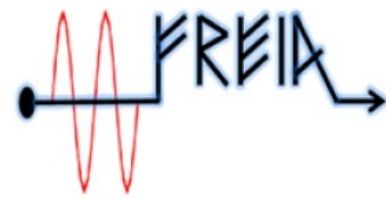




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# Accelerators and Magnets

Volker Ziemann, FREIA



<https://ziemann.web.cern.ch/ziemann/gadget/freia/>

2011



210421, V. Ziemann  
<https://cern.ch/ziemann>

Accelerators and Magnets

Today



# Roadmap

- Accelerators
  - key components
- FREIA
- Iron-dominated magnets
- Current-dominated magnets
- (Permanent magnets)
- (Measurements of magnets)



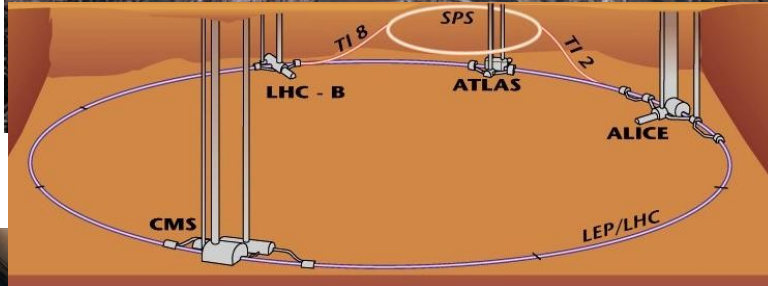


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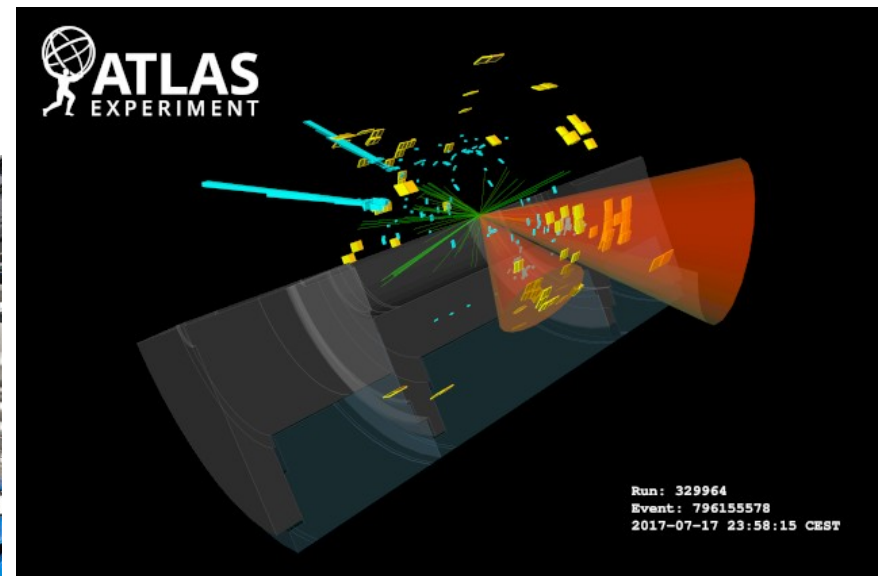
# LHC at CERN



[https://www.youtube.com/watch?v=sB0\\_ohLM3Kg](https://www.youtube.com/watch?v=sB0_ohLM3Kg)



All images  
from CERN



Higgs  $\rightarrow$  bb (jets) +  $\gamma\gamma$

Run: 329964  
Event: 796155578  
2017-07-17 23:58:15 CEST

Accelerating and bending

210421, V. Ziemann



Accelerators and Magnets

LHC does high-energy particle physics



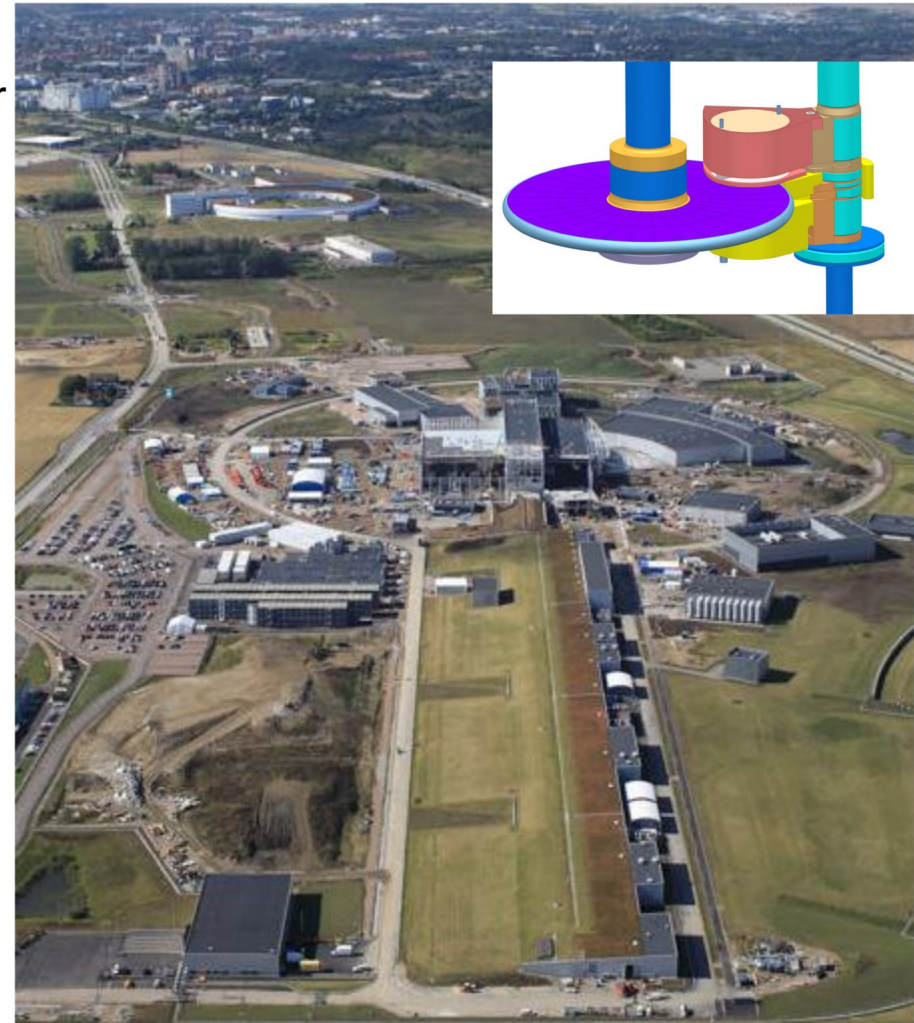
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# MAX IV and ESS



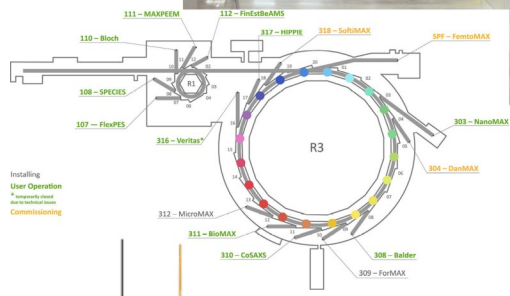
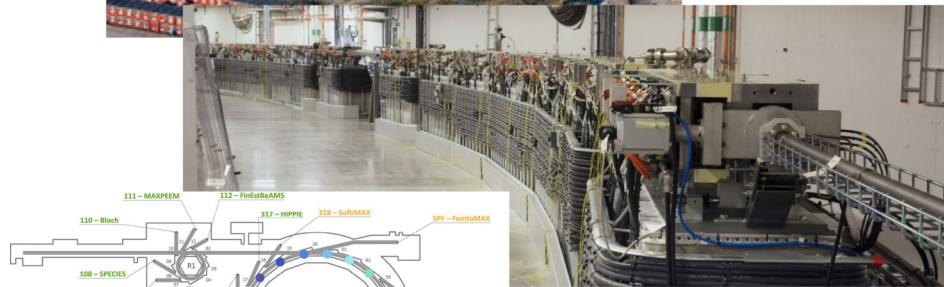
MAX IV:  
electron ring

