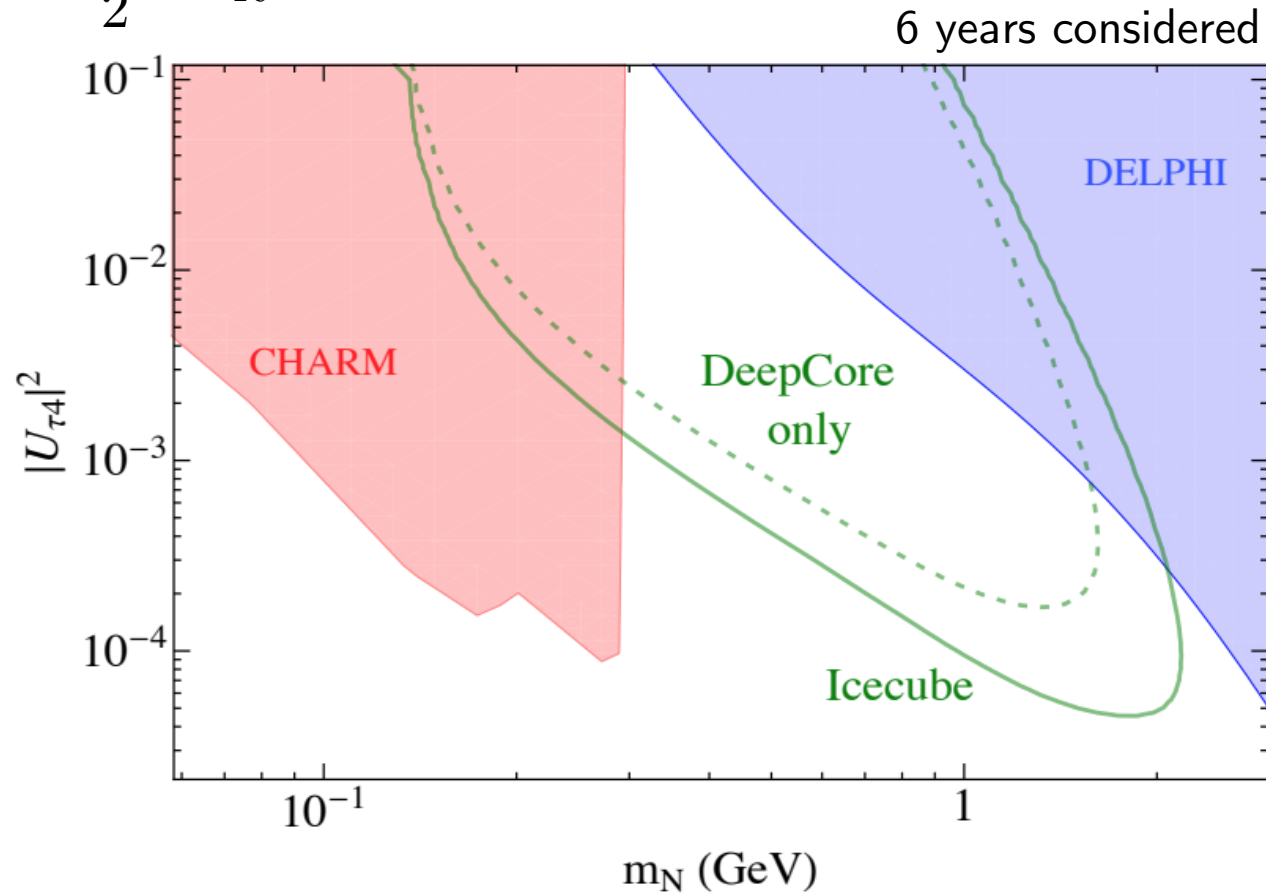
$$N_{sig} \sim \mathcal{O} \left( \frac{|U|^2}{10^{-4}} \right) \text{ events/yr}$$