ESS:  
proton linear  
accelerator

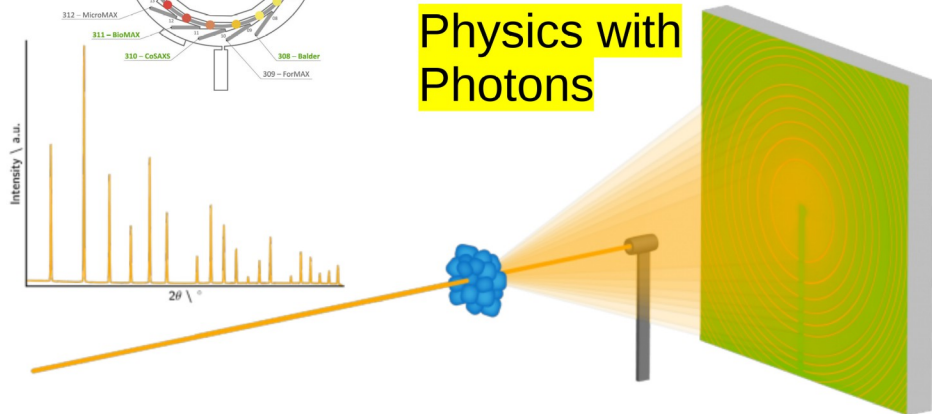


Material and life sciences

Physics with neutrons  
- neutral particles  
- magnetic moment



Physics with  
Photons



210421, V. Ziemann

Accelerators and Magnets

More magnets

More acceleration structures



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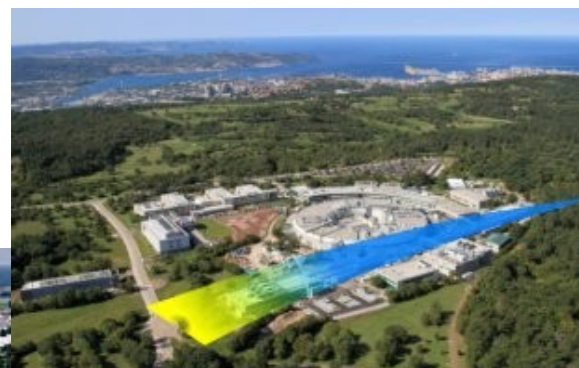
# Accelerators in Europe

...a small selection...



XFEL

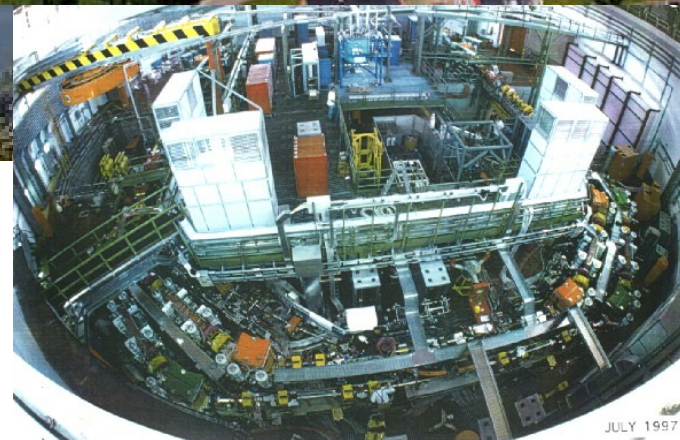
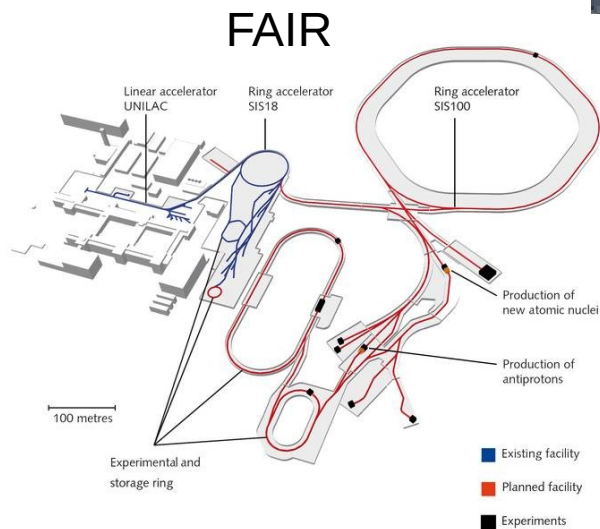
Elettra  
FERMI



Soleil



ESRF



DAFNE

JULY 1997



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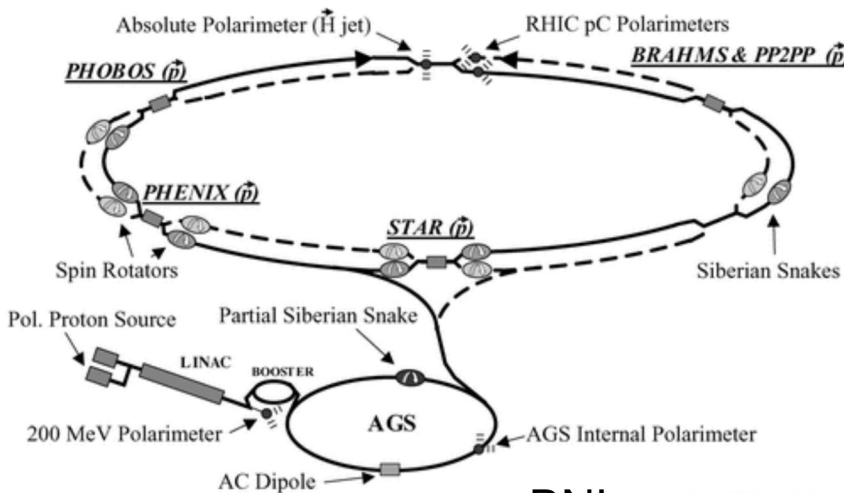
# and the rest of the world



Fermilab



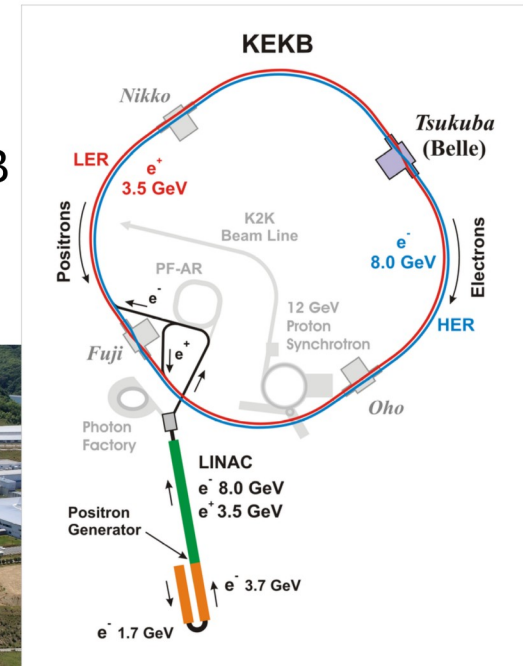
SLAC



BNL



SACLA



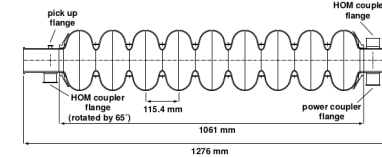
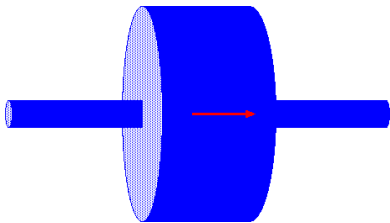
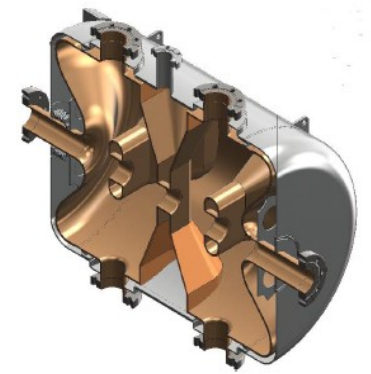
KEKB



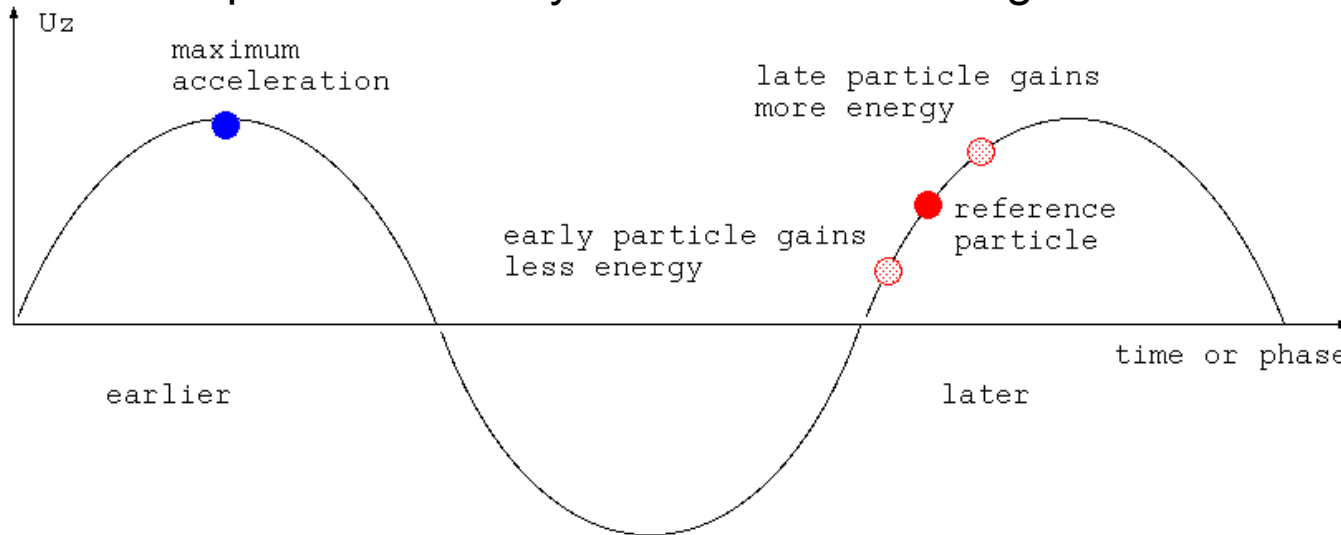
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# Acceleration