Order-of-magnitude estimates **for the background**:

$$N_{bg} \sim \mathcal{O} (0.05) \text{ events/yr}$$

# Vanilla scenario: only mixing

$$Y_\nu \bar{L}_L \tilde{\phi} \nu_R + \frac{1}{2} M \bar{\nu}_R^c \nu_R$$

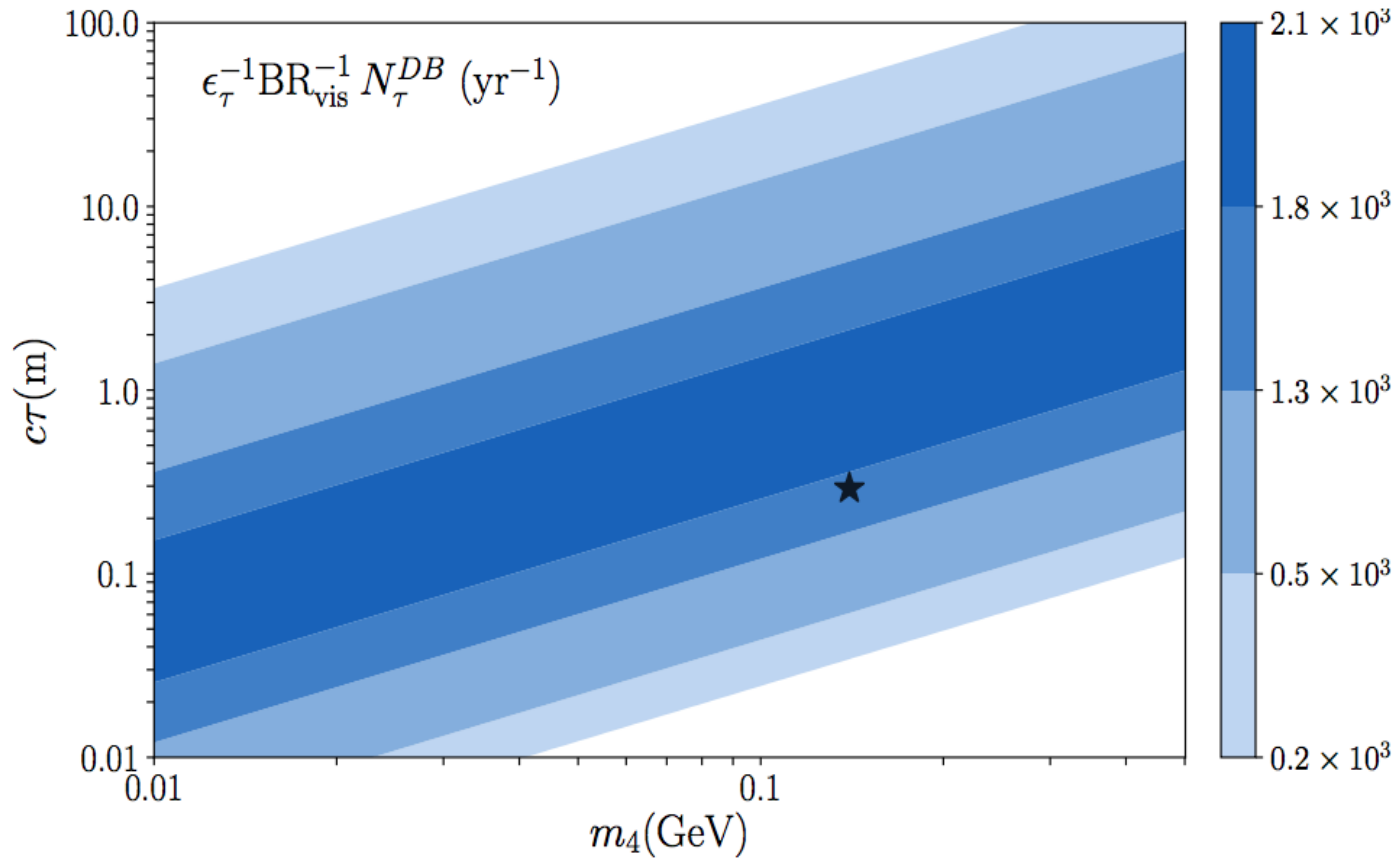


Coloma, Machado, Martinez-Soler and Shoemaker, 1707.08573

# Light $Z'$ boson + mixing

$$\dots + U_{\alpha 4} g' \bar{\nu}_\alpha \gamma^\mu P_L \nu_4 Z'_\mu$$