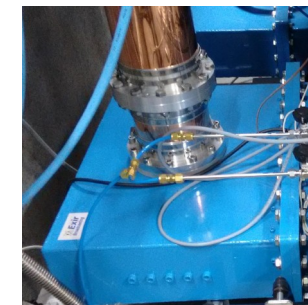
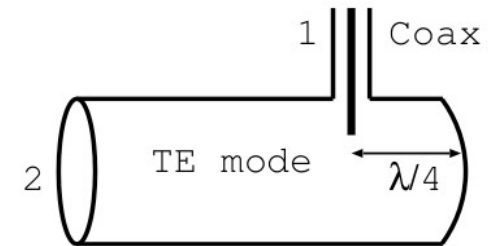
through standing E&M waves, called RF  
electrical field accelerates



Time the particles so they “see” an accelerating field



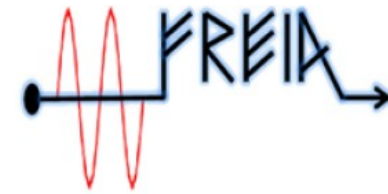
Antenna = Coupler



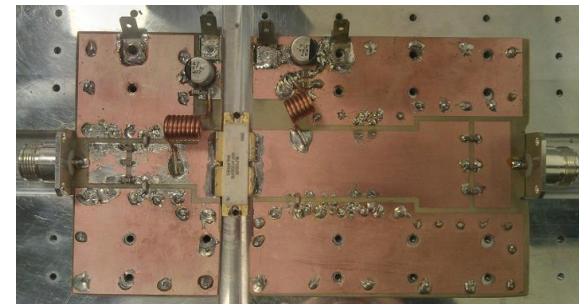
There are normal and superconducting cavities



# Power generation



- Transform “wall-plug power” to radio-frequency power
- Power supplies
  - DC or pulsed (“modulators”)
- Amplifiers
  - Preamplifiers are often based on transistors
  - Tetrodes (old TV and radio vacuum tubes)
  - Klystrons (other types of tubes)



images: FREIA







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# FREIA RF-distribution



Third power station  
with tetrode in the  
background (now  
dismanteled)

Rectangular  
and coaxial  
waveguides



Circulator och  
dummy load  
(resistor)

RF-power  
switchyard

Two RF power stationer  
with two tetrodes each  
400 kW in 3 ms at 17 Hz

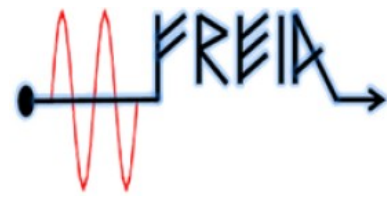
Power  
combiner



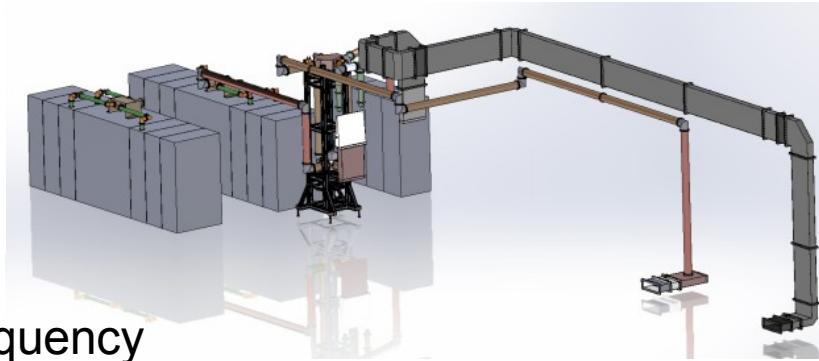
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# What do we do in FREIA?

## Test of Spoke-cavities for ESS



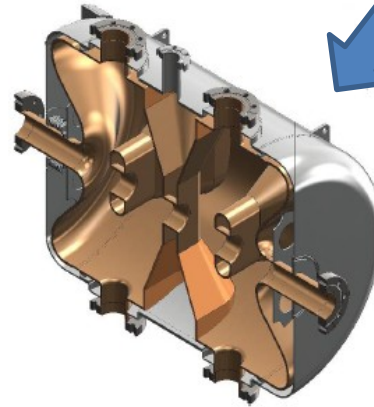
Radio-frequency  
power amplifier



Cryogenics for LHe

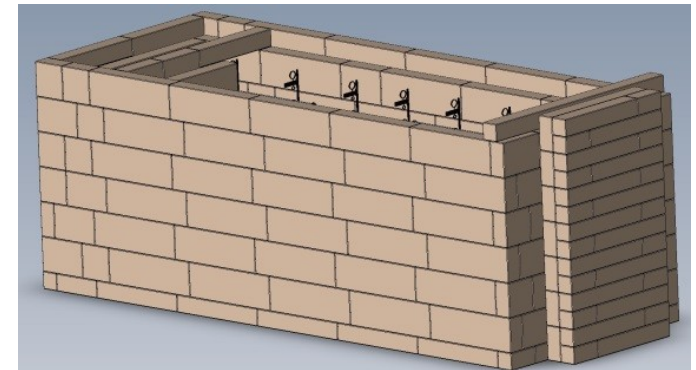


Cryostat



Spoke-cavity  
superconducting

+Diagnostics+Control+...

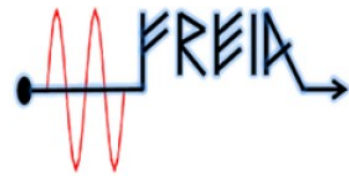


Bunker for radiation protection



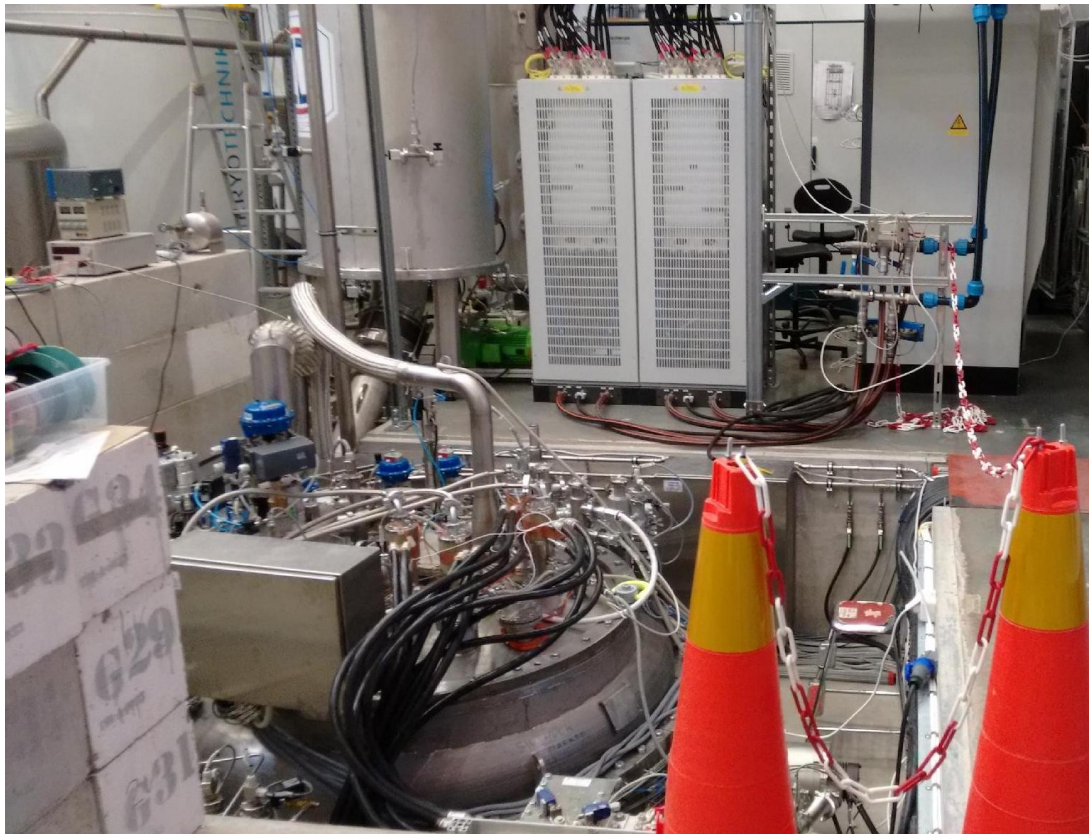
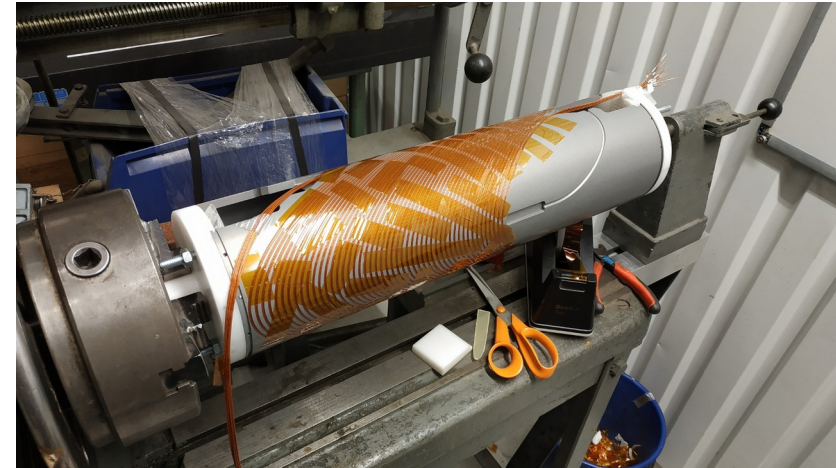
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# FREIA and Magnets