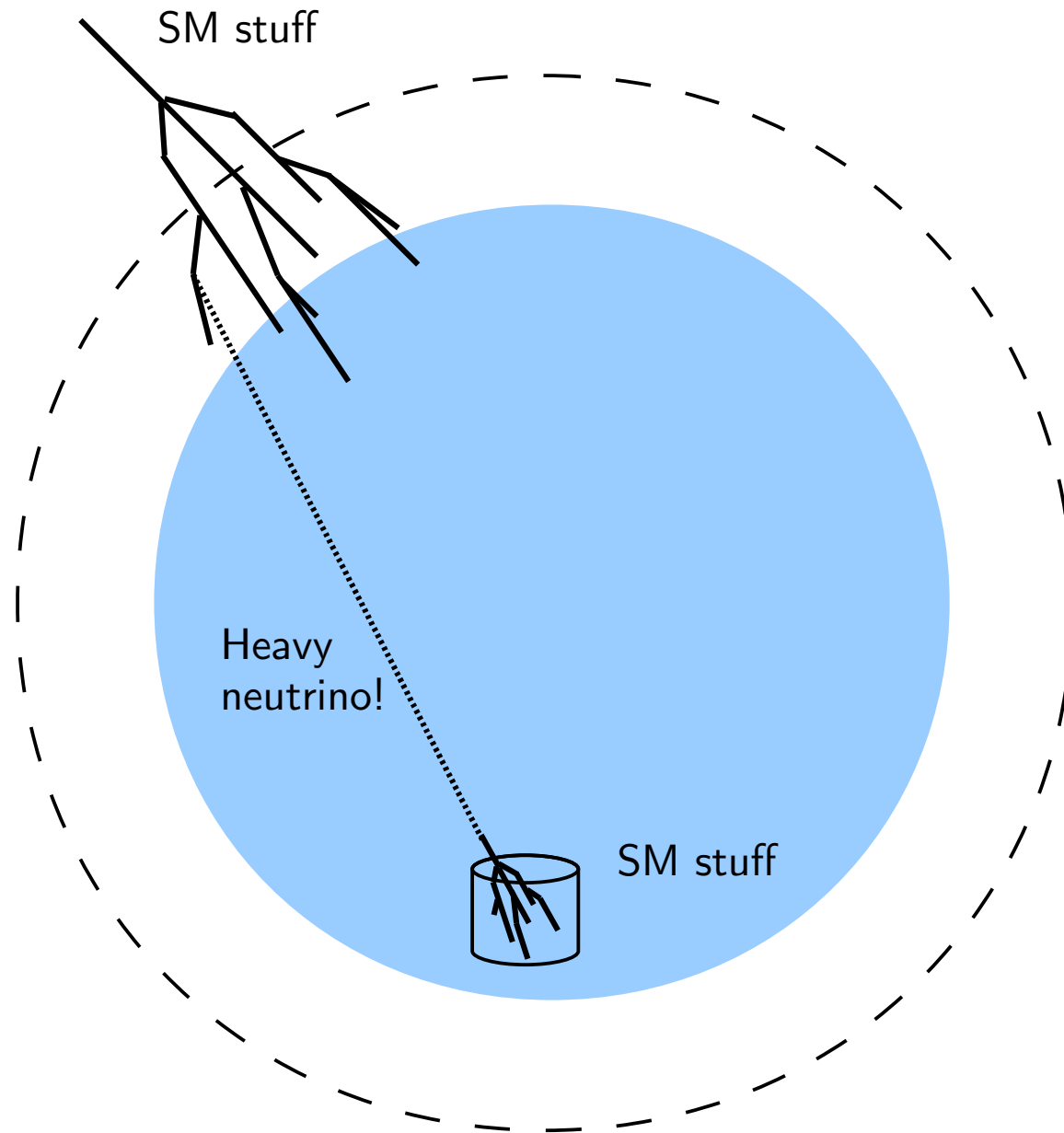
$$N_{\alpha, Z'} = N_{\alpha, Z} \epsilon_\alpha$$



Coloma, 1906.02106



Heavy neutrinos produced  
in the atmosphere



# Signal computation

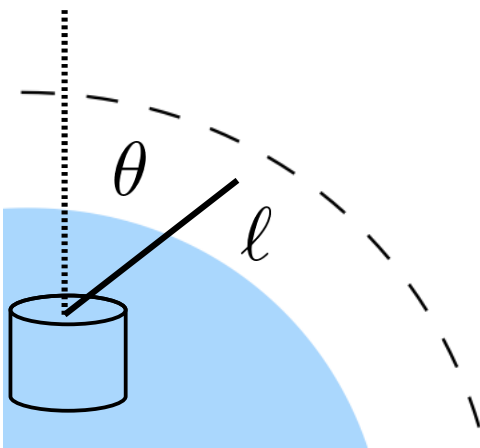
$$\frac{d\Phi_N}{dE_N d\cos\theta} = \int_0^{\ell_{\max}} d\ell \left[ \frac{d\Pi_N}{dE_N d\cos\theta d\ell} \right] e^{-\frac{\ell}{\ell_{\text{decay}}}}$$