Test of magnets in FREIA's vertical cryostat GERSEMI

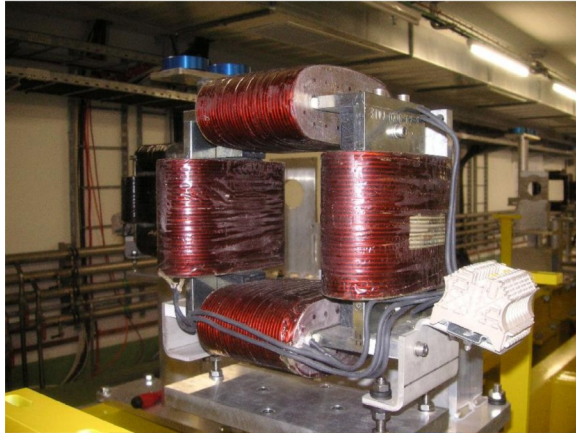
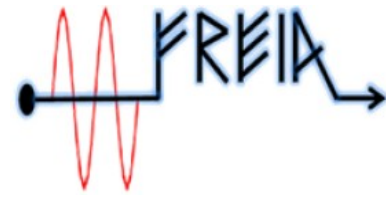
- \* HL-LHC dipoles
- \* CCT magnets



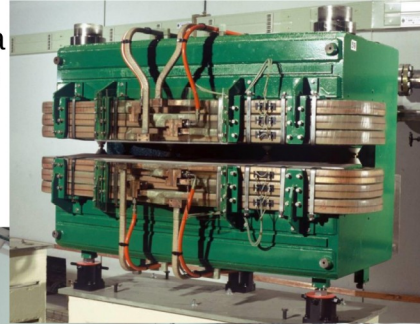


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# Magneter



Elettra  
dipole



from CERN web site



from CERN web site



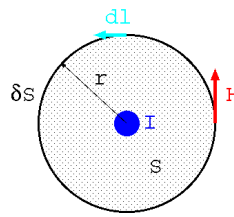
Photos: R. Ruber



from XFEL web site



# A little theory in 2D



Maxwell's equations lead to Ampere's law

$$\begin{aligned} \vec{\nabla} \cdot \vec{D} &= \rho, & \vec{\nabla} \times \vec{E} &= -\frac{\partial \vec{B}}{\partial t}, \\ \vec{\nabla} \cdot \vec{B} &= 0, & \vec{\nabla} \times \vec{H} &= \vec{J} + \frac{\partial \vec{D}}{\partial t} \end{aligned}$$

$\int_S \vec{J} \cdot d\vec{S} = \int_S (\vec{\nabla} \times \vec{H}) \cdot d\vec{S} = \int_{\partial S} \vec{H} \cdot d\vec{l}$   
 Stoke's theorem

$$B(r) = \frac{\mu_0 I}{2\pi r}$$

"current I makes field B"

In 2D derive this from **complex potential F(z)**

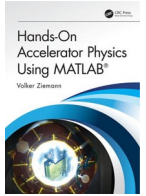
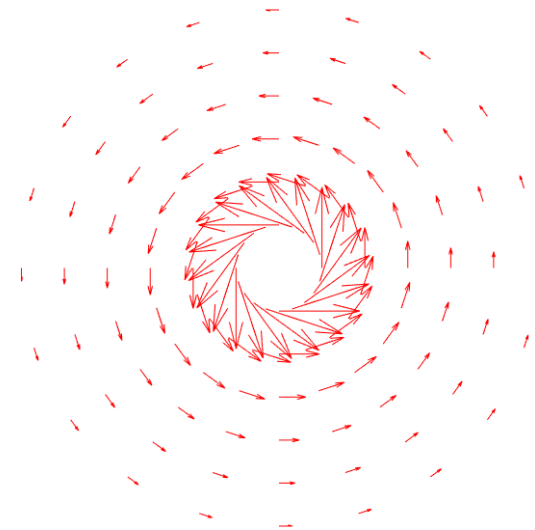
$$B_x - iB_y = \frac{\mu_0 I}{2\pi iz}$$

$$F(z) = (-\mu_0 I / 2\pi) \log(z) \longrightarrow B_x(x, y) - iB_y(x, y) = i \frac{dF(z)}{dz}$$

```

% field_of_current_filament.m
clear all; close all
r=0.2:0.2:1.2; phi=pi/10:pi/10:2*pi;
[R,PHI]=meshgrid(r,phi);
XX=R.*cos(PHI); YY=R.*sin(PHI)
B=1./(XX+i*YY);
By=real(B); Bx=imag(B);
quiver(XX,YY,Bx,By,'Color','r');
axis square; axis off;
xlim([-1.2,1.2]); ylim([-1.2,1.2])

```





# Multipoles

$$z = x + iy$$

- Expand any  $F(z)$  in a power series

$$F(z) = -B_0 R_0 \sum_{m=1}^{\infty} \frac{b_m + ia_m}{m} \left( \frac{z}{R_0} \right)^m$$

- with multipole coefficients  $a_m$  and  $b_m$

- Magnetic fields

$$iw(z) = B_y + iB_x = -\frac{dF}{dz} = B_0 \sum_{m=1}^{\infty} (b_m + ia_m) \left( \frac{z}{R_0} \right)^{m-1}$$

- $b_m$  describe upright magnets

- $a_m$  describe skew-magnets

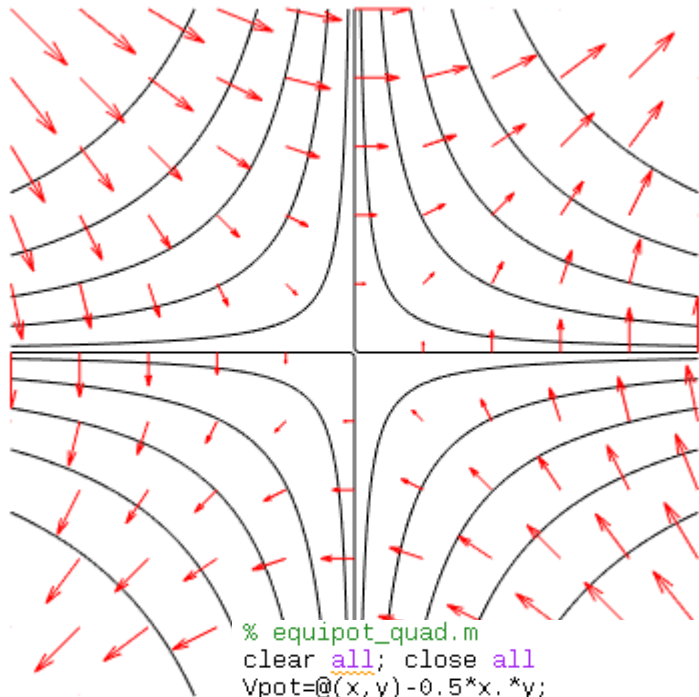
Accelerator people  
use this a lot!





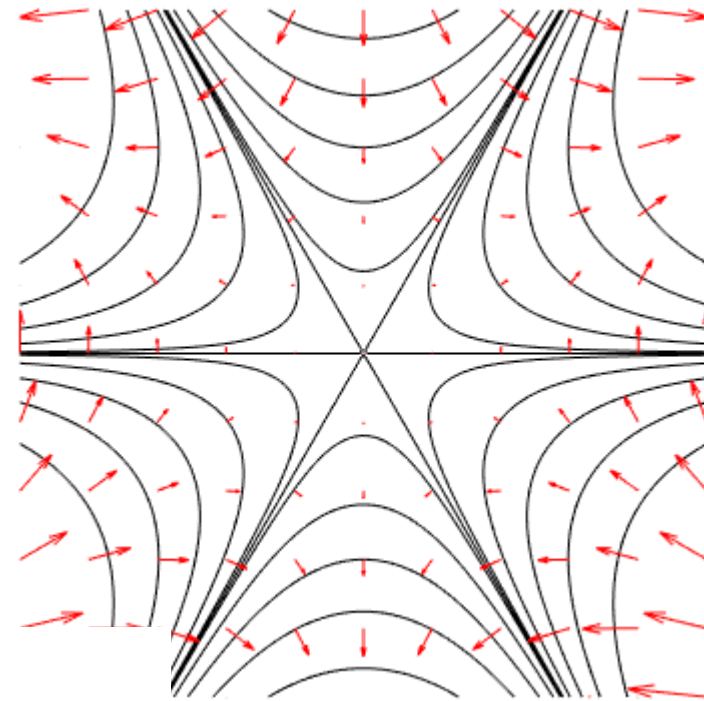
# Equipotential and field lines

- Quadrupole ( $m=2$ )



```
% equipot_quad.m
clear all; close all
Vpot=@(x,y)-0.5*x.*y;
[r=-5:0.05:5; [X,Y]=meshgrid(r,r); V=Vpot(X,Y); % Pot on fine grid
contour(X,Y,V,'LevelList',[0,quantile(V(:),10)],'Color','k');
r=-5:1:5; [XX,YY]=meshgrid(r,r); % Fields on coarse grid
V=Vpot(XX,YY); [Bx,By]=gradient(-V);
hold on; quiver(XX,YY,Bx,By,'Color','r');
axis square; axis off;
axis([-5,5,-5,5])
```

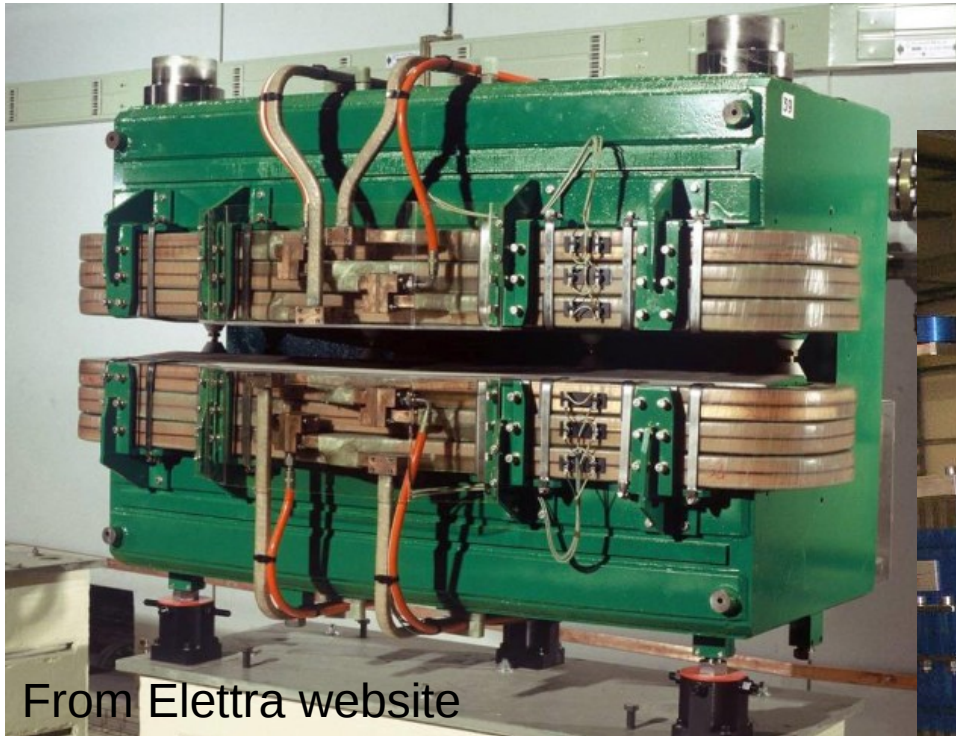
- Sextupole ( $m=3$ )



This gives us a first idea about the shape of the pole faces



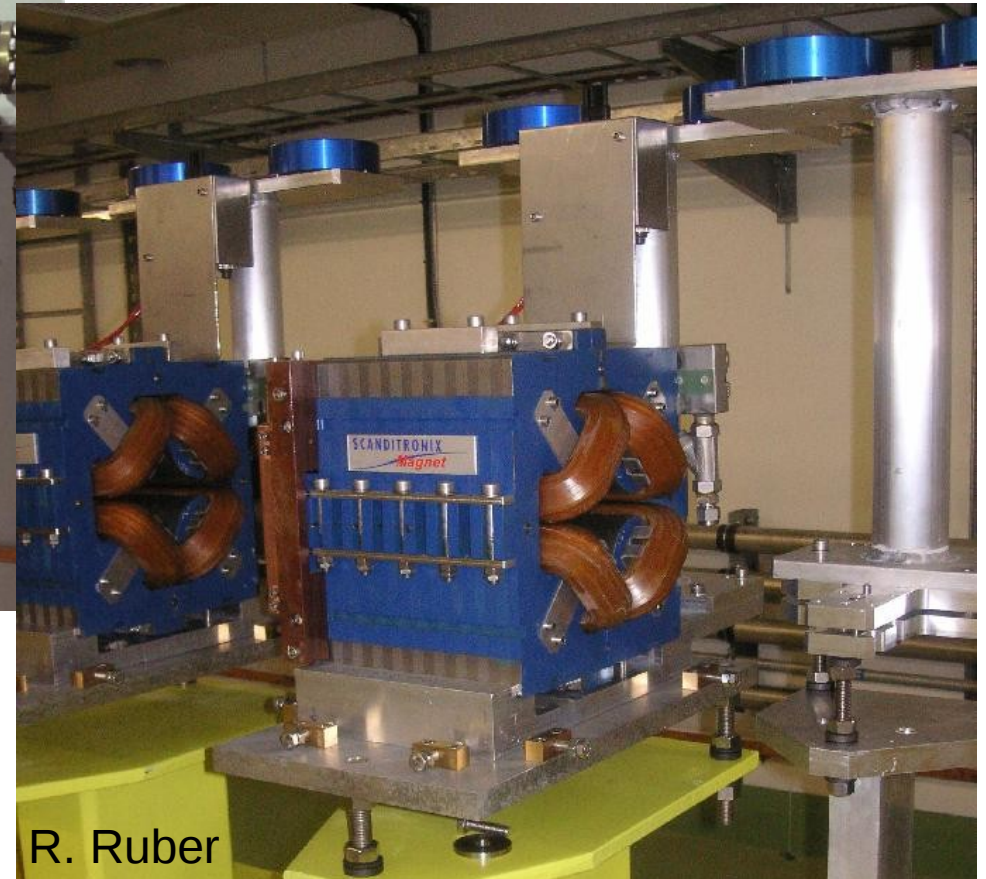
# Iron-dominated magnets



From Elettra website

Iron shapes the field

“Coils pump flux lines into the iron”