Computed as in the SM, but considering GeV-scale neutrinos

$$N_i^\alpha = \text{Br}(e\text{-like}) \int_{E_N} \left[ \epsilon^{\alpha\beta}(E_i, E_N) \right] \int_{\theta} \left[ A_{\text{decay}}^{\text{eff}}(E_N, \cos\theta) \right] \frac{d\Phi_N}{dE_N d\cos\theta}$$

Detector efficiencies  
(estimated)

Geometric acceptance  
(geometry)

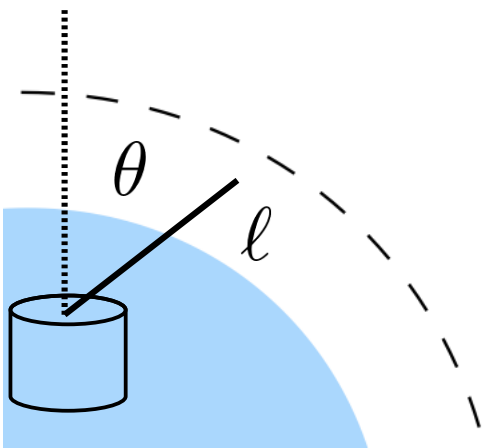


Arguelles, Coloma, Hernandez, Muñoz, 1910.XXXXXX

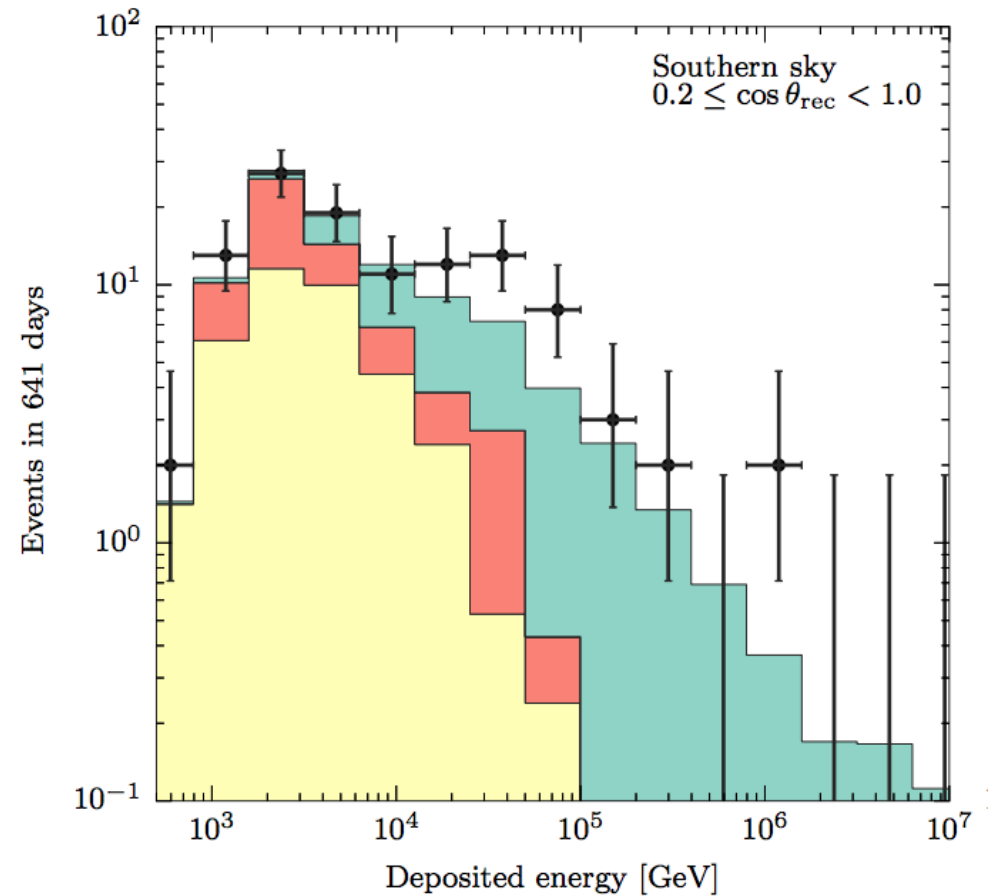
# Backgrounds and data

Data release from 1410.1749 used  
(MESE data, 641 days)