R. Ruber

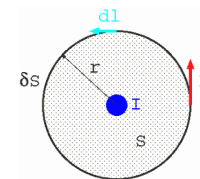




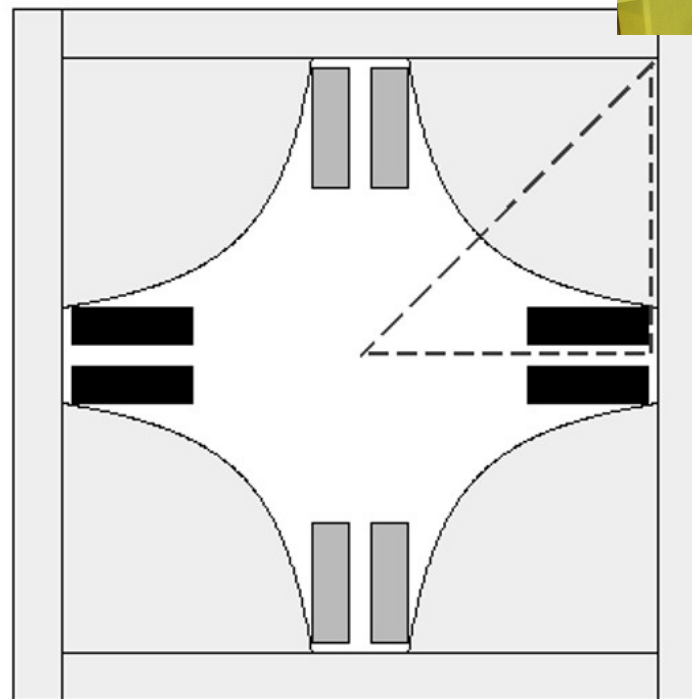
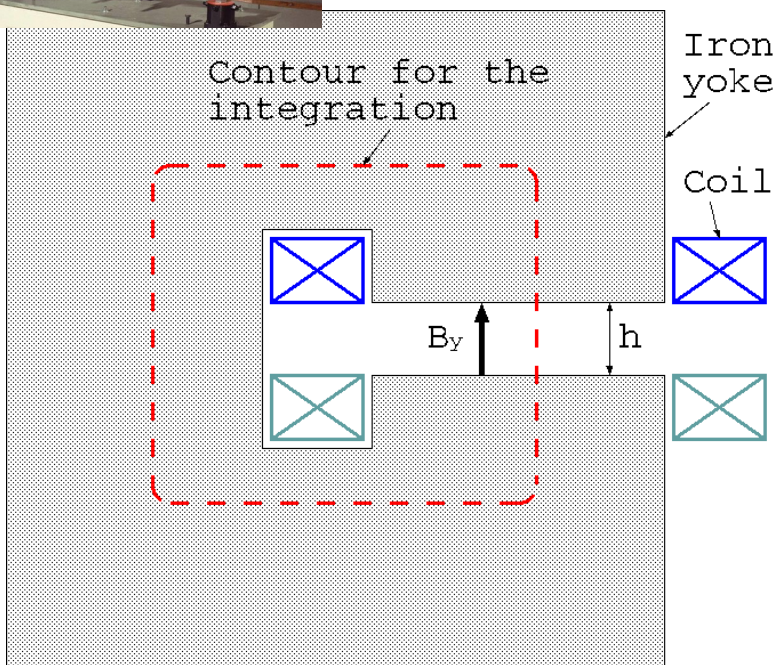
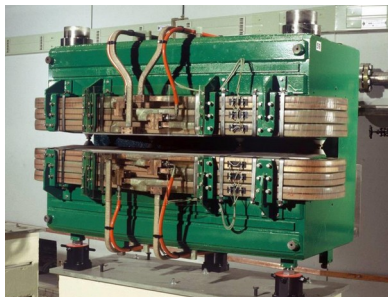
# Calculate with Ampere's law

Enclosed current gives  $Hdl$

$$\int_S \vec{J} \cdot d\vec{S} = \int_S (\vec{\nabla} \times \vec{H}) \cdot d\vec{S} = \int_{\partial S} \vec{H} \cdot d\vec{u}$$



Deflects the beam and  
treats all particles equally



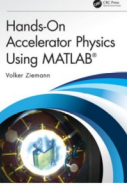
Magnetic lens  
focuses the  
particles

$$2NI = \int_{\text{gap}} H_y dy + \int_{\text{iron}} H_{\text{iron}} dl = \int_{\text{gap}} \frac{B_y}{\mu_0} dy + \int_{\text{iron}} \frac{B_{\text{iron}}}{\mu_0 \mu_r} dl \approx \frac{B_y h}{\mu_0}$$

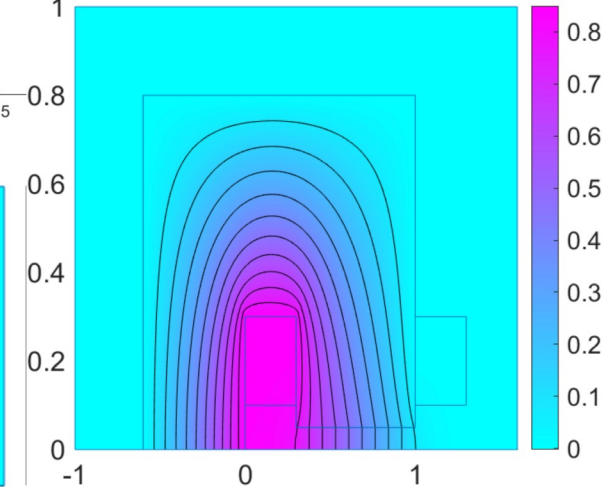
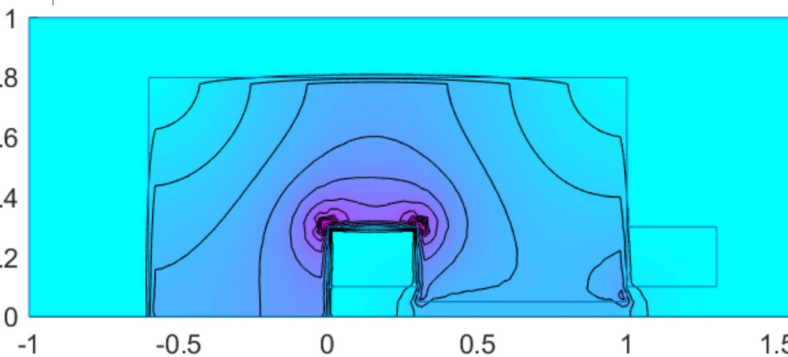
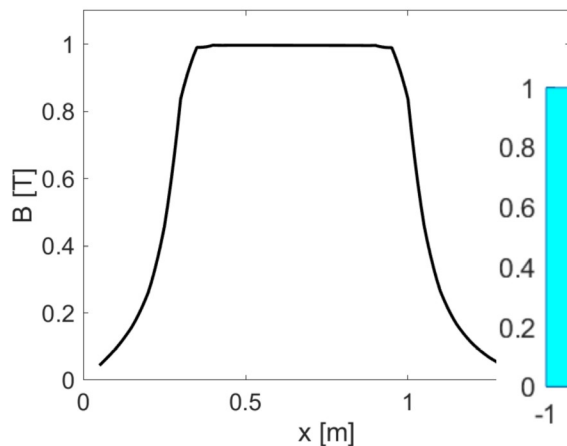
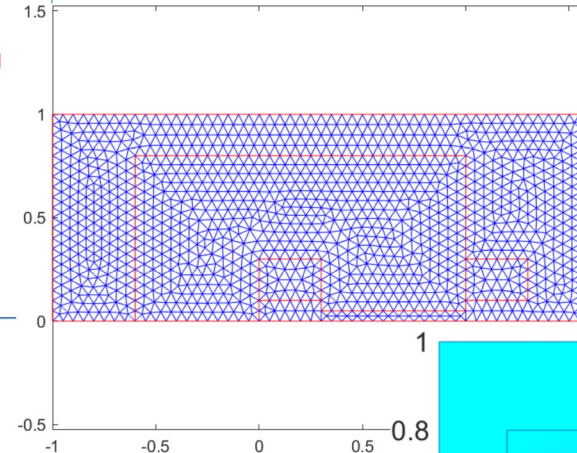
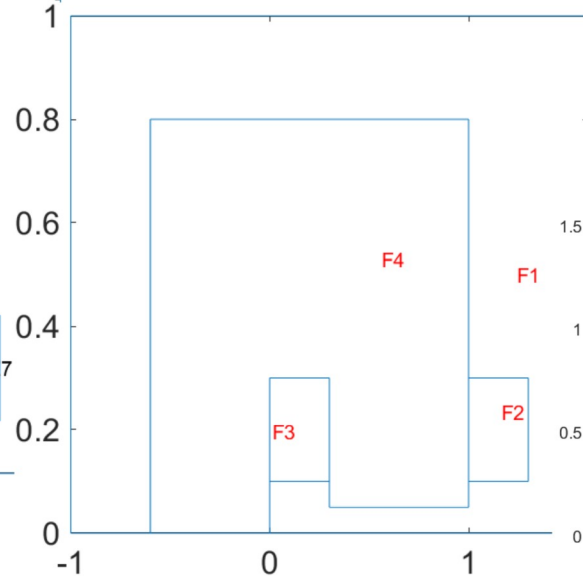
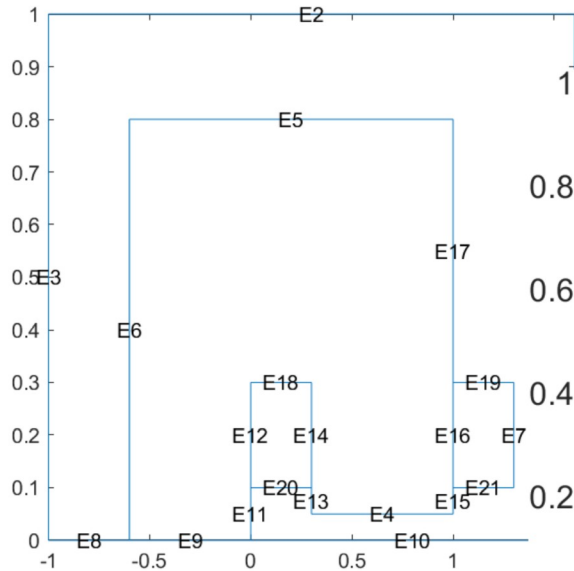
$$NI = \int_{\text{horiz}} \frac{B dl}{\mu_0} + \int_{\text{iron}} \frac{B dl}{\mu_0 \mu_r} + \int_0^a \frac{g r dr}{\mu_0} \approx \frac{g a^2}{2 \mu_0}$$