Only down-going events considered  
(signal region)

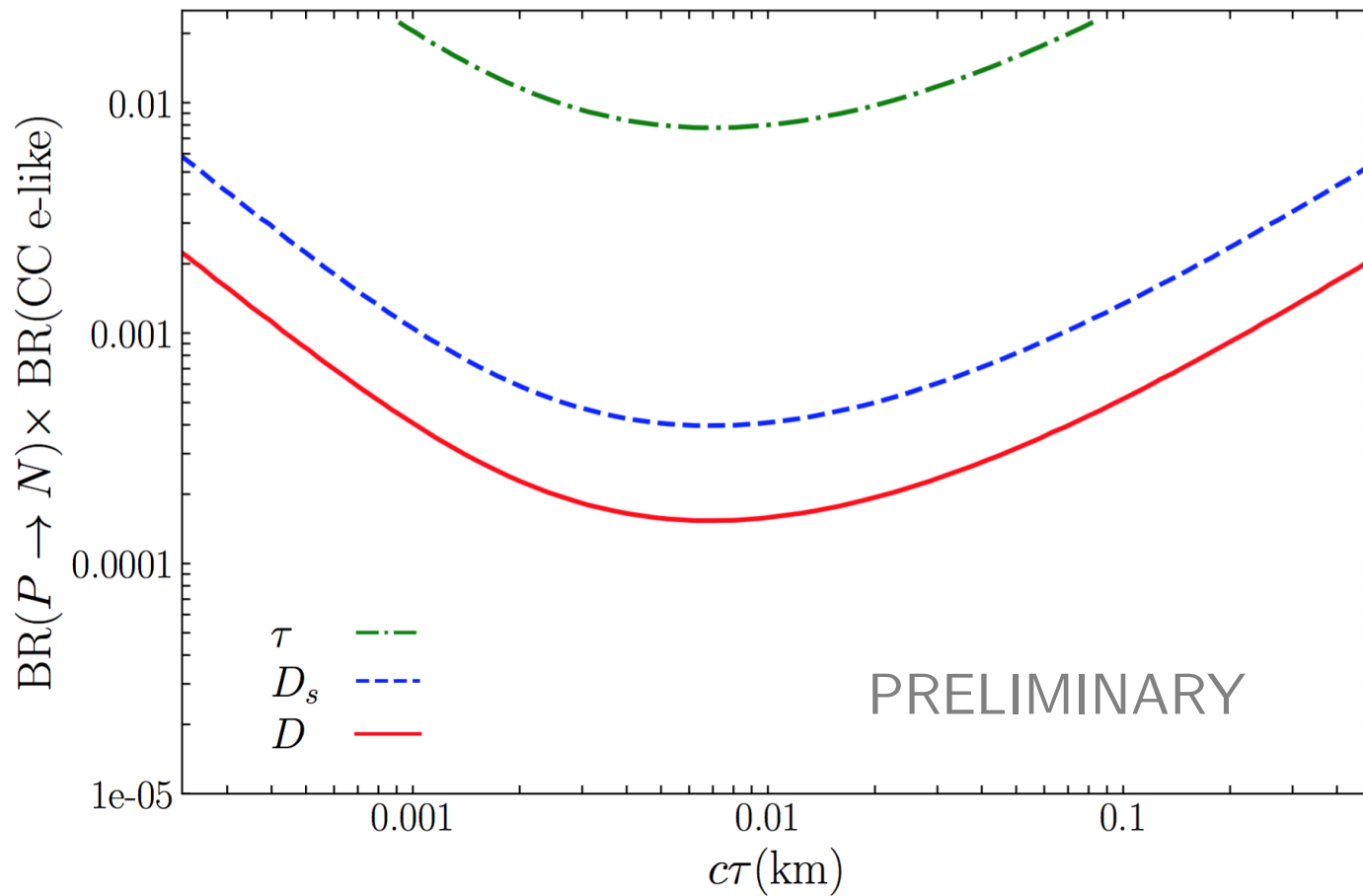


Pilar Coloma - IFIC



Icecube collaboration, 1410.1749

# Exclusion regions



Arguelles, Coloma, Hernandez, Muñoz, 1910.XXXXX

# Summary

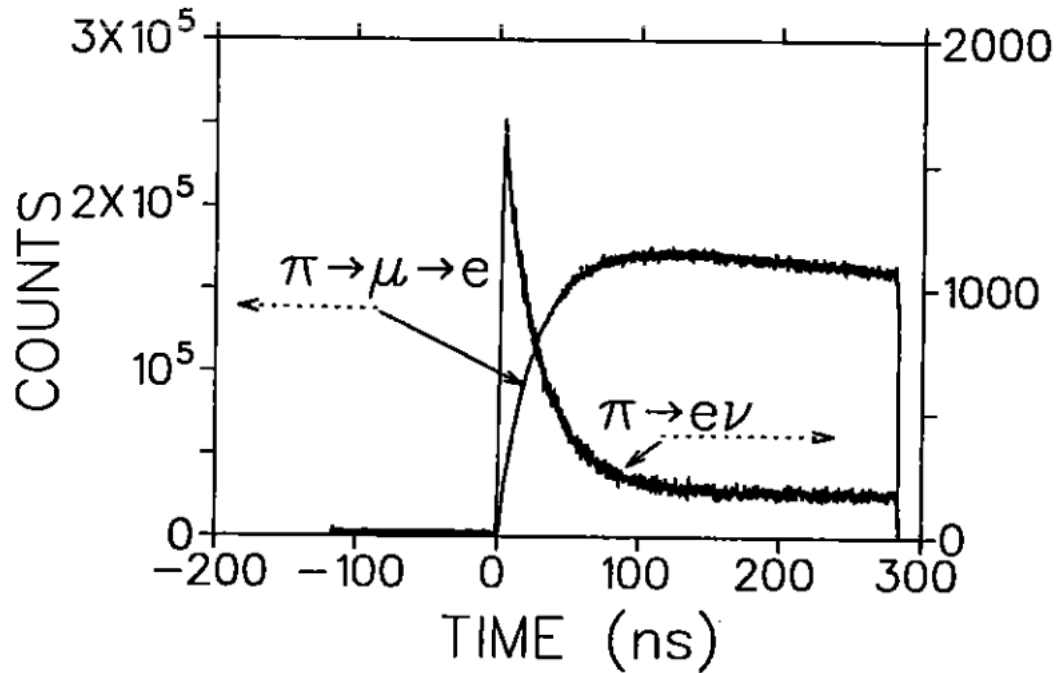
- GeV neutral leptons are a well-motivated extension of the SM, with a rich phenomenology.
- Neutrino telescopes are sensitive to heavy neutrinos produced in several ways:
  - produced and decayed **inside the detector** (double-bangs at low energies)
  - produced **in the atmosphere**, decayed in the detector