# Next step: numerical



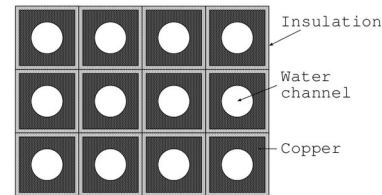
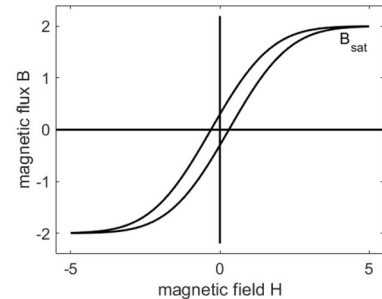
- \* Define geometry
- \* What carries current, what is iron?
- \* Boundary conditions
- \* Generate mesh
- \* Solve the system
- \* Post-processing





# Technological aspects

- Hysteresis and standardization
- Maximum field up to about 2 T.
- Eddy currents and laminated (transformer) steel
- Water cooling through holes in the conductors (up to  $\sim 10 \text{ A/mm}^2$ )
- Air-cooled magnets up to  $\sim 1 \text{ A/mm}^2$





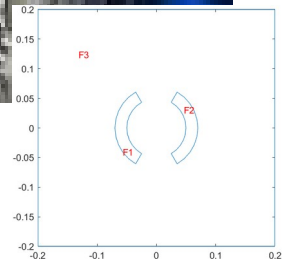
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# Current-dominated magnets

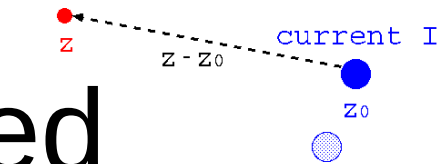
## super-conducting magnets

Required for fields exceeding 2 T such as the dipoles and quadrupoles for LHC

But also in high-field magnets for magnetic imaging



# Multipoles from a displaced current filament with



$$F(z) = \frac{\mu_0 I}{2\pi} \log(z - z_0)$$

Field

$$\underline{\hat{B}}^* = \hat{B}_x - i\hat{B}_y = \frac{\mu_0 I}{2\pi} \frac{i}{z - z_0} = \frac{-i\mu_0 I}{2\pi z_0} \sum_{n_0}^{\infty} \left( \frac{z}{z_0} \right)^n$$

Introduce cylindrical coordinates

$$z = r e^{i\phi}$$

$$z_0 = r_0 e^{i\phi_0}$$

$$\underline{\hat{B}}^*(\phi_0) = \frac{-i\mu_0 I}{2\pi r_0} \sum_{n_0}^{\infty} \left( \frac{r}{r_0} \right)^n e^{in\phi} e^{-i(n+1)\phi_0}$$

For a  $\cos(m\phi)$  current distribution

$$dI/d\phi_0 = \hat{I} \cos(m\phi_0 + \hat{\phi})$$

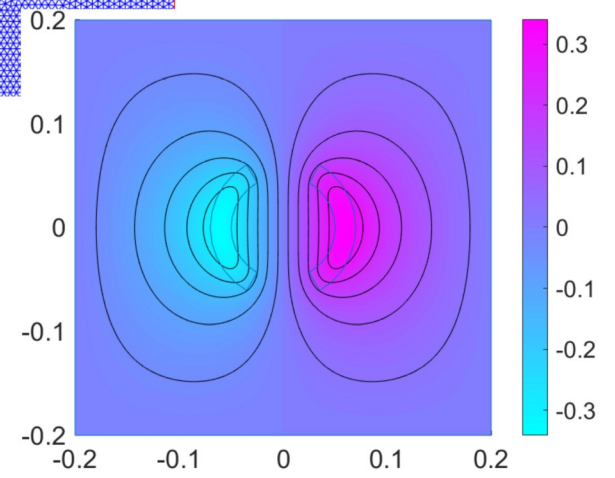
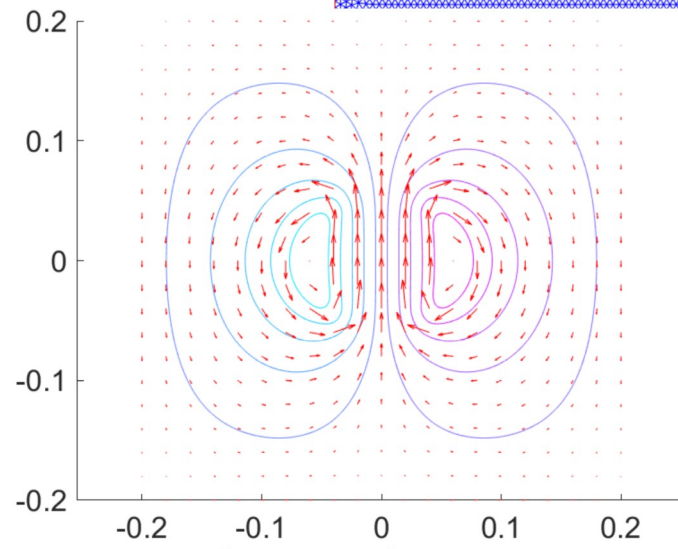
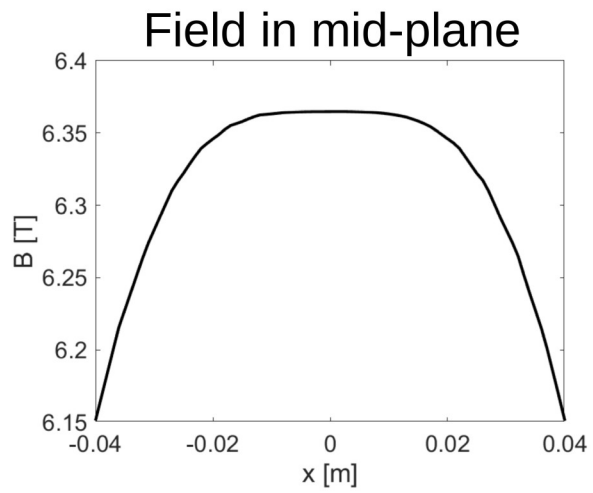
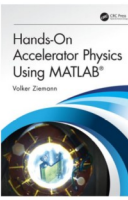
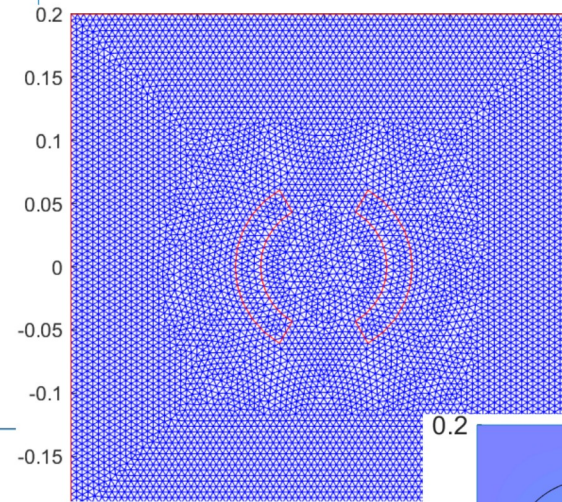
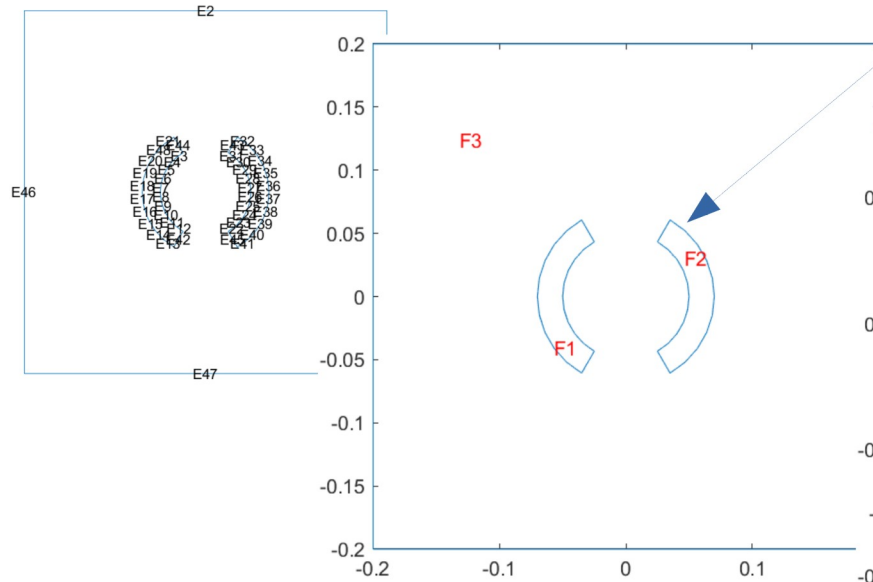
$$\underline{B}^* = \frac{-i\mu_0 \hat{I}}{2\pi r_0} \left( \frac{r}{r_0} \right)^n e^{in\phi} \int_0^{2\pi} e^{-i(n+1)\phi_0} \cos(m\phi_0 + \hat{\phi}) d\phi_0$$