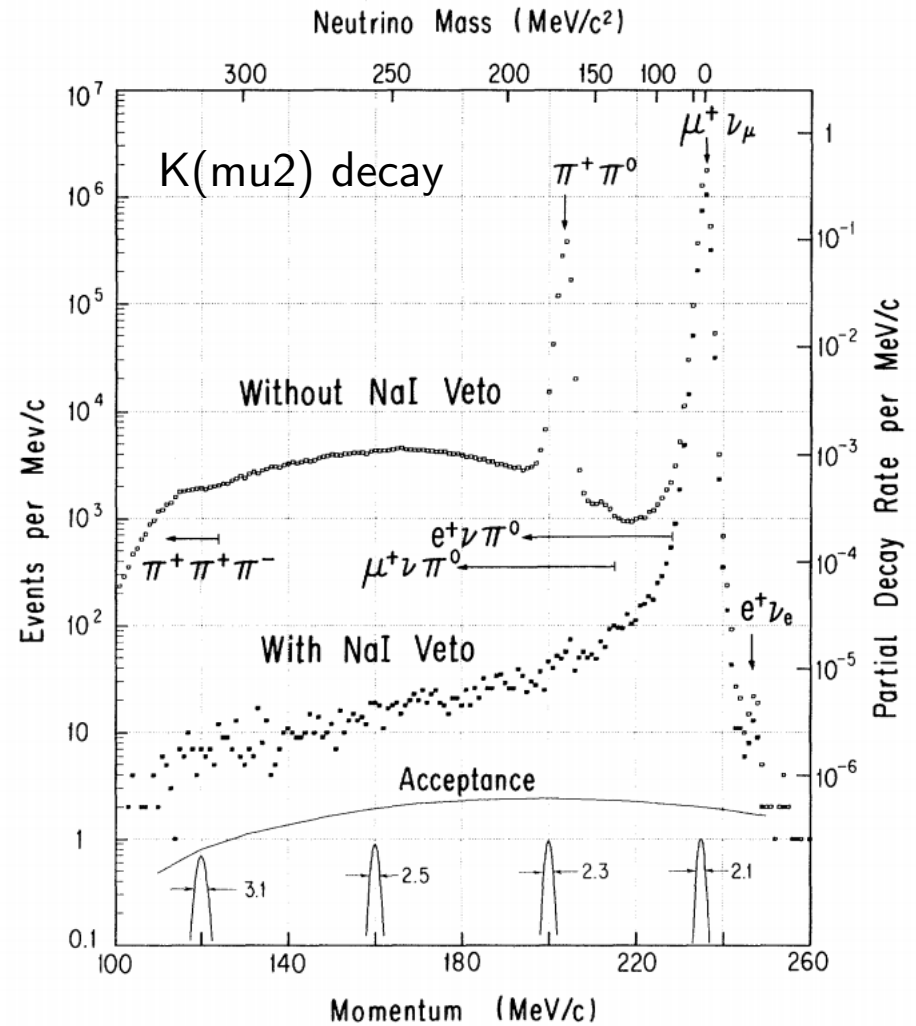
Thank you!!

Backup slides

# Peak searches



Britton et al (TRIUMF),  
 Phys.Rev.Lett. 68 (1992) 3000-3003



Hayano et al, Phys.Rev.Lett. 49 (1982) 1305

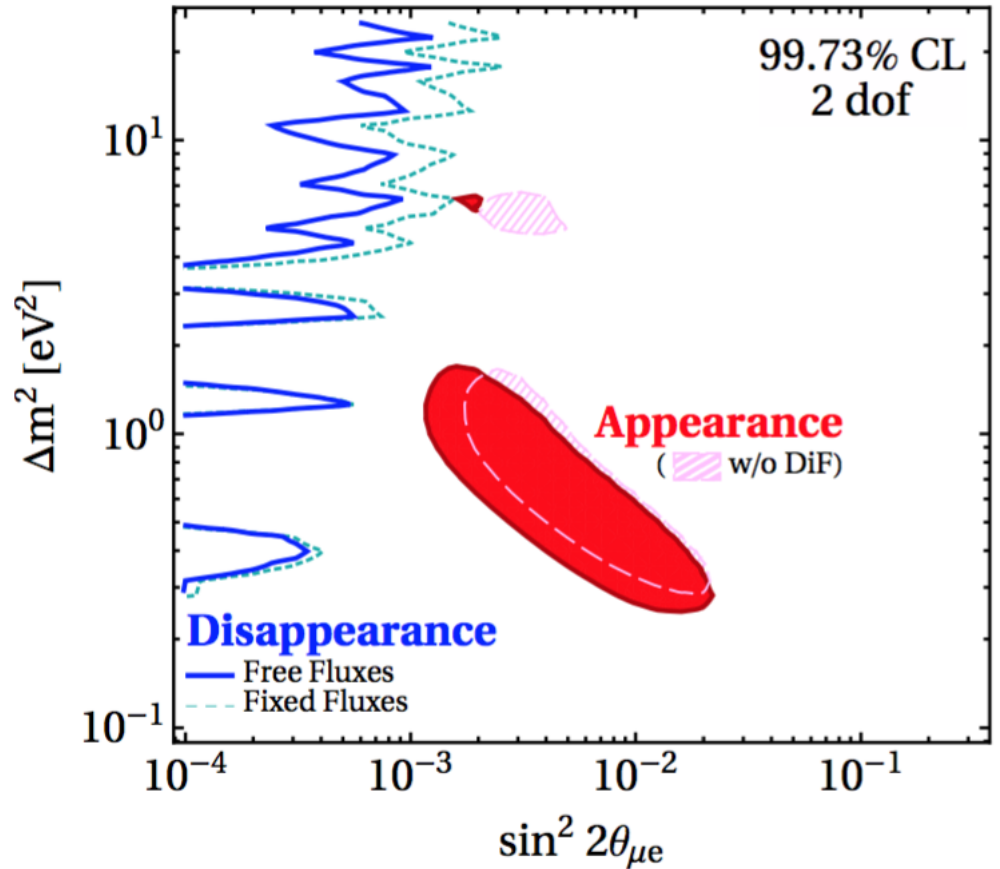
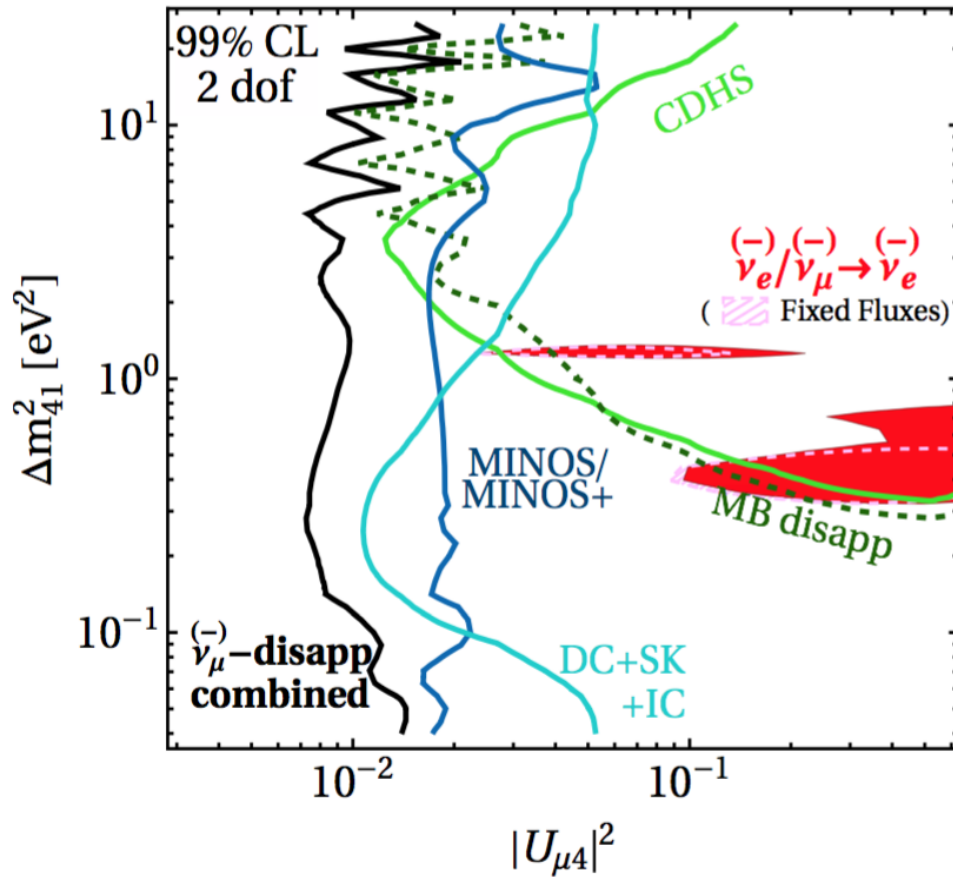
# How to search for GeV neutral leptons

- At low energies, we can describe the interactions with mesons through the effective Lagrangian:

$$\begin{aligned}
 \mathcal{L}_{chiral+N} &= \sum_{P^0} G_F V_{\alpha 4}^* \mathcal{F}_{P^0} \bar{N}_R^c p_\mu P_L \nu_\alpha \\
 &+ \sum_{V^0} G_F V_{\alpha 4}^* \mathcal{F}_{V^0} \bar{N}_R^c \epsilon_\mu P_L \nu_\alpha \\
 &+ \sum_{P^\pm} \sqrt{2} G_F V_{\alpha 4}^* \mathcal{F}_{P^\pm} \bar{N}_R^c p_\mu P_L \ell_\alpha \\
 &+ \sum_{V^\pm} \sqrt{2} G_F V_{\alpha 4}^* \mathcal{F}_{V^\pm} \bar{N}_R^c \epsilon_\mu P_L \ell_\alpha + \text{h.c.}
 \end{aligned}$$

$$P = \pi, K, \dots \quad V = \eta, \omega, \dots$$

# Tensions in a 3+1 scenario



Dentler et al, 1803.10661  
(see also Giunti and Lasserre, 1901.08330)