Only non-zero for  $m=n+1 \rightarrow \cos((n+1)\phi)$  and  $r^n$



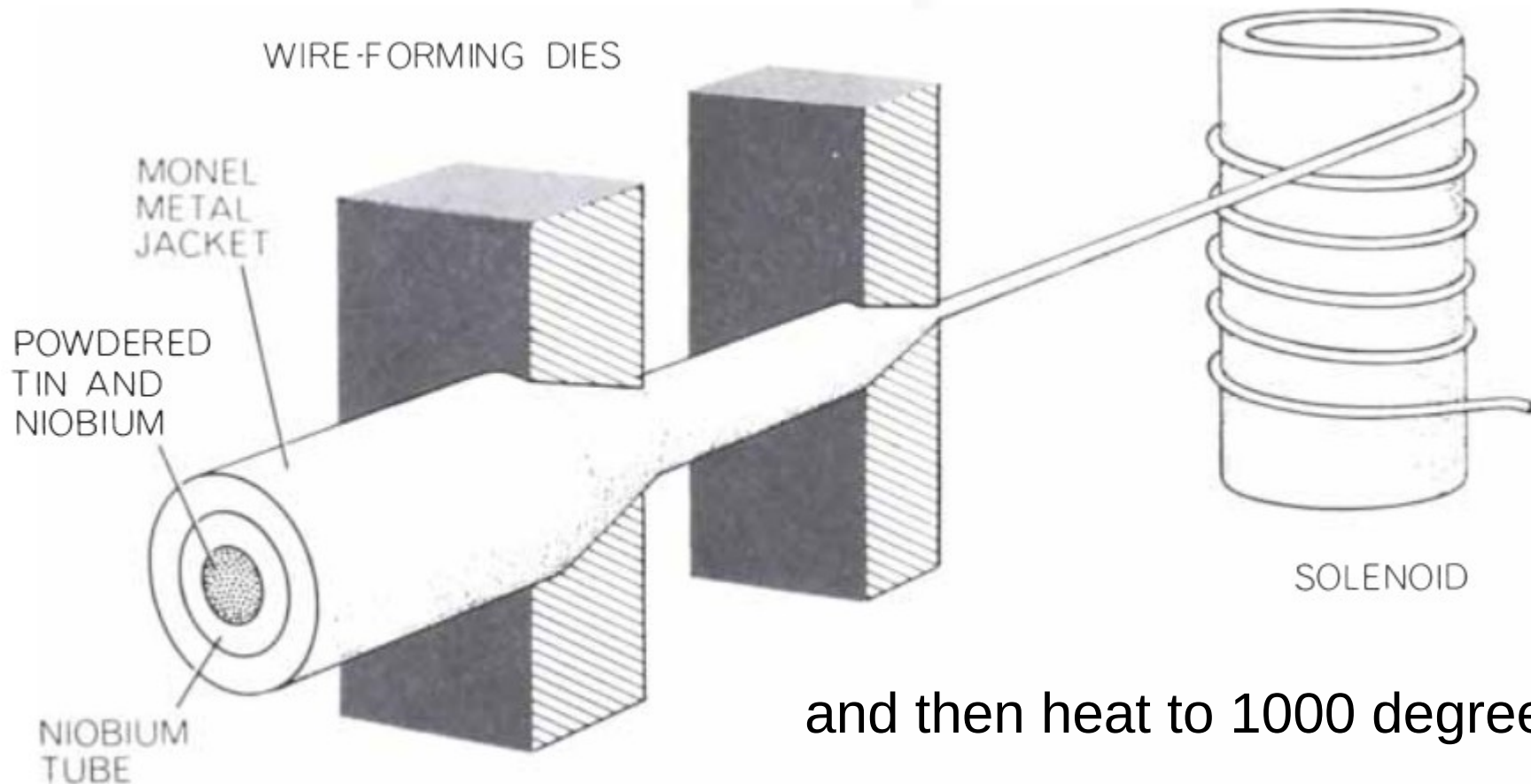
# Dipole with PDE toolbox

define a function that writes a  
circle segment for the geometry description





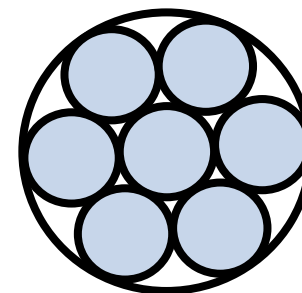
# Super-conducting cables



and then heat to 1000 degrees C

from Scientific American, June 1962

Your CCT cable...





# Technical aspects of super-conducting magnets

- Cooled with liquid helium at 4.2 or 1.9 K.
- Zero DC-conductivity.
- Type-II super-conductors
- reach higher fields  $H_{c2}$  and partially expel fields between  $H_{c1}$  and  $H_{c2}$ .
- Nb-Ti can carry  $J \sim kA/mm^2$ .
- Eddy-currents in strands limit ramp-speed.
- Forces between strands in the super-conducting cables cause friction and subsequent quenches.

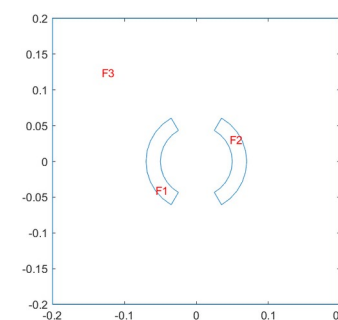
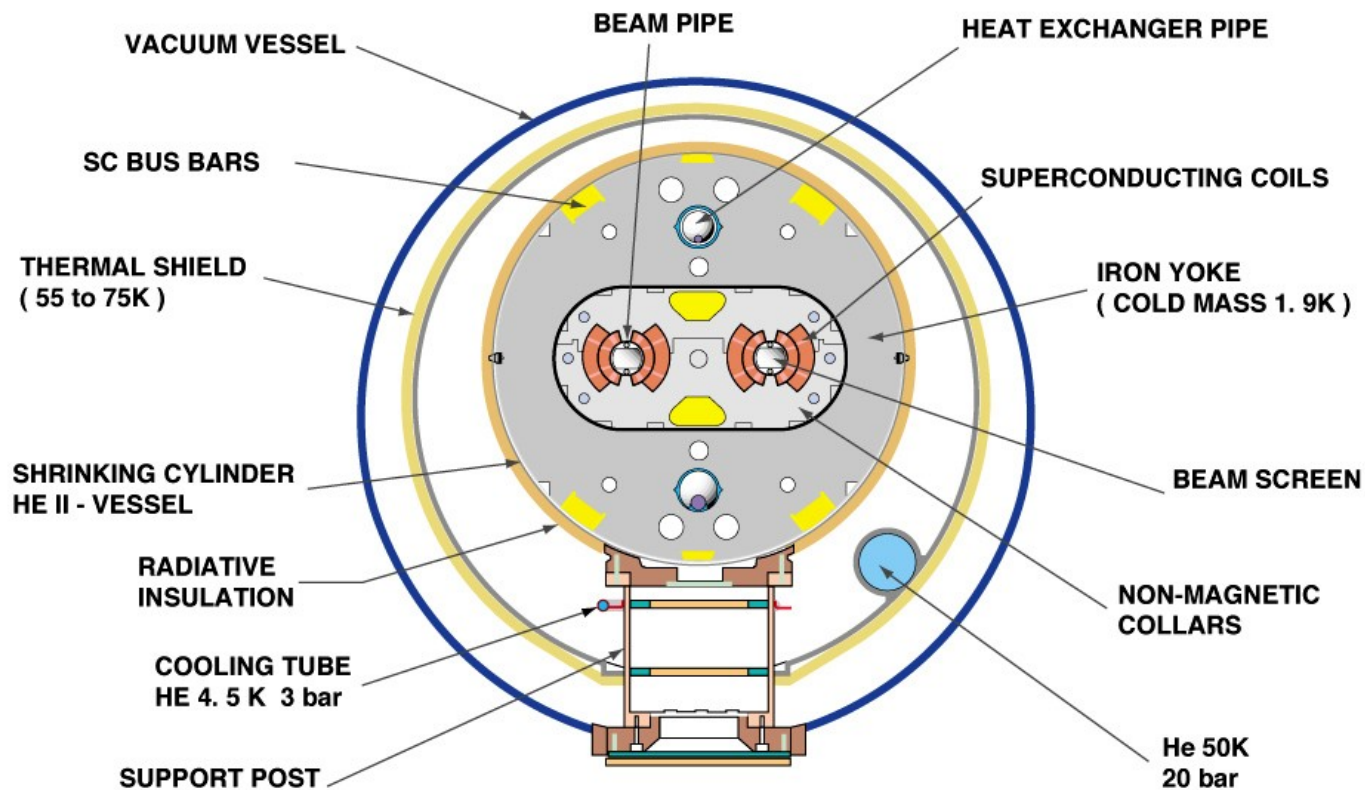




# LHC dipoles



## CROSS SECTION OF LHC DIPOLE



from CERN website

CERN AC\_HE107A\_V02/02/98



# Next time

- Short recap of today's material
- Permanent magnets
- Magnet alignment
- Magnetic measurements

